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# A lumped parameter model to evaluate the relevance of excess air as a tracer of exchanged flows between transmissive and capacitive compartments of karst systems

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The objective of this study was to evaluate the relevance of using excess air (EA) for the characterization of drain/matrix exchange in karst systems using a rainfall discharge model coupled with the simulation of EA measured at the outlet of the studied system. The conceptual model assumes a linear relationship between the formation of EA and the increase of hydrostatic pressure in the capacitive part of the aquifer. The simulated EA at the spring consists of the mixing of water circulating in the different compartments of the aquifer, with their own EA signature. The analysis is performed taking as an example the Durzon karst system (Larzac, France). The modeling is applied using daily rainfall discharge time series and 18 EA measurements at the main outlet of the karst system within 3 hydrological cycles. The main modeling results show that EA variations measured at the karst spring can be explained by recharge processes and exchange between conduit and matrix. EA measurements at the spring thus contain valuable information about the flow dynamics within the aquifer. Furthermore, results show that the use of EA measurements, despite their sparse temporal resolution, allows for reducing uncertainties in the estimation of some parameters of the reservoir model used for the simulation of karst spring discharge.

## KEYWORDS

karst hydrology, air excess, lumped parameter modeling, recharge, karst aquifer

## Introduction

Karstic systems consist of calcareous formations (carbonate rocks), often conceptualized through the co-existence of three types of porosity: (i) inter-granular porosity, (ii) fracture porosity, and (iii) large, wide conduits (Palmer, 1991). The relative proportions of each type of porosity in karst aquifers induce a strong spatial variability

in hydraulic conductivity (Jeannin, 2001) with significant implications for flow kinetics (Ford and Williams, 2013). This results in a heterogeneous permeability field and non-linear hydraulic behavior. The drainage in karst aquifers consists of a self-organized structure leading to a preferential flow path through a conduit network embedded in either a porous matrix which sometimes is highly permeable or a fissured matrix with low permeability (Ford and Williams, 2013). Preferential flow paths can be connected to the surface through sinkholes or stream losses, allowing rapid infiltration of surface water which may have an impact on groundwater chemistry (Aquilina et al., 2005) as well as imply an important risk of karst groundwater contamination (Hartmann et al., 2021; Lukač Reberski et al., 2022).

While the drainage system controls contaminant transport within the aquifer, the matrix may store contaminants because of lateral exchange between the preferential flow path and the surrounding matrix. Karst aquifers are therefore highly vulnerable because of rapid contaminant spread and potential storage in the matrix over time. Understanding and considering lateral exchange between these two kinds of karst porosities seems to be necessary for proper qualitative (chemical vulnerability) and quantitative (resource availability) groundwater resource management of karst aquifer.

Cholet et al. (2017) proposed a coupled deconvolution of spring hydrograph and mass chemograph to assess lateral contributions in terms of flow and mineralization in a conduit-flow-dominated karst system. The study shows that low-mineralized water flow may occur from conduit to matrix after a major precipitation event in the non-saturated zone. Mitrofan et al. (2015) also highlighted flood-induced flow from the karst conduit to the surrounding matrix evidenced by chloride fluxes. Martin and Dean (2001) showed that water exchange from the conduit to the matrix may also have significant implications in karstification processes due to the intrusion of water undersaturated concerning calcite into the matrix. Numerical modeling may also bring insight into lateral exchange dynamics (Bailly-Comte et al., 2010; Reimann et al., 2011; Sivelle et al., 2019; Dal Soglio et al., 2020; Shirafkan et al., 2021). The contribution of exchange from slow-flow to fast-flow systems in three karst catchments located in southern China was estimated between 64 and 87% (Zhang et al., 2017). Within the Norville chalk aquifer (northern France), the fluxes from matrix to conduit represent approximately 10% of the spring discharge (Duran et al., 2020). Within the highly karstified Baget and Aliou watersheds (French Pyrenees), the annual matrix contribution to the total spring flow is ~3% and it can increase up to 25% during periods with low rainfall (Sivelle et al., 2019). Also, the numerical modeling of exchange flow may benefit from auxiliary variables such as temperature (Berglund et al., 2020) or  $\text{Ca}^{2+}$  and  $\text{Mg}^{2+}$  ions concentration (Zhang et al., 2017) for model calibration. Another approach consists in analyzing simulation results together with physicochemical properties

such as electrical conductivity (Duran et al., 2020), to better evaluate the physical significance of simulated internal fluxes.

Among these, the study of dissolved noble gases in the subsurface may be of interest to improve the understanding of groundwater dynamics in karsts by providing information on flow path and connectivity between aquifer components. Historically, noble gases have been used as environmental tracers for a variety of purposes, such as reconstructing paleoenvironmental and paleoclimatic conditions or identifying recharge zones (Heaton and Vogel, 1981; Wilson and McNeill, 1997; Aeschbach-Hertig et al., 2000; Peeters et al., 2002; Ingram et al., 2007; Sano et al., 2017; Tomonaga et al., 2017; Popp et al., 2021; Schilling et al., 2022). While the solubility equilibrium of atmospheric gases is primarily controlled by the temperature of the recharge zone, many studies indicate that noble gas concentrations in groundwater often differ from these atmospheric equilibria and appear to be modified by another ubiquitous phenomenon called “excess air.” The work of Holocher et al. (2002) was able to show experimentally that the formation of excess air can be explained by the complete or partial dissolution of trapped air bubbles under the influence of hydrostatic pressure (Heaton and Vogel, 1981; Stute and Schlosser, 2000; Aeschbach-Hertig, 2004). Therefore, excess air can be considered as a tracer that can provide additional useful information on the type of recharge, that is, saturated or unsaturated and the amplitude of groundwater level fluctuations (Clark et al., 2005; Massmann and Sültenfuß, 2008; Schilling et al., 2021). The excess air is usually expressed as a volume under standard temperature and pressure (STP) conditions, the unit is  $\text{cm}^3 \text{ STP.kg}^{-1}$ . The work of Pérotin et al. (2021) showed on the Durzon karst with a large vadose zone (>200 m) that water transfer along the vadose zone causes a slight excess of air (+ 0.8  $\text{cm}^3 \text{ STP.kg}^{-1}$ ) that remains small compared to that observed at the Durzon spring (+ 5.5  $\text{cm}^3 \text{ STP.kg}^{-1}$ ), suggesting that the excess air is mainly formed by variations in the water level of the saturated zone of this karst.

In this study, we investigated the relevance of using EA to characterize water exchange between transmissive and capacitive compartments in the Durzon karst system. Using a lumped parameter rainfall discharge model, the EA was simulated assuming formation in the capacitive part of the aquifer, when significant recharge occurs, and by addition from exchange fluxes between the transmissive and capacitive parts of the Durzon karst system. We first present the geological, hydrogeological, and meteorological context of the study area. Then, the formation of EA, as well as the information about recharge conditions provided by noble gas concentration measurements are described. Finally, the lumped parameter rainfall discharge model and the way to consider EA in the model are explained, before presenting the calibration results considering both spring discharge and EA and their implication in terms of parametric uncertainty of the considered lumped parameter rainfall discharge model.

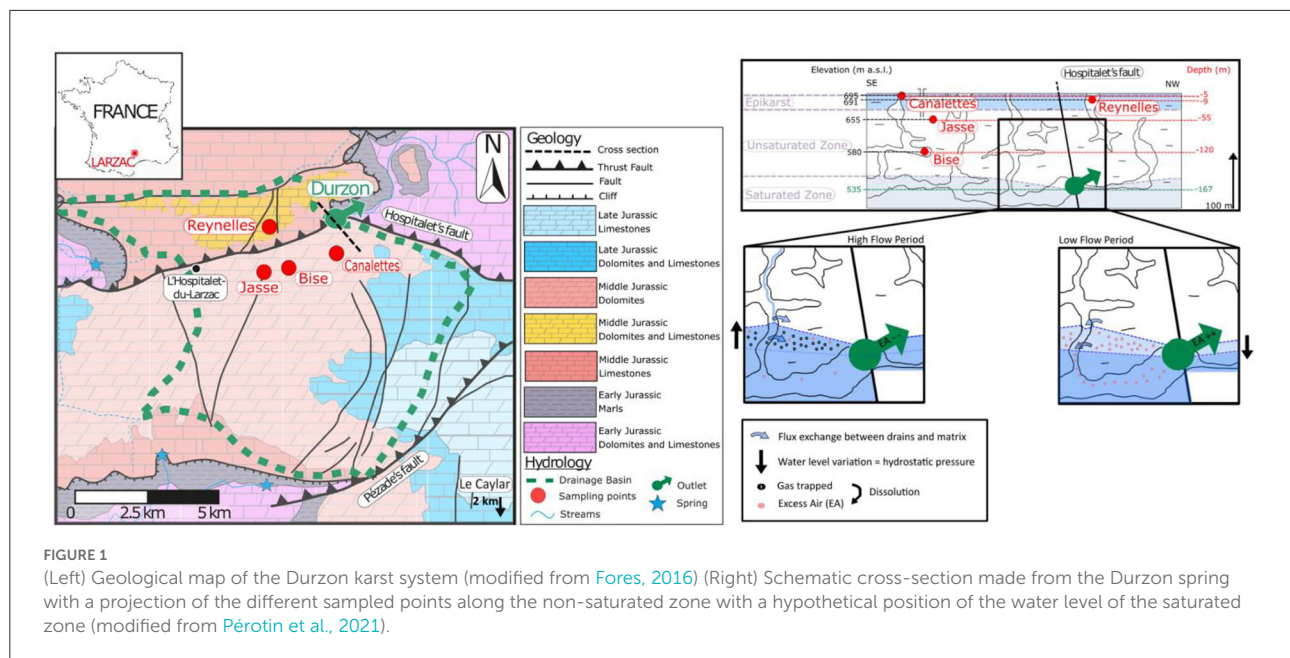


FIGURE 1

(Left) Geological map of the Durzon karst system (modified from Fores, 2016) (Right) Schematic cross-section made from the Durzon spring with a projection of the different sampled points along the non-saturated zone with a hypothetical position of the water level of the saturated zone (modified from Pérotin et al., 2021).

## Study area

### Local geology and hydrogeology

The Durzon karst system is located in the Grand Causses area in the Southern Massif Central (France), with a mean altitude of approximately 700 m asl, varying between 500 and 900 m asl (Fores, 2016). The study area is part of the Natural Regional Park of the Grand Causse (PNRGC). It is one of the study sites of the French National Observation Services H+ (<http://hplus.ore.fr/>), which is part of the Elter European Infrastructure (<https://www.lter-europe.net/elter-esfri>). The geology in the area consists of Jurassic karstified carbonates: The Hettangian carbonate series is separated from the Dogger and Malm carbonate series by a Toarcian marl series, which is considered an impervious geological layer (Figure 1). The marly formations play an important role in the hydrogeological functioning of the hydrogeological system by separating the two main aquifers, which are the two carbonate units mentioned above. The Hospitalet thrust puts the Lias marl into contact with the Dogger–Malm aquifer and allows the emergence of the spring (Figure 1). The Durzon spring is the only known outlet of the unary karst system, that is, the impluvium is only composed of karstifiable rocks and consists of a Vaclousian spring, that is, the drainage is organized below the level of the spring (Marsaud, 1996). The spring is located at an altitude of 533 m asl and drains a recharge basin whose area, estimated from mass balances, varies between 98 and 117 km<sup>2</sup> (Ricard and Bakalowicz, 1996; Tritz et al., 2011; Fores, 2016). The spatial limits of the basin were determined based on artificial tracer tests and

geological features that impose the boundary at the interface with impermeable marls.

The system is characterized by a 200 m thick infiltration zone and no borehole over the area reaches the saturated zone. Therefore, only the spring discharge variation can bring the required information for hydrodynamic characterization of the saturated zone.

### Climate

The area is characterized by a Mediterranean climate and oceanic influence. The annual cumulative rainfall varies from 929 to 1,729 mm (2007–2019) with a mean annual rainfall of  $\sim 1,180$  mm distributed over  $\sim 92$  rainy days per year (rainfall  $> 1$  mm.day<sup>-1</sup>). The mean air temperature varies between 2.5°C in winter and 18°C in summer, with a mean annual air temperature around 10°C. The mean annual potential evapotranspiration (PET) computed using the Penman–Monteith's formula (Penman, 1948) by the national meteorological agency MétéoFrance at the Caylar station varies from 830 to 1,020 mm (2007–2019) with a mean annual PET of approximately 927 mm. The mean PET is lower than 1 mm.day<sup>-1</sup> during winter and rises to approximately 5 mm.day<sup>-1</sup> during the summer period. Autumn is characterized by intense rainy episodes called “Cevenol episodes” where cumulative precipitation amount may reach 286 mm.day<sup>-1</sup> (17 September 2014). Such meteorological events greatly contribute to the aquifer recharge and could generate flooding, the dynamics of which depend on the hydric

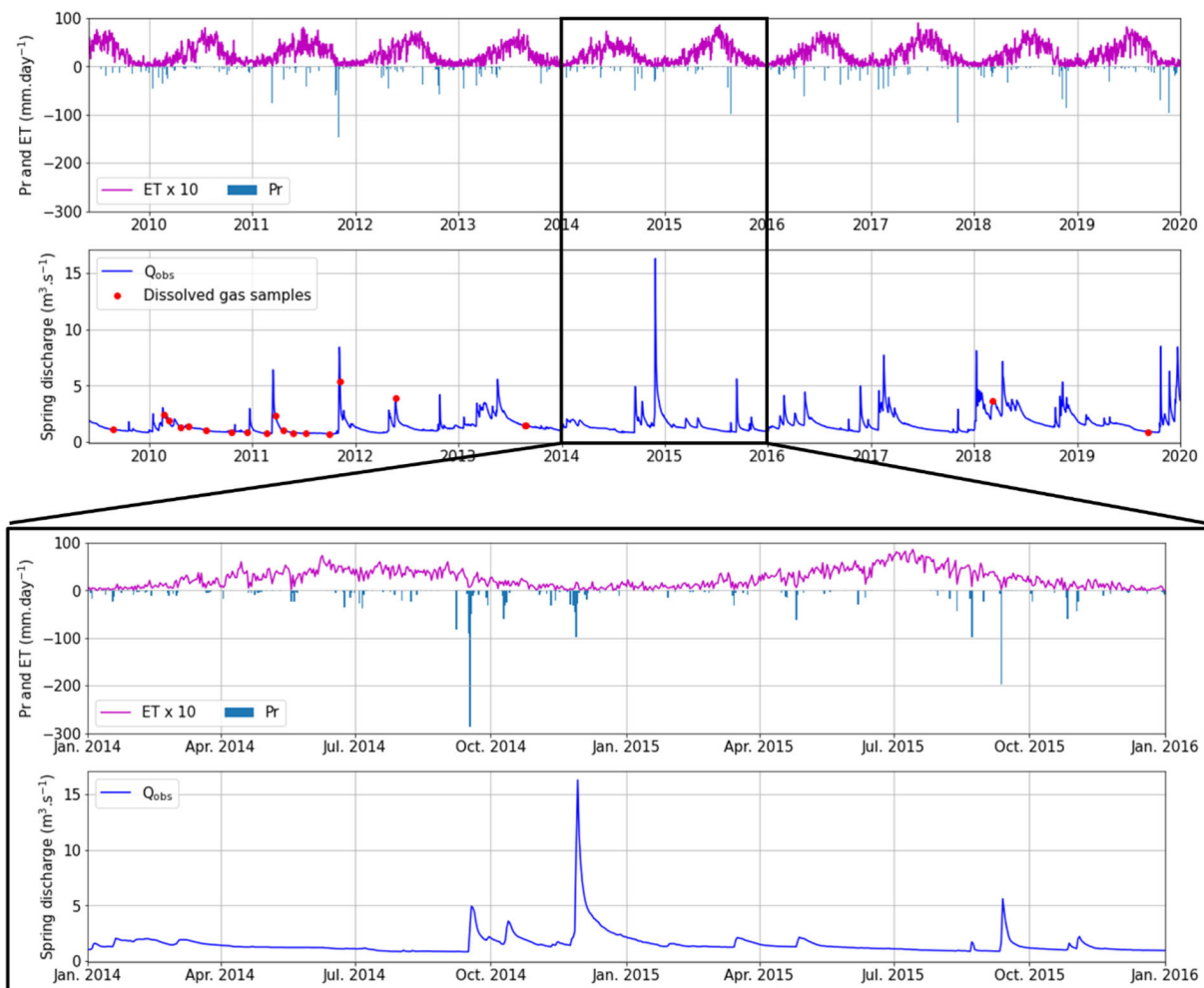


FIGURE 2

(Top) Precipitation (Pr) and evapotranspiration (ET) time series measured at the Caylar station and discharge time series measured at the Durzon spring and sampling of dissolved gases—(Bottom) Zoom on the period from 2014 to 2016.

state of the aquifer. As an example, the 2014 cevenol episode occurred at the end of a long recession period and generated a maximum spring discharge of only  $5 \text{ m}^3 \cdot \text{s}^{-1}$  at the Durzon spring. The following rainfall episode, much less intense with a maximal rainfall intensity of approximately  $100 \text{ mm} \cdot \text{day}^{-1}$ , induced a maximum spring discharge of  $16 \text{ m}^3 \cdot \text{s}^{-1}$  (Figure 2).

## Materials and methods

### Air excess formation

Excess air results from the trapping of air bubbles in a porous medium following a significant recharge event. The hydrostatic pressure then applied to the air bubbles forces their dissolution causing an influx of air larger than what can be expected from

equilibrium exchange between air and water (Holoche et al., 2002; Aeschbach-Hertig, 2004; Klump et al., 2007). EA was calculated based on Ne-Ar-N<sub>2</sub> concentrations, as N<sub>2</sub> is neither produced nor consumed in significant amounts at the studied site (Heaton and Vogel, 1981; Plummer et al., 2004; Chatton, 2017). Due to its formation process, EA can be used as a qualitative tracer for piezometric head variation in the aquifer.

Wilson and McNeill (1997) investigated the relationship between air bubble entrapment and lithology by comparing mean EA values (calculated based on Ne isotopes) in groundwater samples from aquifers developed in limestones, sandstones, and granites. Samples from granitic aquifers were almost devoid of EA, while samples from carbonate aquifers had the highest EA values ( $\sim 1.0 \text{ cm}^3 \text{ STP} \cdot \text{kg}^{-1}$ ). Several authors (Herzberg and Mazor, 1979; Heaton and Vogel, 1981; Mazor et al., 1983) reported high EA values in karst aquifers. Several



conceptual models have been proposed to explain the EA formation during recharge. According to [Heaton and Vogel \(1981\)](#), the presence of air bubbles trapped under the action of capillary forces in the capillary fringe (microscopic scale) could explain the EA formation. When surface water infiltrates, the air bubbles are carried to the level of the saturated zone. When sufficient hydrostatic pressure is applied to the trapped air bubbles, they are dissolved within the groundwater. [Herzberg and Mazar \(1979\)](#), suggested that the highly permeable fractured nature of most of the karst aquifers is often associated with rapid recharge that could be responsible of EA formation at a macroscopic scale. Carbonates thus appear to be suitable lithologies for the EA formation, though the physical process remains hypothetical.

[Aeschbach-Hertig et al. \(2002\)](#) suggest that variations in Neon concentration (and thus EA) can be a proxy of the piezometric head variation in the aquifer. Such statement is supported by many other studies ([Beyerle et al., 2003](#); [Kulongoski et al., 2004](#); [Ingram et al., 2007](#); [Manning and Caine, 2007](#); [Osenbrück et al., 2009](#); [Zhu and Kipfer, 2010](#)). [Holocher et al. \(2002\)](#) experimentally investigated EA formation in quasi-saturated porous media using vertical sand columns. When simulating water level variations, a correlation between these variations and the measured EA was observed. Then EA varied between 0.14 and 1.62 cm<sup>3</sup> STP.kg<sup>-1</sup> with a maximal piezometric head variation of 1 m. According to [Holocher et al. \(2002\)](#), the total pressure exerted on the entrapped air is the dominating parameter responsible for the total amount of dissolved air in groundwater. The EA formation would thus depend on the number of trapped air bubbles and on the hydrostatic pressure that forces their dissolution. In some carbonate rocks such as dolomites, the porous matrix appears as a suitable media for the air bubble trapping considering the inter-granular porosity and significant and fast piezometric head variation with rapid transfer from the surface to the saturated zone. For simplicity, and to avoid over-parameterization in the modeling approach (see section Including air excess in the lumped parameter model), it is assumed that the evolution of excess air is linearly related to piezometric fluctuations.

### Assessing recharge conditions with noble gas temperature and excess air

Groundwater recharge temperature (NGT, Noble Gas Temperature) and EA ([Heaton and Vogel, 1981](#); [Plummer et al., 2004](#); [Chatton, 2017](#)) can be estimated using the concentration of dissolved noble gases (Ne, Ar) and nitrogen (N<sub>2</sub>). The latter is a major atmospheric gas, which can furthermore be produced by biological reactions such as denitrification in anoxic systems ([Mariotti, 1986](#); [Smith et al., 1991](#); [Blicher-Mathiesen et al., 1998](#)). In karst systems, anoxic conditions are rarely observed, especially in an unconfined aquifer. Considering the Durzon karst system, the measured dissolved oxygen is always above

7.6 mg.l<sup>-1</sup>, which excludes the existence of denitrification. For this reason, the concentration of dissolved N<sub>2</sub> is assumed to remain stable during the hydrological cycle and was used to calculate both NGT and EA.

Estimations were performed using the Unfractionated excess Air (UA) by the complete dissolution of air bubbles conceptual model, following the methodology described by [Aeschbach-Hertig and Solomon \(2013\)](#). This model does not include degassing or gas partitioning. The model equations developed by [Aeschbach-Hertig and Solomon \(2013\)](#) are not detailed in this article. The complete methodology in estimation of EA is presented in former studies ([Pérotin, 2021](#); [Pérotin et al., 2021](#)). The average recharge elevation was set at 700 m asl. The variability in elevation of the study area induces < 1°C and 0.2 cm<sup>3</sup> STP.kg<sup>-1</sup> of uncertainty in the NGT and EA parameters, respectively.

### Sampling and analysis

Water samples were collected at various depths within the vadose zone of the study site (see [Pérotin, 2021](#)). Canalettes and Reynelles samples taken at 5 and 9 m depth (691 and 685 m asl, respectively) allow characterization of the epikarst zone, and a 45 m deep borehole (655 m asl) allows characterization of water transfer in the unsaturated zone (Jasse). An underground river at 120 m depth in the unsaturated zone (Bise, 580 m asl) allows characterizing in a more integrative way the flows within the vadose zone, this perched underground river is located nearly 55 m above the level of the Durzon spring (535 m asl). Two sampling campaigns were undertaken during high and low water conditions (March 2018 and Sept 2019, respectively) to evaluate the seasonal hydrological influence on the gas transfers. Two sampling campaigns were conducted during high and low water periods (March 2018 and Sept 2019, respectively) to characterize dissolved gas signatures within the vadose zone. Dissolved gas samples were collected at Durzon Spring from 2009 to 2019 with a monthly time step between February 2010 and November 2011.

The noble gases dissolved in water were collected in 500 ml glass bottles closed with rubber stoppers and sealed with a metal ring. A small submerged pump (Mini Twister 12 volts standard PVC - sdec-France) was used to sample the waters of the Reynelles cave (9 m deep) and the monitoring was carried out at the Durzon spring. At Canalettes cave (5 m deep), a peristaltic pump was required due to the very low water quantity in the natural pool. The sampling is done directly in the cave, so the pump and the tank containing the water were at the same altitude to avoid any loss of pressure and therefore any decompression in the pump. Considering the Jasse borehole (45 m deep), a Grundfos MP1 pump was used, and the dissolved gas samples were collected underwater (directly in the well or in a bucket) to avoid air contamination, after having renewed the water 3 times. Two or three replicates were collected for

each sample. Also, physicochemical parameters (temperature, pH, specific conductivity, and dissolved oxygen) were recorded using a HACH HQ40D portable multimeter for each dissolved gas sample. The groundwater temperature at the Jasse borehole was measured with a CTD Diver located in the borehole, rather than after sampling.

Dissolved gases (Ne, Ar, N<sub>2</sub>) of the water samples were analyzed by the analytical platform CONDATE Eau (Rennes University) with the gas chromatography method. Ne-Ar-N<sub>2</sub> gases were extracted by the head-space method (Sugisaki and Taki, 1987) and measured by gas chromatography using a catharometric detector ( $\mu$ GC 3000 - SRA), with an accuracy of 8% for Ne and 5% for Ar and N<sub>2</sub>.

## Modeling approach

### Lumped parameter modeling

The proposed model (Figure 3) considers three compartments, each corresponding to a physical entity (epikarst, conduit, matrix). The compartments are organized as a two-level structure: the upper level of the model structure comprises the compartment E (epikarst) while the lower level comprises two interconnected compartments M (matrix) and C (conduit). As proposed in a former study (Tritz et al., 2011), the compartment E represents the soil and epikarst subsystem and constitutes the interception zone where part of the received precipitation  $Pr$  goes back to the atmosphere *via* an evapotranspiration rate  $ET$ . The water contained in compartment E flows to the lower compartment M. During high water level periods, part of the water contained in compartment E produces a hysteretic overflow  $Q_{hy}$  (Tritz et al., 2011) reaching the compartment C or the spring. The flow repartition for  $Q_{hy}$  to compartment C or the spring is parameterized by the sharing coefficient  $X_{hy}$ . In the lower part of the model structure, the exchange between compartments M and C is proportional to the water level difference between the two compartments. It allows the reproduction of matrix-conduit interactions as supported by direct observations in other karst systems (Bailly-Comte et al., 2010). The model can be implemented with the KarstMod modeling platform, developed within the framework of the Karst observatory network SNO KARST ([www.sokarst.org](http://www.sokarst.org)), initiated by the French institute INSU/CNRS (Jourde et al., 2018). The model is based on the mass balance equations given by Mazzilli et al. (2019), assuming a linear discharge law for each compartment of the model:

$$\frac{dE}{dt} = Pr - ET - Q_{EM} - Q_{hy} \quad (1)$$

$$\frac{dM}{dt} = Q_{EM} - Q_{MC} \quad (2)$$

$$\frac{dC}{dt} = (X_{hy} \times Q_{hy}) + Q_{MC} - Q_{CS} \quad (3)$$

where:

$$Q_{EM} = k_{EM} \times \frac{hE(t)}{L_{ref}} \text{ if } E(t) > 0, \text{ otherwise } Q_{EM} = 0 \quad (4)$$

$$Q_{MC} = k_{MC} \times \frac{[hM(t) - hC(t)]}{L_{ref}} \quad (5)$$

$$Q_{CS} = k_{CS} \times \frac{hC(t)}{L_{ref}} \quad (6)$$

$$Q_{hy} = \varepsilon_{hy} \times k_{hy} \times \left( \frac{hE - E_{hy}}{L_{ref}} \right) \quad (7)$$

where  $hE$ ,  $hM$ , and  $hC$  are the water levels [mm] in the compartments E, M, and C, respectively,  $k_{EM}$ ,  $k_{MC}$ ,  $k_{CS}$ , and  $k_{hy}$  [mm.day<sup>-1</sup>] are the specific discharge coefficients for the discharge  $Q_{EM}$ ,  $Q_{MC}$ ,  $Q_{CS}$ , and  $Q_{hy}$  [mm.day<sup>-1</sup>], respectively,  $X_{hy}$  [-] is the sharing coefficient that splits the fluxes  $Q_{hy}$  to compartment C or directly to the spring,  $L_{ref}$  is a unit length [mm] and  $\varepsilon_{hy}$  [-] is the binary indicator of the activation of the hysteretic discharge (Tritz et al., 2011).

The simulated spring discharge  $Q_{sim}$  [m<sup>3</sup>.s<sup>-1</sup>] depends on the recharge area such as:

$$Q_{sim} = [Q_{CS} + (1 - X_{hy}) \times Q_{hy}] \times \frac{Ra}{(10^3 \times 86400)} \quad (8)$$

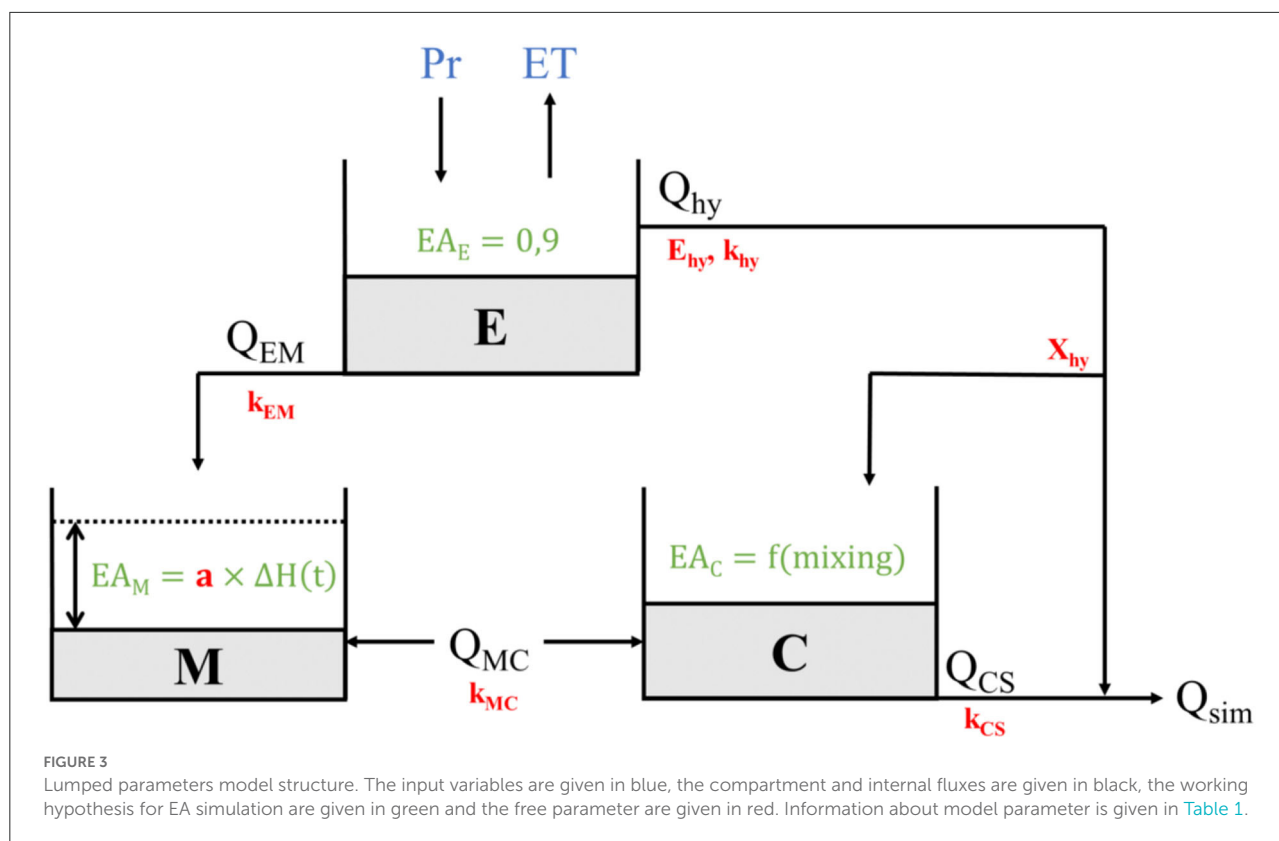
where  $Ra$  [km<sup>2</sup>] is the catchment recharge area divided by a constant for unit constancy to convert the discharge per unit surface [mm.day<sup>-1</sup>] into discharge [m<sup>3</sup>.s<sup>-1</sup>].

### Including air excess in the lumped parameter model

Including EA in the proposed model considers that the EA formation mainly occurs within compartment M as a result of an increasing piezometric head. Indeed, the EA formation occurs by trapping air bubbles in a porous medium, which are dissolved when hydrostatic pressure significantly increases. In the case of the studied karst system, the relatively high porosity of dolomite formation allows considering this hypothesis. To represent such a physical process, the model assumes that the EA formation in compartment M is proportional to the piezometric head variations:

$$EA_M^*(t) = a \times \Delta h(t) \text{ if } \Delta h(t) > 0 \text{ else } = 0 \quad (9)$$

where  $EA_M^*(t)$  is the concentration of EA [cm<sup>3</sup> STP.kg<sup>-1</sup>] in compartment M at time  $t$ , resulting from a piezometric head variation  $\Delta h(t)$  at time  $t$ , depending on the proportional coefficient  $a$ . If  $\Delta h(t)$  is negative, there is no EA formation and  $EA_M^*(t) = 0$ . The EA signature in compartment M depends



on the mass balance between inflow from compartment E and outflow toward compartment C:

$$EA_M(t+1) = \frac{h_M(t) \times EA_M(t) + Q_{EM}(t) \times EA_E \times \Delta t - Q_{MC}(t) \times EA_M(t) \times \Delta t}{h_M(t) + Q_{EM}(t) \times \Delta t - Q_{MC}(t) \times \Delta t} + EA_M^*(t) \text{ if } Q_{MC}(t) > 0 \quad (10)$$

$$EA_M(t+1) = \frac{h_M(t) \times EA_M(t) + Q_{EM}(t) \times EA_E \times \Delta t - Q_{MC}(t) \times EA_C(t) \times \Delta t}{h_M(t) + Q_{EM}(t) \times \Delta t - Q_{MC}(t) \times \Delta t} + EA_M^*(t) \text{ if } Q_{MC}(t) < 0 \quad (11)$$

where  $Q_{EM}$  and  $Q_{MC}$  [ $\text{mm.day}^{-1}$ ] are internal discharge rates per unit surface area and  $EA_E$ ,  $EA_M$ , and  $EA_C$  are the concentrations of EA [ $\text{cm}^3 \text{STP.kg}^{-1}$ ] in compartments E, M, and C, respectively, and  $\Delta t$  is the unit time step of the model [day].

The EA signature within compartment C depends on the mass balance between inflow from C and outflow toward the spring.

The EA signature at the spring depends on the mass balance between outflow from C and outflow from compartment E

directly to the spring:

$$EA_{sim}(t+1) = \frac{Q_{CS}(t) \times EA_C(t) + Q_{hyES}(t) \times EA_E(t)}{Q_{CS}(t) + Q_{hyES}(t)} \quad (14)$$

where  $Q_{hyES}$  is the discharge from compartment E to the spring.

$$EA_C(t+1) = \frac{h_C(t) \times EA_C(t) + Q_{MC}(t) \times EA_M \times \Delta t + Q_{hyEC}(t) \times EA_E(t) \times \Delta t - Q_{CS}(t) \times EA_C(t) \times \Delta t}{h_C(t) + Q_{MC}(t) \times \Delta t + Q_{hyEC}(t) \times \Delta t - Q_{CS}(t) \times \Delta t} \text{ if } Q_{MC}(t) > 0 \quad (12)$$

$$EA_C(t+1) = \frac{h_C(t) \times EA_C(t) + Q_{MC}(t) \times EA_C \times \Delta t + Q_{hyEC}(t) \times EA_E(t) \times \Delta t - Q_{CS}(t) \times EA_C(t) \times \Delta t}{h_C(t) + Q_{MC}(t) \times \Delta t + Q_{hyEC}(t) \times \Delta t - Q_{CS}(t) \times \Delta t} \text{ if } Q_{MC}(t) < 0 \quad (13)$$

where  $Q_{hyEC}$  is the discharge per unit surface area from compartment E to C [ $\text{mm.day}^{-1}$ ],  $Q_{CS}$  is the discharge per unit surface area from compartment C to the spring [ $\text{mm.day}^{-1}$ ].

## Model calibration

The proposed rainfall discharge model counts 6 parameters ([Table 1](#)) and the EA simulation at the spring brings one more



TABLE 1 Description of the model parameters.

Parameter	Unit	Description
Ra	km <sup>2</sup>	Recharge area
k <sub>EM</sub>	mm.day <sup>-1</sup>	Discharge coefficient for Q <sub>EM</sub>
k <sub>hy</sub>	mm.day <sup>-1</sup>	Discharge coefficient for Q <sub>hy</sub>
E <sub>hy</sub>	mm	Activation level for Q <sub>hy</sub>
X <sub>hy</sub>	-	Partition coefficient for Q <sub>hy</sub> to spring or compartment C
k <sub>MC</sub>	mm.day <sup>-1</sup>	Discharge coefficient for Q <sub>MC</sub>
k <sub>CS</sub>	mm.day <sup>-1</sup>	Discharge coefficient for Q <sub>CS</sub>
a	-	Coefficient for the linear relation $EA_M^* \propto \Delta h_M$ (t)

parameter. The conceptual representation of EA formation is considered as simple as possible to avoid overparameterization and to limit potential parametric equifinality (Perrin et al., 2001). Based on former studies, the recharge area was identified as  $Ra = 116 \text{ km}^2$ . The model is implemented in the python environment (Summerfield, 2010), while the parameter estimation was performed using a particle swarm optimization (PSO) procedure (Eberhart and Kennedy, 1995) implemented in the “pyswarm” package (Lee, 2014). Furthermore, a Latin Hypercube sampling (LHS) is applied to account for parametric uncertainties and investigate the parameter distribution according to the model performance criteria. The calibration is made over 10 years (from September 2009 to December 2019). Because the model is not dedicated to forecasting purposes or extrapolation out of the investigated climatic conditions, no validation period is considered in this study. The calibration is made considering the widely used Nash–Sutcliffe Efficiency (NSE; Nash and Sutcliffe, 1970) criterion on both spring discharge ( $Q_{obs}$  and  $Q_{sim}$ ) and excess air ( $EA_{obs}$  and  $EA_{sim}$ ) measurement at the spring. Then, the objective function  $\Phi$  consists of a weighted sum:

$$\Phi = \omega \times \text{NSE}(Q_{sim}) + (1 - \omega) \times \text{NSE}(EA_{sim}) \quad (15)$$

where  $\omega$  varies from 0 to 1 to account for the different importance of both  $Q_{obs}$  and  $EA_{obs}$  in the calibration procedure. One should note that, within the calibration period,  $Q_{obs}$  consists of a continuous daily time series (3,760 observations), whereas the  $EA_{obs}$  time series counts 18 observations within 3 hydrological cycles.

## Results

### EA measurement analysis

The Durzon spring EA shows important variations during the hydrological cycle. The lowest EA values are lower than  $1 \text{ cm}^3 \text{ STP.kg}^{-1}$  and were observed during floods characterized

by flows between  $3.9$  and  $5.4 \text{ m}^3 \cdot \text{s}^{-1}$ . Recession periods are characterized by high values of excess air up to  $11 \text{ cm}^3 \text{ STP.kg}^{-1}$  when the spring discharge is lower than  $0.8 \text{ m}^3 \cdot \text{s}^{-1}$ . Dissolved gas measurements taken within the shallow epikarst zone during high water level periods indicate that EA is low, around  $0.5 \text{ cm}^3 \text{ STP.kg}^{-1}$  for Canalettes and Reynelles (Pérotin et al., 2021). Low values are also measured for the Bise underground river ( $1.3 \text{ cm}^3 \text{ STP.kg}^{-1}$ ), which makes it possible to more generally qualify the signature of the flows flowing in the drainage network of the vadoze zone. The EA measurement at different depths of the unsaturated zone over the Durzon catchment suggests that the EA formation, along with the water transfer toward the saturated zone, is low (lower than  $2 \text{ cm}^3 \text{ STP.kg}^{-1}$ ) and appears insufficient to explain the EA values measured at the Durzon spring, where the EA values increase up to  $11 \text{ cm}^3 \text{ STP.kg}^{-1}$  during periods of low water. A decreasing exponential-type relationship appears to be observed between the measured EA and the spring discharge (Figure 4). In karst environments, rapid and significant recharge events from the surface to the saturated zone can generate piezometric head variation from meters to tens of meters within the transmission zone. When such recharge occurs, flow reversals are observed between the main drainage features (i.e., caves, conduits, unclogged fractures) and the surrounding porous matrix (Martin and Dean, 2001; Meyerhoff et al., 2012; Mitrofan et al., 2015).

When pressure gradient induces flow from the main drainage network to the surrounding porous matrix, it favors matrix recharge in the vicinity of drainage features (i.e., conduit, unclogged fractures, cave), which induces entrapment of air bubbles in the porous media due to hydrostatic pressure increase, following rapid water fluxes through the main drainage features toward the saturated zone. The greater the piezometric head increase, the greater the hydrostatic pressure applied to the air bubbles, and thus the higher the EA formation. Such a process is assumed to explain the EA formation in the porous matrix during piezometric head increase following recharge events. During recession periods, the spring discharge is mainly supplied with groundwater flow from the matrix, with negligible flow from the conduit networks (Plagnes, 1997). Such behavior explains both the high spring EA values when spring discharge is low and the decreasing exponential relationship between the measured EA and the measured spring discharge (Figure 4). During low water level periods, the matrix mostly contributes to the spring discharge so the EA values at the spring increase due to the arrival of water characterized by higher EA values.

During high flow periods (i.e., high piezometric level in the saturated zone), the spring discharge is mainly supplied by fast flow through the main drainage network from the surface to the saturated zone. During these periods, the measurements performed on water circulation through the infiltration zone seem to indicate that the infiltrated water has a low EA

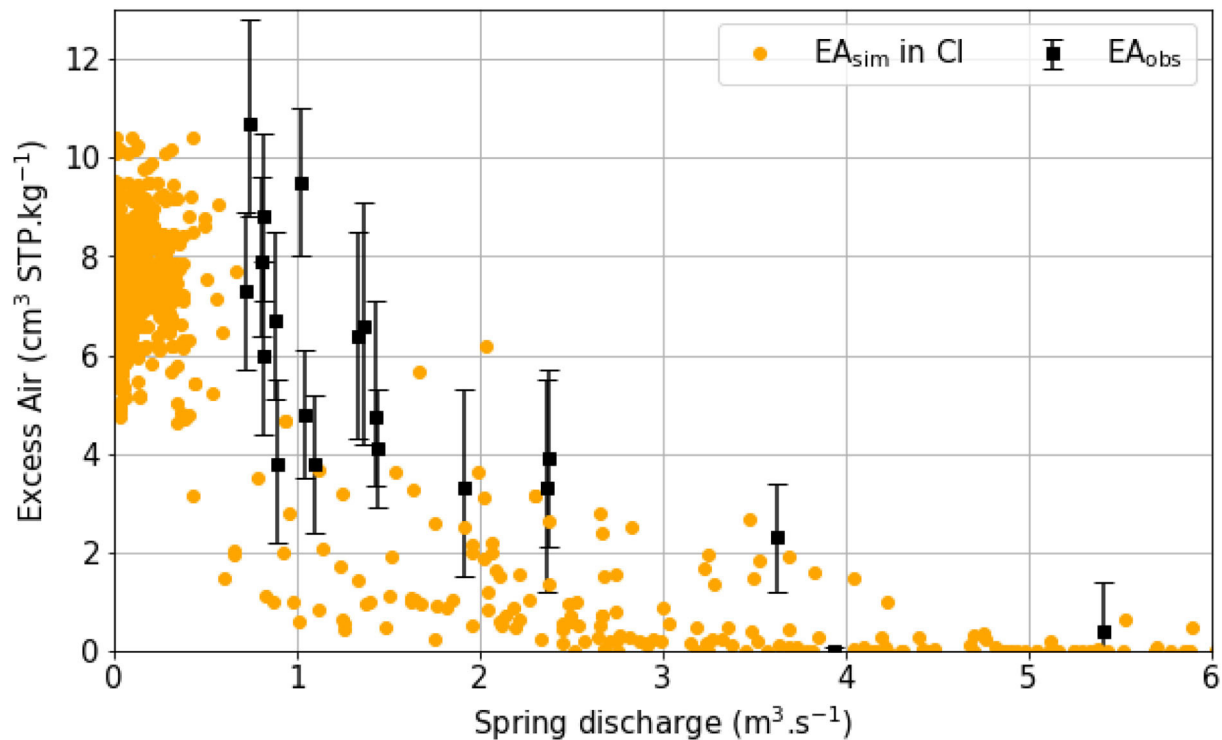


FIGURE 4

Evolution of the EA signature at the Durzon spring as a function of the spring discharge.  $EA_{obs}$  correspond to the EA measurements at the Durzon spring and  $EA_{sim}$  is in the confidence interval (CI) with the 0.1% proportion out of 50,000 simulations within the parameter space sampled with the Latin Hypercube Sampling (LHS) procedure.

signature ( $<1.3 \text{ cm}^3 \text{ STP.kg}^{-1}$ ), and thus EA formation in fast flow groundwater. This would explain the low EA values measured in the Durzon spring during high flow periods assuming no degassing. When the piezometric head decreases at the scale of the aquifer and within the karst drainage network, the lateral exchange with the surrounding porous matrix reverses, causing drainage of groundwater stored in the matrix through the drainage network. This groundwater flow from the matrix contributes to discharge at the spring and thus generates an increase in EA values at the spring. This process of formation and transfer of EA from the subsurface to the spring suggests that the lower EA values measured at the spring will be representative of the EA values of epikarst water.

Additionally, four dissolved gas measurements performed at different depths within the non-saturated zone have shown that the EA varies between  $0.5$  and  $1.2 \text{ cm}^3 \text{ STP.kg}^{-1}$  (Pérotin et al., 2021). This range of variation is one order of magnitude lower than the variation of EA measured at the spring, where measurements indicate variation up to  $11 \text{ cm}^3 \text{ STP.kg}^{-1}$ . For the modeling approach, we will assume that the excess air of the water in the non-saturated zone is equal to  $0.9$ , that is, an intermediate value to those measured in the epikarst and at the Bise.

TABLE 2 Parameter space and values for the parameters estimated with a Particle Swarm Optimization (PSO) procedure implemented in the "pyswarm" package (Lee, 2014) according various objective functions.

	Min	Max	NSE(Q)	$0.5 \cdot \text{NSE}(Q) + 0.5 \cdot \text{NSE}(EA)$	NSE(EA)
$k_{EM}$	0.0001	0.01	0.00029	0.00016	0.00010
$k_{hy}$	0.01	0.5	0.18	0.28	0.24
$E_{hy}$	150	450	231.66	295.85	384.29
$X_{hy}$	0.6	0.9	0.85	0.85	0.66
$k_{MC}$	0.0001	0.01	0.0045	0.0053	0.0097
$k_{CS}$	0.0001	0.01	0.00280	0.00281	0.00275
$a$	0.005	0.5	-	0.10	0.10

## Coupled flow and EA simulation

The model is calibrated considering both spring discharge and EA measurements (Table 2). Figure 5 shows the various time series, which allow for analyzing the internal functioning of the Durzon aquifer. The periods focus on 2010–2012, with a higher EA measurement density although the calibration periods lasted until 2019 (the last EA measurement at the Durzon spring).

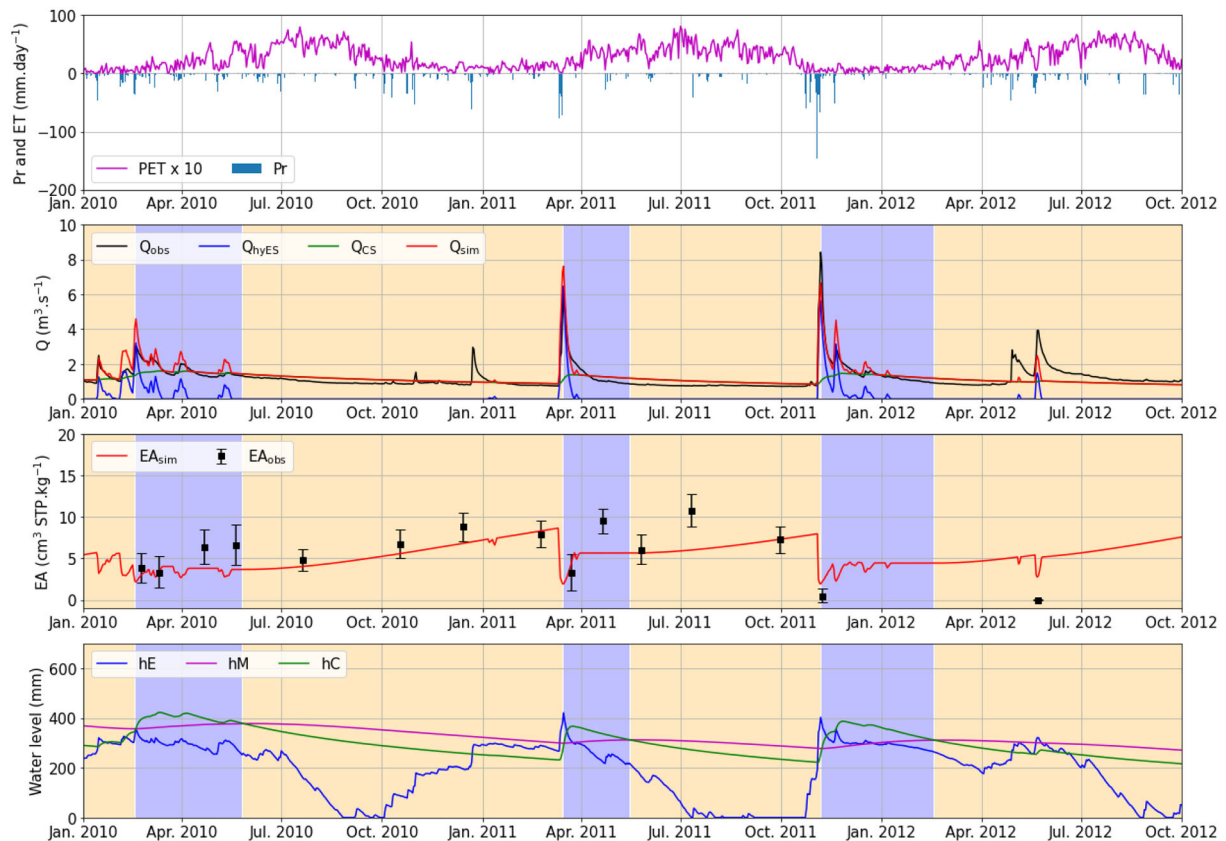


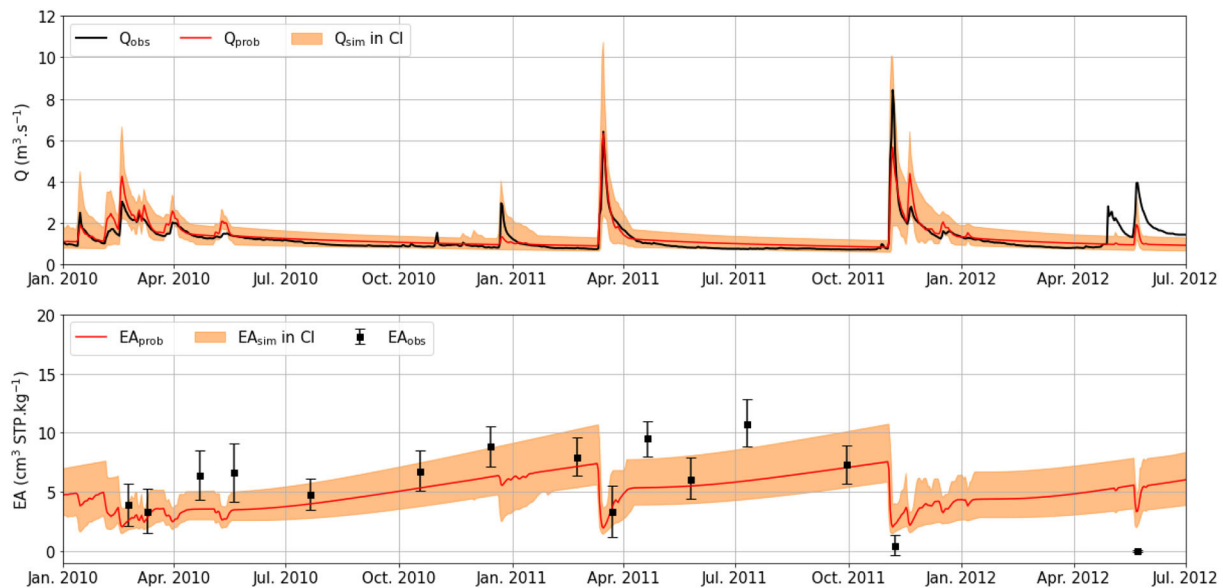
FIGURE 5

(i) Precipitation and potential evapotranspiration time series, (ii) observed and simulated spring discharges, (iii) observed and simulated spring EA time series, (iv) water level time series in the compartments E, M, and C. Periods with blue background color correspond to periods when the water level in C is greater than the water level in M so  $Q_{MC} > 0$  (compartment C feeds the compartment M). Conversely, periods in orange background color correspond to periods when compartment M feed compartment C.

The model calibration shows a quite satisfactory agreement with  $NSE(Q_{sim}) = 0.75$  and  $NSE(EA_{sim}) = 0.53$ . The simulated discharge appears to reproduce the major flood event but lacks the reproduction of some minor floods as in December 2011 and during the May–June period in 2012. The simulated spring discharge is composed of  $Q_{hyES}$  flowing from the compartment directly to the spring, which occurs when the water level in compartment E exceeds the activation threshold  $E_{hy} = 295.9$  mm. This quick flow constitutes the major contribution to the spring discharge during floods, especially following episodes with significant cumulative precipitation. Also, part of the overflow from compartment E contributes to the recharge of compartment C, through the  $Q_{hyEC}$  fluxes and constitutes the major part of the overflow from compartment E. Indeed, the repartition coefficient  $X_{hy}$  is estimated to be 0.84, which means that 84% of the  $Q_{hy}$  contributes to the recharge of compartment C and that 16% flow directly to the spring and thus participate in floods observed after significant precipitation events. During the rest of the year, the fluxes from compartment C to the

spring  $Q_{CS}$  ensure the spring baseflow. The model appears to overestimate the spring discharge during recession periods. The complex flow behavior within the Durzon spring, first explained by considering a hysteretic behavior (Tritz et al., 2011) was better reproduced while considering an infinite characteristic time transfer function (Guinot et al., 2015). However, this model structure does not allow simulating the EA formation as well as exchanges between capacitive and transmissive parts of the aquifer allowing to account for the delay between recharge (EA formation) and increasing spring EA values in periods of low flow (Figure 4).

Despite the sparse temporal resolution of the EA time series, the spring EA variation appears quite well-reproduced. Indeed, the model allows reproducing the observed spring EA decrease following the precipitation episodes in March and November 2011. This occurs simultaneously with the occurrence of  $Q_{hyES}$ , which corresponds to fast flow from the surface to the spring without significant formation of EA, and thus allows reproducing of the EA signature observed at



**FIGURE 6**  
Observed and simulated time series in the confidence interval (CI) with the 0.1% proportion out of 50,000 simulations within the parameter space sampled with the Latin Hypercube Sampling (LHS) procedure: (top) spring discharge  $Q$  and (bottom) spring air excess  $EA$ .  $Q_{prob}$  and  $EA_{prob}$  consists of the weighted mean of the simulation in the confidence interval. The weight corresponds to the objective function  $\Phi$  such as the lower the prediction errors the higher the weight.

the spring because water from compartment E has a low EA concentration, estimated to be around  $= 0.9 \text{ cm}^3 \text{ STP.kg}^{-1}$  (Pérotin et al., 2021). Considering the flow reversal in March 2011, the spring EA increases after a short period and even though the model underestimates the spring EA signature, the dynamics are quite well-reproduced.

The water level in E (hE) shows an important annual variability, with completely dry periods during the late summer periods. Conversely, the water levels in M (hM) and C (hC) present smooth variations, although compartment C presents a much more dynamic response to recharge. The proposed model considers that EA can be formed in compartment M linearly with the recharge intensity. The coefficient is estimated to be equal to 0.1 then the model considers that a recharge intensity of  $1 \text{ mm.day}^{-1}$  will produce an EA of  $0.1 \text{ cm}^3 \text{ STP.kg}^{-1}$ . One should note that the maximum recharge intensity in compartment M is estimated at around  $3 \text{ mm.day}^{-1}$  so the maximum creation of EA during a single day is  $0.3 \text{ cm}^3 \text{ STP.kg}^{-1}$ . In the proposed model, the EA formed during a time step (here, equal to 1 day) is mixed with the water stock available at the considered time step, with its inherited EA signature (i.e., calculated at the previous time step). The numerical model considers an instantaneous perfect mixing, which constitutes a quite simple mixing model but appears reasonable considering the poor temporal resolution of the EA time series measured at the outlet.

## Parametric uncertainties

The LHS procedure is applied to investigate the parameter space with 50,000 simulations where only the 0.1% portion of the best  $\Phi$  value is considered to evaluate the confidence interval (Figure 6). The parametric uncertainties lead to significant predictive errors. Indeed, the confidence interval can cover a wide range of values regarding the spring discharge time series. As an example, considering the flood in November 2011, the confidence interval (CI) covers a range of peak values from  $4.5$  to  $11 \text{ m}^3.\text{s}^{-1}$ , for an observed spring discharge around  $8.5 \text{ m}^3.\text{s}^{-1}$ . Also, some of the simulations lack to reproduce floods of moderate amplitude such as in January 2011. Considering the EA data measured at the outlet, the confidence interval shows consistent dynamics with the observed time series. Although the parametric uncertainties lead to predictive errors, the errors do not change the dynamics of the simulated time series. Also, considering the previously discussed episode in March and November 2011, all the simulations allow the reproduction of the EA decrease observed at the outlet.

Such results show that building the rainfall discharge model considering the concept of EA formation appears consistent with the simulation of spring EA but could introduce significant bias in the simulation of the spring discharge. Also, the parametric uncertainties can be less sensitive considering the EA

time series due to the low temporal resolution of the spring EA time series available for parameter estimation.

The LHS procedure is applied considering three  $\Phi$  so that the calibration considers either the spring discharge alone, both spring discharge and spring EA, or the spring EA alone. Accordingly, the affected weight  $\omega$  to each calibration variable (spring discharge and spring EA) is modified such as  $\omega = [0, 0.5, 1]$ . Figure 7 shows the scatter plots of the values of the objective function (calibration period) against the values of the parameters and the associated kernel density estimation (KDE) for each objective function. The KDE is normalized to 1 for an easier comparison of the distribution according to the different objective functions. According to all the model parameters, different cases can be observed regarding the KDE:

- The KDE shows different distributions if the model is calibrated according to the spring discharge or the spring EA, as observed for the parameters  $k_{hy}$  and  $k_{CS}$ .
- The KDE shows a close position for the distribution peak while the distribution shows a different shape, as observed for the parameters  $k_{EM}$ ,  $E_{hy}$ ,  $X_{hy}$ , and  $k_{MC}$ . Considering  $k_{EM}$ , the KDE appears bimodal when accounting only for spring discharge while the KDE appears unimodal when considering spring EA for model calibration. Also, the KDE appears centered close to  $4 \times 10^{-4} \text{ mm.day}^{-1}$ . Such results also concern  $k_{MC}$  where the distribution spread appears to reduce when considering spring EA for parameter estimation while the optimum value occurs between  $7$  and  $9 \times 10^{-2} \text{ mm.day}^{-1}$ .

Analysis of the parameter distribution shows that the introduction of spring EA measurements may bring uncertainty in the estimation of parameters governing directly the flow dynamics at the spring, that is,  $k_{hy}$  and  $k_{CS}$ . Conversely, considering spring EA seems to reduce the parametric uncertainty for the parameters governing the dynamics within compartment M (representative of the capacitive part of the aquifer), that is,  $k_{EM}$  and  $k_{MC}$ . Such results appear consistent with the working hypothesis because the proposed rainfall discharge model is based on a conceptual representation of the EA formation in karst aquifers.

## Discussion

A parameter estimation procedure accounting only for the EA measurements at the spring shows that the simulated hydrographs still reproduce the main components of the spring discharge flow behavior (Figure 8). Indeed, the major floods appear quite well-reproduced when considering only spring EA for parameter estimation. Nonetheless, some floods such as in January 2010 and January 2011 are not reproduced when considering spring EA only for parameter estimation but

are reproduced when considering only spring discharge for parameter estimation. The proposed rainfall discharge model presents some weaknesses in spring hydrograph simulation compared with the model structure proposed in former studies (Tritz et al., 2011; Mazzilli et al., 2012, 2013; Guinot et al., 2015). However, the proposed model allows reproducing the dynamics of observed spring EA variations and provides consistent spring hydrographs. Such results bring answers to the following questions:

- What is the information content (i.e., data worth) of additional observations such as spring EA?

In karst hydrology, rainfall discharge models are generally calibrated based only on the spring discharge observations. However, several studies have shown an interest in considering additional observations such as  $\text{NO}_3$  and  $\text{SO}_4$  concentrations (Hartmann et al., 2013) or electrical conductivity (Chang et al., 2021), particularly to identify realistic structures of the rainfall discharge model. Within the Durzon karst system, Mazzilli et al. (2013) attempted to calibrate a rainfall discharge reservoir model considering ground-based gravity measurements. The results of the study show that the gravity measurements do not help to better constrain the rainfall discharge model calibration. Among other reasons for such results, the question of the data representativity is risen by the authors. Other studies attempted to include piezometric head variations with conflicting results. As an example, Sivelle and Jourde (2020) identified a complex relationship between piezometric head time series and spring discharge time series within the Oeillal's spring catchment (southern France) and so, did not consider the piezometric head variation to calibrate the rainfall discharge model. Conversely, for the Lez karst system (southern France), considering the piezometric head variation in the estimation of the parameters, besides the discharge, allowed proposing a parsimonious model with better performance than former models (Mazzilli et al., 2011; Cousquer and Jourde, 2022). Also, piezometric head variation derived from ground-based gravity measurement can be of interest for the calibration of a rainfall discharge model but the information content of such data can be disputable on spatial consideration as well as the geometry of the aquifer (Mazzilli et al., 2013).

In our study, we attempted to calibrate the proposed rainfall discharge model under consideration of spring EA measurements, considered as a proxy of the recharge processes at the catchment scale. Nonetheless, the underlying hypothesis for the numerical modeling considers a linear relationship between water level variations in compartment M (representative of the capacitive function of the aquifer) and the formation of EA. Also, the study does not account for observed piezometric head variation because no borehole over the area reaches the saturated zone. Thus, the results lack information to discuss the physical realism in the relationship



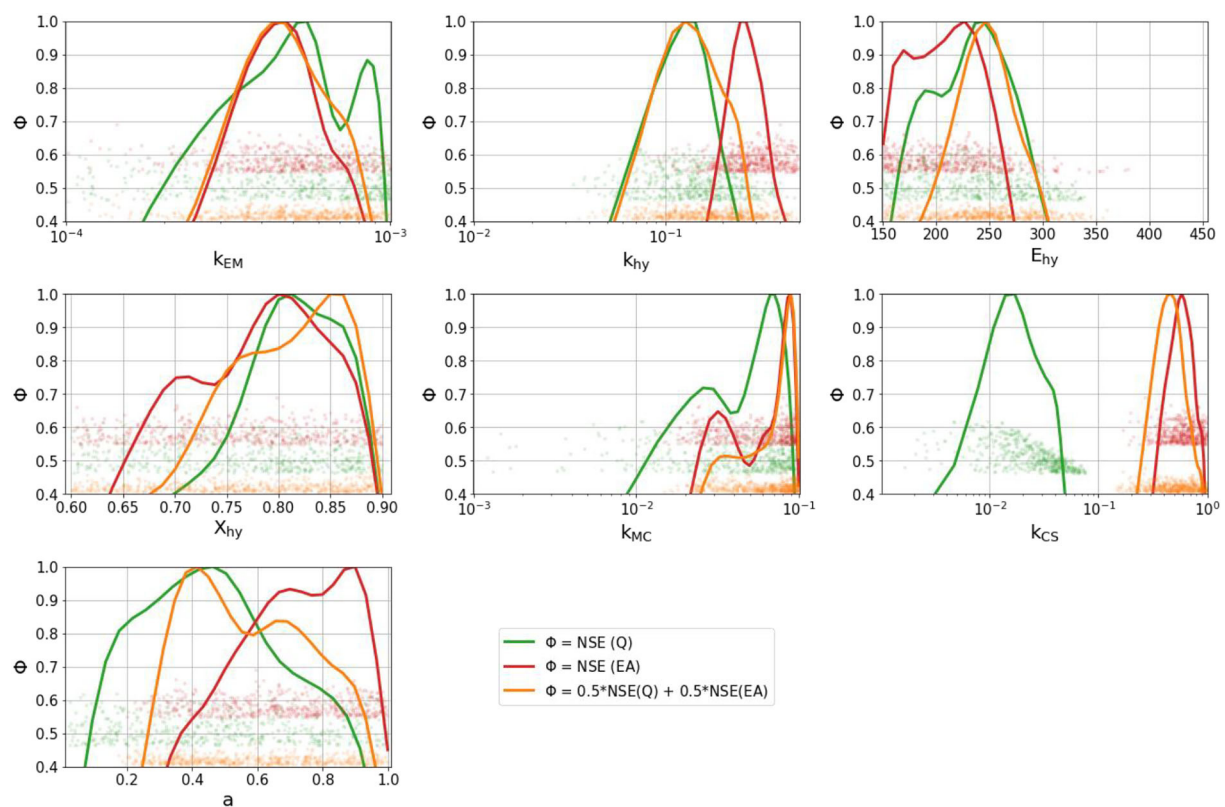


FIGURE 7

Kernel density estimation for each model parameter depending on the objective function  $\Phi$ . Here, only the weight affected by the calibration variables (either spring discharge or spring EA measurement) is modified.

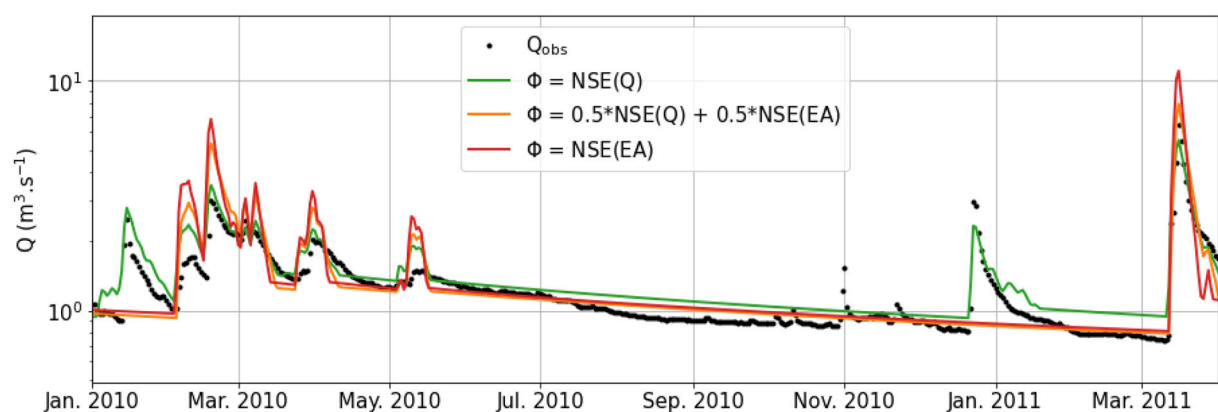


FIGURE 8

Observed and simulated spring discharge time series considering different objective functions.

between EA formation and piezometric head variations. Although some experiments at the laboratory scale (Holocher et al., 2002) bring some relationship between piezometric head variation and EA signature on sand columns, no study has investigated such a relationship in natural hydrogeological

systems, especially in karst aquifers. As suggested by Pérotin et al. (2021) “it is unknown whether there is a direct link between EA and rainfall recharge rates.” The results of the present study show that considering a linear relationship between EA and recharge rates appears efficient to reproduce

the EA variations observed in the spring with a rainfall discharge model.

- How can we choose a suitable model structure?

Several rainfall discharge models have been proposed to simulate the Durzon spring discharge. Deville et al. (2013) proposed a two-level model structure where the upper compartment (C1) is exposed to precipitation and evapotranspiration while the lower level (C2) received flow when water content exceeds the activation threshold in C1. The lower compartment C2 supplies the spring discharge. The model is calibrated according to the ground-based gravity measurements and shows consistent results with the observation. Additional analysis is made, over a longer period and shows that gravity alone cannot constrain a rainfall discharge model because the gravity measurement in the area appears more sensitive to the slow transfer (Fores et al., 2017). Tritz et al. (2011) proposed a hysteretic model to better account for both the suction curve hysteresis in the soil/epikarst zone and the influence of the connectivity of this zone, allowing quick spring discharge increase in periods of important recharge. The authors highlight that using the hysteretic function allows for better reproduction of first flood peaks, which bring complementary information to the ground-based gravity measurements. Nonetheless, Mazzilli et al. (2013) showed that the gravity-derived information might not constrain the parameter estimation of the rainfall discharge model. Guinot et al. (2015) introduced a transfer function with an infinite characteristic time providing satisfactory performance criteria with a lower number of parameters. In this study, the calibration was performed accounting for the spring discharge alone, and no calibration is performed yet accounting for additional observations such as ground-based gravity measurement or hydrochemical observations.

The present study brings a new model structure for the Durzon karst system. The model structure is based on water exchanges between the capacitive and the transmissive compartments of the karst aquifer, considered a prerequisite to explaining the spring EA variation. The introduction of spring EA in the parameter estimation procedure allows both reproducing of the spring discharge dynamics and reducing the parametric uncertainties, especially for the parameter governing the dynamics of the capacitive compartments emptying, namely,  $k_{EM}$  and  $k_{MC}$ . Conversely, the parameters governing the simulated spring discharge, namely,  $k_{hy}$  and  $k_{CS}$  appear not better constrained when considering EA in the calibration, and the model performance for the simulation of spring hydrographs appears lower than the one obtained in former studies (Tritz et al., 2011; Mazzilli et al., 2013; Guinot et al., 2015).

## Conclusion

The study evaluates the relevance of using EA for the characterization of matrix–conduit interactions in karst systems using a rainfall discharge model coupled with the simulation of EA measurements at the outlet of the studied system. The main results show that considering EA measurement can reduce the parametric uncertainties. Indeed, the parameters governing the dynamics of the capacitive compartment (E and M) appear better constrained when, besides spring discharge, EA measurements at the spring are considered as complementary observations in the model calibration. Also, the rainfall discharge model allows for the reproduction of the spring hydrographs even when considering only EA measured at the spring in the objective function. Such results corroborate that the EA variations measured in the spring are informative about the internal dynamics of the Durzon karst aquifer.

Finally, the proposed numerical model allows for validating the conceptual model proposed to explain the EA formation in karst aquifers. Also, the study provides encouraging results on the use of EA as a tracer of flow exchanges. Research perspectives might include the following:

- Specifying the relationship between the EA measured at the spring and the piezometric variations within the saturated zone (which remains unknown for the Durzon system due to the absence of deep drilling),
- Combining the information content of the EA measurements at the spring with other observations such as ground-based gravity measurements, electrical conductivity, and temperature,
- Assessing both structural and parametric uncertainties related to the various types of observations available for rainfall discharge model calibration,
- Assessing the sensitivity of rainfall discharge model predictions to observations as well as assessing the information content of each type of observation for rainfall discharge model calibration,
- Developing more robust rainfall discharge models dedicated to karst aquifers by using hydrochemical observations.

## Data availability statement

The data analyzed in this study is subject to the following licenses/restrictions: The data are provided by the Regional Natural Park of the Grands Causses (PNRGC) and the French National Observation Services H+. Requests to access these datasets should be directed to <http://hplus.ore.fr/>.

## Author contributions

VS: conceptualization, methodology, software, formal analysis, and writing—original draft. LP: conceptualization, data curation, methodology, formal analysis, and writing—original draft. BL: project administration, methodology, supervision, and writing—review and editing. VM: project administration, funding acquisition, and supervision. VB-C: methodology and supervision. CC: methodology, supervision, and writing—review and editing. HJ: conceptualization, methodology, and writing—review and editing. All authors contributed to the article and approved the submitted version.

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## Conflict of interest

The authors declare that the research was conducted in the absence of any commercial or financial relationships that could be construed as a potential conflict of interest.

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