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1	A generic analytical solution for modelling pumping tests in wells
2	intersecting fractures
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10	Abstract
11	The behaviour of transient flow due to pumping in fractured rocks has been studied for at
12	least the past 80 years. Analytical solutions were proposed for solving the issue of a well
13	intersecting and pumping from one vertical, horizontal or inclined fracture in homogeneous
14	aquifers, but their domain of application-even if covering various fracture geometries-was
15	restricted to isotropic or anisotropic aquifers, whose potential boundaries had to be parallel or
16	orthogonal to the fracture direction. The issue thus remains unsolved for many field cases. For
17	example, a well intersecting and pumping a fracture in a multilayer or a dual-porosity aquifer,
18	where intersected fractures are not necessarily parallel or orthogonal to aquifer boundaries,
19	where several fractures with various orientations intersect the well, or the effect of pumping
20	not only in fractures, but also in the aquifer through the screened interval of the well.
21	Using a mathematical demonstration, we show that integrating the well-known Theis
22	analytical solution (Theis, 1935) along the fracture axis is identical to the equally well-known
23	analytical solution of Gringarten et al. (1974) for a uniform-flux fracture fully penetrating a
24	homogeneous aquifer. This result implies that any existing line- or point-source solution can
25	be used for implementing one or more discrete fractures that are intersected by the well.
26	Several theoretical examples are presented and discussed: a single vertical fracture in a dual-
27	porosity aquifer or in a multi-layer system (with a partially intersecting fracture); one and two
28	inclined fractures in a leaky-aquifer system with pumping either only from the fracture(s), or

also from the aquifer between fracture(s) in the screened interval of the well. For the cases

30 with several pumping sources, analytical solutions of flowrate contribution from each 31 individual source (fractures and well) are presented, and the drawdown behaviour according 32 to the length of the pumped screened interval of the well is discussed. Other advantages of 33 this proposed generic analytical solution are also given.

The application of this solution to field data should provide additional field information on fracture geometry, as well as identifying the connectivity between the pumped fractures and other aquifers.

To protect this original concept of a generic solution for modelling pumping tests in fractured
media, a patent application has been deposited on parts of this work (French National Institute
of Industrial Property).

40 <u>Key words:</u> analytical solution, pumping in discrete fractures, fractured rocks, pumping test

41

#### 42 **1. Introduction**

Since the late 1930s (e.g., Strelsova, 1988), much work has been carried out to characterize 43 44 the transient flow of pumping tests carried out in naturally or artificially fractured aquifers (e.g. Muskat, 1937; Warren and Root, 1963; Russell and Truitt, 1964; Bertrand et al., 1980; 45 46 Barker, 1988; Moench, 1984; Hamm and Bidaux, 1996; Jourde et al., 2002; Tiab, 2005; Delay et al., 2007; Rafini and Larocque, 2012; Dewandel et al., 2014; Roques et al., 2016). This led 47 48 to the development of several analytical solutions for understanding the flow behaviour created by a well intersecting and pumping one vertical, horizontal or inclined fracture with 49 infinite or finite hydraulic conductivity embedded in a homogeneous aquifer (Gringarten and 50 Ramey, 1973, 1974; Gringarten et al., 1974; Cinco-Ley et al., 1975, 1998, Thiery, 1980; see 51 also the PetroWiki website). These solutions were obtained by applying the Green's and 52 source functions and the Newman's product method (Newman, 1936; Gringarten and Ramey, 53 1974). However, their domain of application, though proposed for a variety of fracture 54 geometries, is restricted to isotropic or anisotropic infinite aquifers that may be limited in 55 space by no-flow or constant-head boundaries (Gringarten et al., 1974), leaving several 56 possibilities unsolved. For example, pumping a fracture in a multilayer or dual-porosity 57 aquifer, where the fracture is not necessarily parallel or orthogonal to the aquifer boundaries, 58 or a well intersecting and pumping fractures with various orientations, or the effect of 59

pumping from both fractures and also the aquifer in front of the screened interval of the well,etc.

The aim of our research was to seek a new and alternative solution for computing drawdown 62 while pumping in one or several fractures, based on existing line or point source solutions. 63 We first demonstrate that the solutions developed for pumping in a vertical fracture can 64 directly be found by integrating the Theis analytical solution (Theis, 1935) along the fracture 65 axis. Then, by extension, line-source solutions for dual-porosity and multi-layer aquifers and a 66 point-source solution for a leaky aquifer are used when pumping discrete fractures in such 67 aquifers. Theoretical examples are given for pumping in: i) A fracture in a dual-porosity 68 aquifer; ii) A fracture in a multi-layer system that fully or partially intersects one of the 69 aquifer layers; *iii*) An inclined fracture in a homogeneous aquifer; *iv*) Two inclined fractures 70 71 in a leaky-aquifer system with pumping only in the fractures (i.e. the well is only screened in 72 front of the fractures); and v) The same as iv), but now pumping in both the fractures and the 73 aguifer trough the well itself (screened interval). As in iv) and v), the flowrate contributions of each pumped source (i.e., the two fractures and the screened interval of well) to the total 74 pumping rate vary over time, analytical solutions for evaluating their relative contributions are 75 also presented. 76

We do not suggest that the proposed solution should replace existing models used for modelling drawdown in pumping tests performed in homogeneously fractured media (e.g., Barker, 1988; Moench, 1984; Hamm and Bidaux, 1996), or for pumping in fractures in homogeneous aquifers (e.g., Gringarten et al, 1974; Thiéry, 1980). Rather, they are meant to supplement existing models by providing additional hydrogeological information.

82

#### 83

#### 2. Mathematical demonstration

Here, we demonstrate that the well-known Theis analytical solution (1935), defined for a well fully penetrating an isotropic aquifer and integrated along a fracture axis, is strictly identical to the analytical solution of pumping in one vertical fracture proposed by Gringarten et al. (1974), with uniform flux distribution along the fracture plane. In this conceptual aquifer model (Fig. 1), the well intercepts the middle of a vertical fracture of length 2xf and negligible thickness, intersecting a homogeneous and infinite aquifer of transmissivity *T* and storage

coefficient S. Hereafter, we briefly show how we achieved this demonstration; more details 90 can be found in the Supplemental Materials. 91

Assuming that the conductivity of the fracture can be considered infinite, and integrating the Theis 92 well function along the fracture plane leads to the following equation: 93

94 
$$s_{I}(x, y, t) = \frac{1}{4\pi T} \int_{-xf}^{+xf} q(x) E_{1} \left( \frac{(x - x_{obs})^{2} + y_{obs}^{2}}{4Tt} S \right) dx$$
 Eq.1

95

Theis well function

where q(x) is the rate of pumping per unit length of the fracture, t the time since starting the 96 97 pumping, and  $E_1$  the exponential integral (see Fig. 1 for parameters that are not defined in the text). Assuming that the pumping rate Q is uniformly distributed along the fracture, then q(x)98 takes the following form: 99

100 
$$Q = \int_{-x_f}^{+x_f} q(x) dx \Rightarrow q(x) = q = \frac{Q}{2x_f}$$
 Eq.2

Using this statement and the following dimensionless variables:  $t_D = \frac{Tt}{x_f^2 S}$ ;  $x_D = \frac{x_{obs}}{x_f}$ ; 101

102 
$$y_D = \frac{y_{obs}}{x_f}$$
, Eq.1 becomes:

103 
$$s_{I}(x, y, t) = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-x_{f}}^{+x_{f}} \left( \int_{0}^{t_{D}} \frac{e^{-\left[ \left( \frac{x_{f}}{x_{f}} - x_{D} \right)^{2} + y_{D}^{2} \right] / 4\tau}}{\tau} \right) d\tau d\tau d\tau$$
 Eq.3

According to the Fubini theorem ( $\tau$  and x being independent variables), the order of the 104 integration can be inverted, and as  $e^{-y_D^2/4\tau}$  does not depend on x, Eq.3 can be rearranged as: 105

106 
$$s_{I}(x, y, t) = \frac{Q}{2\pi T} \frac{1}{x_{f}} \int_{0}^{t_{D}} e^{-y_{D}^{2}/4\tau} \left( \int_{-x_{f}}^{+x_{f}} \frac{e^{-\left(\frac{x}{x_{f}} - x_{D}\right)^{2}/4\tau}}{4\tau} dx \right) d\tau$$
 Eq.4  
107 (a)

108 When changing variable  $v = \left(\frac{x_{f}}{x_{f}} - x_{D}\right) / 2\sqrt{\tau}$ , the term (a) can be separated into two terms

109 related to the *Erf* function, such as:

110 
$$\mathbf{a} = \int_{-x_f}^{+x_f} \frac{e^{-\left(\frac{x_f}{x_f} - x_D\right)^2 / 4\tau}}{4\tau} dx = \frac{x_f \cdot \sqrt{\pi}}{4\sqrt{\tau}} \left[ -Erf\left(\frac{(1+x_D)}{2\sqrt{\tau}}\right) + Erf\left(\frac{(1-x_D)}{2\sqrt{\tau}}\right) \right]$$
Eq.5

111 Combining Eq.4 and Eq.5, we obtain:

112 
$$s_{I}(x, y, t) = \frac{Q}{2\pi T} \int_{0}^{t_{D}} e^{-\frac{y_{D}^{2}}{4\tau}} \left[ Erf\left(\frac{1-x_{D}}{2\sqrt{\tau}}\right) - Erf\left(\frac{1+x_{D}}{2\sqrt{\tau}}\right) \right] \frac{\sqrt{\pi}}{4\sqrt{\tau}} d\tau$$
Eq.6

Eq.6 demonstrates that the Theis analytical solution, when integrated along the fracture axis, corresponds exactly to the Gringarten et al. (1974) analytical solution (their Eq.20) using Green's and source functions and the Newman's product method for a vertical fracture with an uniform flux distribution that fully penetrates the aquifer.

This implies that the integration of any line-source solution (i.e. a well function defined for a fully or partially penetrating well) along the fracture axis can be used for computing drawdown caused by pumping from a vertical fracture that partially or fully penetrates the aquifer. In addition, any point-source solution can be used for computing drawdown while pumping an inclined or a vertical fracture that is intersected by a well.

122

#### **3.** Theoretical examples

This section presents theoretical examples based on the integration of known analytical line-124 or point-source solutions along the fracture plane. They are based on equation Eq.1 from 125 which the inner integral (Theis well function in Eq.1) is replaced by another well function, 126 assuming that flux is uniformly distributed along the fracture (Eq.2). Dimensionless 127 128 drawdown (s<sub>D</sub>) and its logarithmic derivative (s<sub>D</sub>') vs. dimensionless time with respect to fracture length ( $t_{Dxf}$  or  $t_{DLl}$ , see appendix A) were computed for creating Log-Log diagnostic 129 130 plots and examining transient flow regimes (Bourdet et al., 1983; Deruyck et al., 1992; Renard et al., 2009; Rafini et al., 2017, etc.). As benchmarks, the solutions were compared to 131 132 existing analytical solutions for pumping in a vertical and a horizontal fracture in a homogeneous aquifer (Gringarten et al., 1974; Thiéry, 1980). 133

#### 134 **3.1 Vertical fractures**

These analytical solutions are found by integrating line-source solutions. Because the analytical integration of a given line-source solution along the fracture can be too difficult, we approximate it by first dividing the fracture length into smaller segments and then placing a line-source solution at each segment. Because of the linear properties of the diffusivity equation, we used the principle of superposition, summing up drawdowns induced by each segment to provide the total drawdown value due to pumping in the fracture. Then, this general formulation is used for integrating the solutions numerically:

142 
$$s_{\varpi_{-}f}(x, y, t, x_f, \lambda, \alpha, ...) = \frac{Q}{4\pi T.M} \sum_{i=1,2,3..}^{M} s_{\varpi}(x_i, y, t, \lambda, \alpha, ...)$$
Eq.7

143 where  $s_{\overline{\alpha}}(x, y, t, \lambda, \alpha, ...)$  is a known line-source solution (e.g. dual-porosity, dual-144 permeability or partially penetrating well solutions) and  $s_{\overline{\alpha},f}(x_i, y, t, x_f, \lambda, \alpha, ...)$  is the solution 145 for pumping from a vertical fracture with half-length xf (Fig. 1) in an aquifer whose 146 parameters are defined in the  $s_{\overline{\alpha}}$  solution  $(\lambda, \alpha, ...)$ . The fracture is along the *x*-axis, *M* is the 147 number of segments, and  $x_i$  is the abscissa of the middle of each segment  $(-x_f \le x_i \le +x_f)$ .

# 3.1.1. Pumping at the centre of a vertical fracture fully penetrating a dual-porosity medium

Figure 2 provides an example of the use of Eq.7 with Moench's (1984) dual-porosity model (a 150 model frequently used for interpreting pumping tests performed in homogeneously fractured 151 media). Here, the implemented fracture is vertical and fully penetrates the dual-porosity 152 aquifer (e.g. Fig. 1). Figure 2 shows how pumping in a fracture behaves for a set of 153 154 interporosity flow coefficients  $\lambda$  (see caption for explanation of  $\lambda$ ). As expected, when  $\lambda=0$ (i.e. the matrix hydraulic conductivity is nil corresponding to a single-porosity aquifer), the 155 solution is equivalent to the Gringarten et al. (1974) analytical solution, with an error of 0.2% 156  $(error = |s_G - s_{\overline{\sigma}_{-f}}|/s_G, S_G$  being Gringarten's solution; Eq.6). The difference between the 157 theoretical solution and its numerical evaluation shows that discretization (division of the 158 159 fracture length into smaller elements) errors are very small.

At the start of pumping, as expected, the flow is linear (half-unit slope of derivative curves) and corresponds to flow from the most permeable medium (i.e. the secondary porosity) to the vertical fracture. For intermediate pumping stages, the derivative curves have a classical 'U' 163 shape, characterizing the flow from primary (the block matrix) to secondary porosity of the 164 dual-porosity aquifer. Then, for late pumping stages, the derivatives form a plateau that 165 corresponds to radial flow from the dual-porosity aquifer to the vertical fracture ( $s_D$ '=0.5). In 166 experimental data, this implies that drawdown values on a semi-logarithmic plot form a 167 straight line from which the aquifer transmissivity can be deduced.

3.1.2. Pumping at the centre of a vertical fracture partially penetrating the deepest layer
of a multi-layer aquifer

Our second example (Fig. 3a) corresponds to a multilayer aquifer, where pumping from the 170 deepest aquifer induces depletion in the upper one. This system is characterized by a lower 171 aquifer of thickness B with transmissivity T and storage coefficient S, and an upper aquifer 172 with transmissivity  $T_0$ , and storage coefficient Sy. Both aquifers are separated by an aquitard 173 of hydraulic conductivity k' and thickness B'. The vertical fracture is located in the deeper 174 aquifer and is characterized by its location in the aquifer ( $z \in$  vertical coordinate of the 175 fracture centre), its height  $h_f$  and its length  $2x_f$ . The line-source solution for this conceptual 176 177 model is an extension of the Hunt and Scott (2007) model for a partially penetrating well.

Figure 3b shows type curves for various degrees of penetration of the fracture into the aquifer 178  $(h_{t'}B$  ratio). In these examples, the fracture is located at the centre of the deepest aquifer 179  $(z_f/B=0.5)$ , and there is no hydraulic conductivity anisotropy (kx=ky=kz). As expected when  $h_f$ 180 equals aquifer thickness (B), the solution is identical to the Gringarten et al. (1974) analytical 181 solution until leakage from the upper aquifers starts. At the start of pumping, derivative 182 curves follow the half-unit slope that shows linear flow from the aquifer to the fracture, before 183 decreasing and following a negative slope tending to -1/2 that corresponds to ellipsoidal flow 184 because of the partial entry of the fracture into the aquifer (low  $h_{t}/B$  ratios, curve 5 in Fig. 3b 185 for instance). For intermediate times, the derivative curves form a first plateau corresponding 186 to radial flow into the lower pumped aquifer (s<sub>D</sub>'=0.5). Later, they have a 'V' shape 187 characterizing leakage from the upper aquifers. Finally, for very late stages of pumping, 188 derivative curves form a second plateau whose value depends on the transmissivity values of 189 both upper and lower aquifers  $(s_D' = T/2(T + T_0))$ . 190

- 191 **3.2 Inclined fractures**
- 192
- *3.2.1 Pumping at the centre of a single inclined fracture*

Figure 4a presents the conceptual model. The aquifer is characterized by transmissivity T and 194 a storage coefficient S, and is anisotropic  $(k_x, k_y, k_z)$ . It is overlain by a leaky aquifer of 195 thickness B' and hydraulic conductivity k'. This conceptual aquifer model is similar to that 196 proposed by Hantush (1961) and assumes that the aquitard does not react to pumping (infinite 197 storage). The fracture crosscuts the aquifer, and is characterized by length L, width l, and 198 angle  $\theta$  with the vertical axis. The analytical solution of drawdown for pumping in a fracture 199 200 in such an aquifer is found by integrating the point-source analytical solution given by Hunt (2005) on the fracture plane: 201

$$s(x, y, z, t) = \frac{1}{2\pi T} \int_{-L/2}^{L/2} \int_{-L/2}^{L/2} q(\xi, \theta) \sum_{n=1}^{\infty} \frac{\cos(\alpha_n \xi) \cos(\alpha_n z/B)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r/B)^2}{4t_D}, \alpha_n \frac{r}{B}\sqrt{k_z/k_{xy}}\right) d\xi d\theta \text{ Eq.8}$$

where *r* is the radial distance to the well, *B* is aquifer thickness,  $k_z/k_{xy}$  its vertical anisotropy in hydraulic conductivity,  $t_D$  a dimensionless time,  $q(\xi, \vartheta)$  the pumping rate per unit area,  $\alpha_n$  the root of an equation, and W(a, b) the Hantush leaky-aquifer well function; all parameters are explained in Appendix B.

As before, this equation is solved numerically by dividing the fracture in small elements along both *L* and *l* while assuming a uniformly distributed pumping rate per unit area (similarly to Eq.2; i.e.  $q(\xi, \vartheta) = Q/(lL)$ ). Then the principle of superposition is applied for computing the drawdown at any location into the aquifer. Therefore, the solution yields:

211 
$$s(x, y, z, t) = \frac{Q}{2\pi T.M.p} \sum_{j=1}^{p} \sum_{i=1}^{M} \sum_{n=1}^{\infty} \frac{\cos(\alpha_n z_{i,j}/B) \cos(\alpha_n z/B)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r_{i,j}/B)^2}{4t_D}, \alpha_n \frac{r_{i,j}}{B} \sqrt{k_z/k_{xy}}\right)$$
Eq.9

where  $r_{i,j}$  is the radial distance between the point (x, y, z) and the *i*, *j*<sup>th</sup> element, and *M* and *p* are the number of segments along *L* and *l*.

Figure 4b presents type curves of Eq.9 for various values of  $\theta$ . For this example, the aquifer has a vertical anisotropy of 10.0 ( $k_{xy}/k_z=10$ ) and we ignore the leaky aquifer (k'/B'=0). The fracture is located at the centre of the aquifer ( $z_{f}/B=0.5$ ) and its width equals the aquifer thickness (l=B).

As expected, the results show that Eq.9 is identical to the benchmark solutions: the Gringarten et al. (1974) solution for a vertical rectangular fracture when  $\theta=0^{\circ}$ , and the Thiéry (1980)

solution for a horizontal rectangular fracture located at the centre ( $z_{t'}B=0.5$ ) of the aquifer, 220 when  $\theta$ =90°. Derivative curves show that, regardless of the  $\theta$  value, the early stages of 221 pumping always describe linear flow within the fracture (half-unit slope of derivative curves). 222 Later, derivatives create a hump, more or less pronounced according to the value of  $\theta$ . This 223 behaviour corresponds to the transition from flow perpendicular to the fracture-controlled by 224 225 the average hydraulic conductivity normal to the fracture plane-to horizontal flow controlled by the horizontal aquifer transmissivity. In some cases (e.g.  $\theta=90^{\circ}$ ), the derivatives follow a 226 near-negative half-unit slope characterizing ellipsoidal flow because of the partial penetration 227 228 of the fracture into the aquifer. Note also that increasing the  $k_{xy}/k_z$  ratio with  $\theta \rightarrow 90$ , will result in a more pronounced hump of the derivative curve, because of the increased resistance to 229 flow induced by the low vertical hydraulic conductivity. When the fracture is vertical and 230 231 fully penetrates the aquifer ( $\theta=0^{\circ}$ ), the hump disappears as the drawdown no longer depends upon the vertical anisotropy in hydraulic conductivity. For the late stages of pumping, 232 derivatives form a plateau characterizing the radial flow induced by flow towards the fracture 233 at the aquifer scale ( $s_D$ '=0.5). As mentioned before, in experimental data the plateau value 234 will depend on the horizontal transmissivity of the aquifer. 235

#### *3.2.2 Pumping intersecting two fractures – Pumping in the fractures only*

One of the flexibilities inherent in our analytical/numerical solution is the possibility to consider several fractures with different locations and orientations intersecting the pumping well. Figures 5 and 6 show an example of application where two fractures intersect the well. The characteristics of fracture 1 are *L*: 10 m; *l*: 20 m,  $\theta$ : 20°, and *z<sub>f</sub>*: 13 m, and of fracture 2 *L*: 10 m; *l*: 45 m,  $\theta$ : 100°, and *z<sub>f</sub>*: 5 m. As before, only fractures are pumped (uniform flux).

242 As shown previously, fractures with different orientations and locations are characterized by different flow behaviours. Consequently, their resistance to flow is different; this implies, 243 when several fractures are pumped simultaneously, that the flowrates from each fracture vary 244 over time. Therefore, the flowrate contributions of each fracture to the total pumping rate of 245 the well have to be considered before evaluating the average drawdown in the well. To solve 246 this problem, we again use the principle of superposition, i.e. the drawdown in a particular 247 location is the sum of drawdown values of several pumping sources. This allows computing 248 the flowrates of each fracture from the combination of dimensionless pumping-source 249 solutions (see similar work in Lolon et al. 2008 and Lashgari et al., 2014). However, the 250 boundary conditions of each individual solution do not include perturbation due to the 251

presence of the other pumping source, which is the only approximation of this solution. Note that this approximation is the same as the one used for modelling drawdown (with analytical solutions) created by several pumping wells.

As the well is screened in front of each fracture only and as it intersects the fracture centre, the average dimensionless drawdown in the well,  $s_{Dtot}$ , can be derived from the following equations:

258 
$$s_{Dtot}Q = \frac{1}{2} [q_1(s_{D1} + s_{D12}) + q_2(s_{D2} + s_{D21})]$$
 Eq.10a

259 
$$Q = q_1 + q_2$$
 Eq.10b

where  $q_1$  and  $q_2$  are flowrates from fractures 1 and 2, Q is the total pumping flowrate.  $s_{D1}$  and  $s_{D2}$  are the dimensionless drawdown values at the centre of fractures 1 and 2,  $s_{D12}$  is the dimensionless drawdown induced by fracture 1 at the centre of fracture 2, and  $s_{D21}$ , the one induced by fracture 2 at the centre of fracture 1.

At the intersection between fracture 1 and the well, the drawdown is:

265 
$$S_{F1} = q_1 s_{D1} + q_2 s_{D21}$$
 Eq.10c

and at the intersection between fracture 2 and the well:

267 
$$S_{F2} = q_2 s_{D2} + q_1 s_{D12}$$
 Eq.10d

Assuming equal drawdown at the centre of both fractures ( $S_{F1}=S_{F2}$ ), thus implicitly assuming uniform drawdown at the well location, and solving the system of equations provided by Eqs. 10, one will find solutions for  $s_{Dtot}$ ,  $q_1$  and  $q_2$ :

271 
$$s_{Dtot} = \frac{1}{2} (A + BC) / (1 + C)$$
 Eq.11a

272 
$$q_1 = Q/(1+C)$$
, or  $q_{D1} = 1/(1+C)$  Eq.11b

273 
$$q_2 = QC/(1+C)$$
, or  $q_{D2} = C/(1+C)$  Eq.11c

274 with  $A = (s_{D1} + s_{D12})$ ,  $B = (s_{D2} + s_{D21})$  and  $C = (s_{D1} - s_{D12})/(s_{D2} - s_{D21})$ , and dimensionless 275 flowrate  $q_{Di} = q_i/Q$ . For the computation, drawdown values  $s_{D1}$ ,  $s_{D2}$ ,  $s_{D12}$  and  $s_{D21}$  were calculated separately with Eq.9 while assuming a unit flow rate (Q=1).

Figures 6a and b present the resulting drawdown and the flowrate contributions of each 278 fracture, and figures 6c, d, e and f show drawdown distribution in the vertical plane at 279 different times. In the early stages of pumping, the drawdown derivative is characterized by 280 the classical half-unit slope corresponding to flow from the aquifer to both fractures 281 (t <0.8 min; Fig. 6a and c). In the beginning, both fractures contribute almost equally to the 282 total pumping rate (Fig. 6b), though in the very early stage ( $t \le 0.1 \text{ min}$ ) fracture 2 contributes 283 slightly more because of its larger area. However, very rapidly flowrate from fracture 1 284 becomes dominant as its orientation offers less restriction to flow compared to fracture 2, 285 which is oriented almost orthogonal to the hydraulic conductivity axis of the lower aquifer. 286 287 Between 0.1 and 100 min, the derivative curve shows transitional flow behaviour, followed by a near -1/2 slope (100<t<2000 min) that characterizes the ellipsoidal flow towards both 288 fractures as drawdown starts to progress in the whole aquifer (Fig. 6d). Later 289 (2000<t<7000 min), the derivative stabilizes, forming a plateau characterizing radial flow 290 from the aquifer to both fractures (Fig. 6e) whose value depends on the aquifer transmissivity. 291 For the last stages (t>7000 min), the derivative decreases because of leakage from the top 292 aquitard (Fig. 6f). At the end of pumping, fracture 1 (the smallest; 200 m<sup>2</sup>) contributes 74% of 293 the total flow rate while fracture 2, the largest  $(450 \text{ m}^2)$ , contributes only 26% (Fig. 6b). 294

As a final point, this drawdown curve shows the difficulty of identifying the contribution of more than one fracture on a diagnosis plot, as this looks similar to one of the cases presented in Fig. 4 until leakage appears. In this case, only drawdown observation on other wells in the aquifer and/or few data on the geometry of fractures can provide unequivocal information about these fractures.

# 300 3.2.3. Pumping intersecting two fractures – pumping in the fractures and in the 301 aquifer through the well itself

This last case is similar to the sketch presented on Figure 5, but pumping affects both fractures as well as the aquifer directly through a screened portion of the well. The well intersects the centres of both fractures. As for the previous case, flowrate contributions of each individual pumping source (the two fractures and the well) have to be evaluated before computing the average dimensionless drawdown,  $s_{Dtot}$ , in the well:

307 
$$s_{Dtot}Q = q_1 s_{D1W} + q_2 s_{D2W} + q_W s_W$$
 Eq.12a

308 with 
$$Q = q_1 + q_2 + q_W$$
 Eq.12b

where  $q_1$  and  $q_2$  are flowrates from fractures 1 and 2,  $q_W$  the flowrate from the aquifer through the well and Q the total pumping flowrate.  $s_{D1W}$  and  $s_{D2W}$  are the dimensionless drawdown values induced by fractures 1 and 2 in the screened well interval.  $s_{D1W}$  and  $s_{D2W}$  are the average drawdown values along the screened interval, computed by integrating the analytical solution (Eq.8) between the upper and lower limits of the pumped section.  $s_W$  is the dimensionless drawdown induced by pumping in the screened interval (Hunt's (2005) analytical solution).

At the intersection between fractures and the well, the drawdown in the well has to be identical to that from the fractures. Therefore, at the intersection between the well and fracture 1:

319 
$$q_W s_W = q_1 s_{D1} + q_2 s_{D21}$$
 Eq.12c

320 and between the well and fracture 2:

321 
$$q_W s_W = q_2 s_{D2} + q_1 s_{D12}$$
 Eq.12d

with  $s_{D1}$  and  $s_{D2}$  the dimensionless drawdown values at the centre of fractures 1 and 2,  $s_{D12}$  the dimensionless drawdown induced by fracture 1 at the centre of fracture 2, and  $s_{D21}$  the one induced by fracture 2 at the centre of fracture 1.

Assuming equal drawdown at the centre of both fractures (Eq.12c=Eq.12d) in the well in front of both fractures, and solving the system of equations provided by Eqs.12, solutions are found for  $s_{Dtot}$ ,  $q_1$ ,  $q_2$  and  $q_W$ :

328 
$$s_{Dtot} = (s_{1W} + Cs_{2W} + D)/(1 + C + D/s_W)$$
 Eq.13a

329 
$$q_1 = Q/(1 + C + D/s_W)$$
, or  $q_{D1} = 1/(1 + C + D/s_W)$  Eq.13b

330 
$$q_2 = QC/(1+C+D/s_W)$$
, or  $q_{D2} = C/(1+CD/s_W)$  Eq.13c

331 
$$q_W = QD/[s_W(1+C+D/s_W)]$$
, or  $q_{DW} = D/[s_W(1+C+D/s_W)]$  Eq.13d

where *C* is defined before (Eqs.11),  $D = (s_{D1}s_{D2} - s_{D12}s_{D21})/(s_{D2} - s_{D21})$ , and  $q_{Di} = q_i/Q$  are the dimensionless flowrates.

Drawdown computations and flowrate of each pumping source were made for a well fully penetrating an aquifer with characteristics and fractures as sketched on Figure 5. Drawdown values were computed separately with Eq.9 for the fractures and with the Hunt (2005) analytical solution for the well (vertical, fully penetrating the aquifer and uniform flux distribution). Figures 7a and b show the results and Figures 7c, d, e and f the drawdown distribution (on a vertical axis) at different times.

340 The drawdown behaviour is very different to the case where only the fractures are pumped (Fig. 6a). At the start of pumping (see also Fig. 7c), the half-unit slope persists as the 341 contribution of the aquifer through the well in terms of flowrate is low compared to that from 342 the fractures (Fig. 7b). After 1 to ~1000 minutes, the derivative curve slightly decreases, 343 following a gentle near-negative slope. This period describes the transitional flow regime 344 during which the contribution though the well starts to contribute significantly to the total 345 flowrate. During this period, ellipsoidal flows from fractures still exist (Fig. 7d), but are 346 347 masked by the well contribution. Between 1000 and ~7000 minutes (Fig. 7e), the derivative curve forms a plateau corresponding to near-radial flow induced by the three pumping 348 sources. At the end of pumping, the well contributes up to 52% of the total flow rate, fracture 349 1 up to 36%, while fracture 2 yields only 12%. 350

The high flowrate contribution from the aquifer to the well, here equivalent to both fractures, 351 is due to two reasons: *i*) The medium to high transmissivity of the aquifer (T= $10^{-4}$  m<sup>2</sup>/s), or 352 more exactly its diffusivity; and *ii*) Because the screened interval covers the entire aquifer 353 thickness. If the aquifer transmissivity had been very low (e.g.  $<10^{-8}$  m<sup>2</sup>/s), or if the well had 354 been screened near the fractures, the aquifer contribution through the well would have been 355 negligible or drastically less, creating an average drawdown in the well similar to the one 356 observed on Fig. 6. Appendix C provides an intermediate figure where the well is screened in 357 front of both fractures only. Fractures and aquifer properties are those of Fig. 5 with a well-358 359 screened interval of 7 m. Drawdown and derivative behaviour are different from the previous case, and closer to those of Fig. 6 because of the lower contribution of the aquifer through the 360 361 well to the total pumping flowrate (at the end of pumping, well: 20%, fracture1: 56% and 362 fracture 2: 22%).

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This result shows that where an aquifer has significant transmissivity, the fracture signature (i.e. the half-unit derivative slope in the early stages of pumping) may be very short, and even not seen. This may explain why pumping tests in fractured media commonly are not interpreted with solutions invoking discrete fractures intersected by the well, but with homogeneous aquifer solutions (e.g. dual-porosity models), even if fractures were seen on field data (e.g. Taylor and Howard, 2000; Maréchal et al., 2004; Dewandel et al., 2011, 2017).

369

#### **4.** Conclusions

Integrating the well-known Theis analytical solution (1935) along a fracture axis is 371 mathematically identical to the equally well-known solution of Gringarten et al. (1974) for a 372 uniform-flux fracture fully penetrating an aquifer, obtained with Green's function and the 373 product solution method. Though this result was mathematically expected (Green's function 374 represents the distribution of instantaneous point-sources over length, area or volume), this 375 had never before been demonstrated or highlighted as far as we are aware. This implies that 376 377 any line- or point-source solution integrated along the fracture plane can be used for computing flow through a fracture. In view of the large number of existing analytical 378 379 solutions, this provides a very wide range of applications for this generic analytical solution, and thus helps modelling discrete fractures that are intersected and pumped by a well, in most 380 381 hydrogeological settings.

We give several theoretical examples for dual-porosity or multilayer aquifer types, even when the pumping well intersects several fractures, and for pumping both in fractures and/or directly in the aquifer through the screened interval of a well. Where several sources are pumped (the aquifer trough the well and/or several fractures), solutions of flowrate contributions of each individual pumping source are given and can be extended to any number of fractures.

As expected when the conceptual hydrogeological models respect the theoretical assumptions of benchmark analytical solutions (here for pumping in a vertical or a horizontal fracture), the proposed solutions do not show significant differences between models.

Where a fracture is vertical and fully intersects the aquifer, the classical flow regime of a fracture can be recognized regardless of the conceptual aquifer model used. However, in the case of an inclined fracture, we show that flow behaviour is not unique. In that case, it depends on the screened-interval length of the well, as well as on whether only the fracture, or
both the fracture and the surrounding aquifer are screened. Even, when aquifer transmissivity
(or diffusivity) is not low, the fracture signature (the classical half-unit slope on a derivative
drawdown curve) can be masked by the contribution of the screened interval of the well.

Our work is based on the assumption that flux is uniformly distributed along the fracture and 398 the screened interval of the well. Even if this hypothesis is the closest approximation for an 399 infinite conductivity fracture -it is nonetheless identical at early stages of pumping- solutions 400 for infinite conductivity fracture can be derived from uniform flux analytical solutions 401 (Gringarten et al., 1974). This can be computed by dividing the fracture into small sections, 402 each with uniform flux per unit area, and then by evaluating the flow contribution of each 403 section to arrive at identical drawdown values along the entire fracture, including the screened 404 405 interval of the well if it is pumped as well.

In addition to the possibility of using any line- or point-source solutions, our generic 406 analytical solution offers several other advantages. For example, by using the superposition 407 408 principle (image well theory) it allows developing solutions where no-flow and constant-head boundaries are not necessarily parallel to the fracture directions. Currently available analytical 409 solutions, however, assume fractures parallel or at right angle to boundaries. In the case of a 410 well field, our solution allows computing the overall drawdown created by several pumping 411 wells intersecting fractures with various geometries, or when only some of the wells intersect 412 413 such fractures, the others directly pumping the aquifer. Finally, though we established analytical solutions for rectangular fractures, these solutions can also be developed for any 414 415 other geometrical configuration.

From a practical point view, the proposed analytical solutions and few diagnosis plots as 416 417 presented provide additional information on flow behaviour and drawdown in fractured media. Their application to real field data is expected soon to characterize fracture geometries, 418 particularly in the case of fractured thermo-mineral aquifers (Maréchal et al., 2014). This 419 420 should help defining, among other points, the location of a well within a given deep fracture, and then to establish if this fracture extends very deep or not. In addition, it should help 421 characterizing fractures of the weathered Stratiform Fractured Layer of hard rock aquifers 422 (Lachassagne et al., 2011) as well as their relationships with overlying saprolite (leakage 423 effects). 424

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- 430

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- 525

#### 526 **Figure captions**

- Figure 1. Conceptual sketch of a well intercepting a single vertical fracture; it corresponds to
  the analytical solution of Gringarten et al. (1974) (plan view).
- 529 Figure 2. A single fully penetrating fracture in a dual-porosity aquifer (dual-porosity
- 530 conceptual model: Moench, 1984), pumping at the centre of the fracture. Type curves for
- various  $\lambda$  coefficients ( $\lambda = K_m r^2 / (K_f b_m)$ ). Dimensionless drawdown  $-s_D$ : plain curves and
- derivatives $-s_D$ ': dotted curves. Circles correspond to the Gringarten et al. (1974) solution.
- 533 Dual-porosity model parameters:  $K_f$  and  $K_m$ , fracture and matrix hydraulic conductivity,  $b_m$ :
- block size, *r*: radial distance from the well,  $s_f$  and  $s_m$ : fracture and matrix storage,  $s_{kf}$ = fracture skin.
- 536 Figure 3. A single partially entering vertical fracture in a multilayer aquifer. Fracture in the
- 537 deepest layer, and pumping at the centre of the fracture. a) Conceptual model with aquifer
- parameters (see text for explanation). b) Type curves for various  $h_{f}/B$  ratios. Circles
- 539 correspond to the Gringarten et al. (1974) analytical solution.
- 540 **Figure 4.** Pumping at the centre of an inclined fracture in a leaky aquifer system. a)
- 541 Conceptual model with aquifer parameters (see text for explanation). b) Type curves for
- various angles of the fracture with the vertical,  $\theta$ , and k'/B'=0 (no-leakage). Dots correspond
- to the Gringarten et al. (1974) and Thiéry (1980) analytical solutions.
- 544 **Figure 5.** Two inclined fractures in a leaky aquifer system: model parameters.
- Figure 6. Two inclined fractures in a leaky aquifer system; the well is screened in front of
  both fractures only. a) Average drawdown and its derivative at the well. b) Dimensionless
  flowrate contributions of each fracture. c), d), e) and f) Drawdown distributions at various
  times along the x-axis. Pumping rate is 1 m<sup>3</sup>/h.
- 549 Figure 7. Two inclined fractures in a leaky aquifer system, pumping in the fractures and in
- the aquifer through a fully penetrating well. a) Average drawdown and its derivative at the
- well. b) Dimensionless flowrate contributions of each fracture and of the aquifer through the
- well. c), d), e) and f) Drawdown distributions at various times along the x-axis. Pumping rate
- 553 is  $1 \text{ m}^3/\text{h}$ .
- 554 Caption for the figure in Appendix C: drawdown and its derivative (a) and flowrates (b)
- behaviour for a well screened in front of both fractures (length: 7 m); both fractures and the

- aquifer though the screened-interval are pumped. Fractures and aquifer properties are the ones
- 557 of Figure 5. Pumping rate is  $1 \text{ m}^3/\text{h}$ .
- 558



Figure 1. Conceptual sketch of a well intercepting a single vertical fracture; it corresponds tothe analytical solution of Gringarten et al. (1974) (plan view).

562



563

**Figure 2.** A single fully penetrating fracture in a dual-porosity aquifer (dual-porosity conceptual model: Moench, 1984), pumping at the centre of the fracture. Type curves for various  $\lambda$  coefficients ( $\lambda = K_m r^2/(K_f b_m)$ ). Drawdown  $-s_D$ : plain curves and derivatives $-s_D$ ': dotted curves. Circles correspond to the Gringarten et al. (1974) solution. Dual-porosity model parameters:  $K_f$  and  $K_m$ , fracture and matrix hydraulic conductivity,  $b_m$ : block size, *r*: radial distance from the well,  $s_f$  and  $s_m$ : fracture and matrix storage,  $sk_f$ = fracture skin.





Figure 3. A single partially entering vertical fracture in a multilayer aquifer. Fracture in the
deepest layer, and pumping at the centre of the fracture. a) Conceptual model with aquifer

- parameters (see text for explanation). b) Type curves for various  $h_{f}/B$  ratios. Circles
- 575 correspond to the Gringarten et al. (1974) analytical solution.



577 **Figure 4.** Pumping at the centre of an inclined fracture in a leaky aquifer system. a)

578 Conceptual model with aquifer parameters (see text for explanation). b) Type curves for

various angles of the fracture with the vertical,  $\theta$ , and k'/B'=0 (no-leakage). Dots correspond

to the Gringarten et al. (1974) and Thiéry (1980) analytical solutions.





582



585

**Figure 6.** Two inclined fractures in a leaky aquifer system; the well is screened in front of both fractures only. a) Average drawdown and its derivative at the well. b) Dimensionless flowrate contributions of each fracture. Pumping rate is  $1 \text{ m}^3/\text{h}$ .



**Figure 6. cont'd.** c), d), e) and f) Drawdown distributions at various times along the x-axis.



**Figure 7.** Two inclined fractures in a leaky aquifer system, pumping in the fractures and in the aquifer through a fully penetrating well. a) Average drawdown and its derivative at the well. b) Dimensionless flowrate contributions of each fracture and, of the aquifer though the well. Pumping rate is  $1 \text{ m}^3/\text{h}$ .



**Figure 7. cont'd.** c), d), e) and f) Drawdown distributions at various times along the x-axis.

#### 601 Appendices:

- 602 **Appendix A: Dimensionless drawdown and time** Dimensionless drawdown,  $s_D$ : 603  $s_D = \frac{2\pi T}{Q}s$ 604 With s, the drawdown, T, the transmissivity, and Q, the pumping rate. 605 606 Dimensionless time,  $t_{Dxf}$  and  $t_{DLl}$ : - For a vertical fracture fully penetrating the aquifer,  $t_{Dxf}$ : 607  $t_{Dxf} = \frac{tT}{Sx_{s}^{2}}$ ; with  $x_{f}$ , the half-fracture length, t, the time and S, the storage coefficient of 608 the aquifer. 609 - For a fracture partially penetrating the aquifer,  $t_{DLl}$ : 610 -  $t_{DU} = \frac{tT}{S(L/2, l/2)}$ ; with L and l, the length and the width of the fracture. 611 612
  - 613

#### <u>Appendix B:</u> point-source solution of Hunt (2005)

Point-source (or sink source) solution for the aquifer described in Figure 4 (Hunt, 2005), from
which 3-D anisotropy in hydraulic conductivity has been implemented:

616 
$$s_{\text{sink}}(x_D, y_D, z_D, t_D) = \frac{Q}{2\pi T} \sum_{n=1}^{\infty} \frac{\cos(\alpha_n Z/B)\cos(\alpha_n z_D)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r/B)^2}{4t_D}, \alpha_n \frac{r}{B}\sqrt{k_z/k_{xy}}\right)$$

617 where  $x_D = \frac{x}{B} \sqrt{k_y/k_{xy}}$ ,  $y_D = \frac{y}{B} \sqrt{k_x/k_{xy}}$ ,  $z_D = \frac{z}{B}$ , are dimensionless variables of the point

618 location (x, y, z);  $k_x$ ,  $k_y$ ,  $k_z$  are the hydraulic conductivity along the x, y, z directions; Z is the z 619 coordinate of the point source where pumping Q takes place;  $k_{xy}$  is the horizontal hydraulic 620 conductivity of the aquifer ( $k_{xy} = \sqrt{k_x k_y} = T/B$ ); B is the aquifer thickness and  $t_D = t k_{xy}/SB$ ; 621 *t* is time;  $\alpha_n$  is the root of equation  $\alpha_n \tan(\alpha_n) = \frac{k'/B'}{k_z/B}$ , with *k*' and *B*' the hydraulic 622 conductivity and the thickness of the leaky aquifer; W(a, b) is the Hantush leaky-aquifer well 623 function.

624

625 <u>Appendix C:</u> drawdown and its derivative (a) and flowrates (b) behaviour for a well screened 626 in front of both fractures (length: 7 m); both fractures and the aquifer though the screened-627 interval are pumped. Fractures and aquifer properties are the ones of Figure 5. Pumping rate is 628  $1 \text{ m}^3/\text{h}$ .



629

### 631 Supplemental material:

**Demonstrating that the well Theis funtion integrated along a fracture plane is identical** 

to the analytical solution of pumping in a vertical fracture proposed by Gringarten et al.
 (1974) with uniform flux distribution.

- 635 Definition conceptual model: the borewell intercepts a vertical fracture of length 2xf in a
- homogeneous and infinite aquifer of transmissivity, T, and storativity, S. The width of the
- 637 fracture is assumed to be negligible.





## 639 Notations

640 
$$r^2 = (x - x_{obs})^2 + y_{obs}^2$$
 with x a point in the fracture ;  $t_D = \frac{Tt}{x_f^2 S}$  ;  $x_D = \frac{x_{obs}}{x_f}$  ;  $y_D = \frac{y_{obs}}{x_f}$  ;

641 
$$u = \frac{r^2 \cdot S}{4Tt}$$
 =>  $u = \frac{(x - x_{obs})^2 + y_{obs}^2}{4Tt} S = \frac{(x/x_f - x_D)^2 + y_D^2}{4t_D}$  and  $a = u \cdot t_D = \frac{(x/x_f - x_D)^2 + y_D^2}{4}$ 

642

## 643 **Properties of the well function** W(u):

644 
$$W(u) = E_1(u) = -E_i(-u) = \int_u^\infty \frac{e^{-y}}{y} dy$$

645 With the following change of variable:  $y = \frac{a}{\tau}$ , it becomes that  $\int_{a/t_D}^{\infty} \frac{e^{-y}}{y} dy = \int_{0}^{t_D} \frac{e^{-a/\tau}}{\tau} d\tau$ 

### 646 **Properties of the** *Erf* **function:**

647 
$$Erf(x) = \frac{2}{\sqrt{\pi}} \int_{0}^{x} e^{-u^{2}} du = -\frac{2}{\sqrt{\pi}} \int_{-x}^{0} e^{-u^{2}} du$$

$$648 \qquad Erf(x) = -Erf(-x)$$

#### 650 Well Theis function for a well

The aquifer is pumped in a single vertical well fully penetrating the aquifer with a pumping rate q. The drawdown at a radial distance r from the well can be expressed as:

$$\begin{cases} s(r,t) = \frac{q}{4\pi T} W(u) = \frac{q}{4\pi T} \int_{u}^{\infty} \frac{e^{-y}}{y} dy \\ u = \frac{r^2 S}{4Tt} \end{cases}$$
[Eq.1]

654

653

#### 655 Gringarten analytical solution (Gringarten et al., 1974)

For a pumping in a single vertical fracture intersecting the pumped well, with uniform fluxdistribution, the drawdown can be expressed as:

658 
$$s(x, y, t) = \frac{Q}{2\pi T} s_D \text{ with } s_D = \int_0^{t_D} e^{-\frac{y_D^2}{4t_D}} \left[ Erf\left(\frac{1-x_D}{2\sqrt{\tau}}\right) - Erf\left(\frac{1+x_D}{2\sqrt{\tau}}\right) \right] \frac{\sqrt{\pi}}{4\sqrt{\tau}} d\tau$$
[Eq.2]

659 Where Q is the pumping rate. As Q is assumed to be uniformly distributed along the fracture, 660 therefore:

661 
$$Q = \int_{-x_f}^{+x_f} q(x) dx \implies q(x) = q = \frac{Q}{2x_f}$$

#### 662 Mathematical demonstration

663 Integration of the Theis' solution along the fracture yields:

664 
$$s_{I}(x, y, t) = \int_{-xf}^{+xf} s(x, y, t) dx = \frac{1}{4\pi T} \int_{-xf}^{+xf} q E_{I} \left( \frac{(x - x_{obs})^{2} + y_{obs}^{2}}{4Tt} S \right) dx$$

with the simplified notation and using the property of the  $E_i$  function, it becomes:

$$666 \qquad s_{I} = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} E_{I}\left(\frac{a}{t_{D}}\right) dx = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} \left(\int_{a/t_{D}}^{+\infty} \frac{e^{-u}}{u} du\right) dx = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} \left(\int_{0}^{t_{D}} \frac{e^{-a/\tau}}{\tau} d\tau\right) dx$$

$$\left(\int_{0}^{+\infty} \left(\int_{0}^{+\infty} \frac{e^{-u}}{\tau} d\tau\right) dx\right) dx = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} \left(\int_{0}^{t_{D}} \frac{e^{-a/\tau}}{\tau} d\tau\right) dx$$

667 which can be rewritten as :  $s_I = \frac{Q}{2\pi T} \frac{1}{4x_f} \int_{-x_f}^{+x_f} \left( \int_{0}^{t_D} \frac{e^{-\left[ \left( \frac{x}{x_f} - x_D\right)^2 + y_D^2 \right] / 4\tau}}{\tau} d\tau \right) dx$ 

According to the Fubini theorem (i.e.  $\tau$  and x are independent), the order of integration can be

669 inverted, yielding to:

$$670 \qquad s_{I} = \frac{Q}{2\pi T} \frac{1}{x_{f}} \int_{-x_{f}}^{+x_{f}} \left( \int_{0}^{t_{D}} e^{-y_{D}^{2}/4\tau} \frac{e^{-\left(\frac{x}{x_{f}} - x_{D}\right)^{2}/4\tau}}{4\tau} d\tau \right) dx = \frac{Q}{2\pi T} \frac{1}{x_{f}} \int_{0}^{t_{D}} \left( \int_{-x_{f}}^{+x_{f}} e^{-y_{D}^{2}/4\tau} \frac{e^{-\left(\frac{x}{x_{f}} - x_{D}\right)^{2}/4\tau}}{4\tau} dx \right) d\tau$$

671 and as  $e^{-y_D^2/4\tau}$  does not depend on x, it can be rearranged as:

672 
$$s_{I} = \frac{Q}{2\pi T} \frac{1}{x_{f}} \int_{0}^{t_{D}} e^{-y_{D}^{2}/4\tau} \left( \int_{-x_{f}}^{+x_{f}} \frac{e^{-\left(\frac{x}{x_{f}} - x_{D}\right)^{2}/4\tau}}{4\tau} dx \right) d\tau$$
[Eq.3]

674 With the following change of variable:  $v = \left(\frac{x_{x_f}}{x_f} - x_D\right) / 2\sqrt{\tau}$ , the (a) term of [Eq.3] can be 675 rewritten as :

676 
$$\int_{-x_{f}}^{+x_{f}} \frac{e^{-\left(\frac{x}{x_{f}}-x_{D}\right)^{2}/4\tau}}{4\tau} dx = \frac{x_{f}}{4\sqrt{\tau}} \int_{-(1+x_{D})/2\sqrt{\tau}}^{(1-x_{D})/2\sqrt{\tau}} dv$$
[Eq.4]

The right part of [Eq.4] can be decomposed into two terms related to the *Erf* function:

680 Combining [Eq.3] and [Eq.5], we obtain:

[Eq.6] proves that the Theis analytical solution integrated along the fracture plane is exactly
the same as the Gringarten analytical solution, for a vertical fracture with a uniform flux
distribution (Eq. 2 or the Eq. 20 in Gringarten et al. 1974).