

# **A generic analytical solution for modelling pumping tests in wells**

Benoît Dewandel, Sandra Lanini, Patrick Lachassagne, Jean-Christophe

Maréchal

# **To cite this version:**

Benoît Dewandel, Sandra Lanini, Patrick Lachassagne, Jean-Christophe Maréchal. A generic analytical solution for modelling pumping tests in wells. Journal of Hydrology, 2018, 559, pp.89 - 99. 10.1016/j.jhydrol.2018.02.013. hal-01814503

# **HAL Id: hal-01814503 <https://brgm.hal.science/hal-01814503v1>**

Submitted on 13 Jun 2018

**HAL** is a multi-disciplinary open access archive for the deposit and dissemination of scientific research documents, whether they are published or not. The documents may come from teaching and research institutions in France or abroad, or from public or private research centers.

L'archive ouverte pluridisciplinaire **HAL**, est destinée au dépôt et à la diffusion de documents scientifiques de niveau recherche, publiés ou non, émanant des établissements d'enseignement et de recherche français ou étrangers, des laboratoires publics ou privés.



 inclined fractures in a leaky-aquifer system with pumping either only from the fracture(s), or also from the aquifer between fracture(s) in the screened interval of the well. For the cases

 with several pumping sources, analytical solutions of flowrate contribution from each individual source (fractures and well) are presented, and the drawdown behaviour according to the length of the pumped screened interval of the well is discussed. Other advantages of this proposed generic analytical solution are also given.

 The application of this solution to field data should provide additional field information on fracture geometry, as well as identifying the connectivity between the pumped fractures and other aquifers.

 To protect this original concept of a generic solution for modelling pumping tests in fractured media, a patent application has been deposited on parts of this work (French National Institute of Industrial Property).

Key words: analytical solution, pumping in discrete fractures, fractured rocks, pumping test

#### **1. Introduction**

 Since the late 1930s (e.g., Strelsova, 1988), much work has been carried out to characterize the transient flow of pumping tests carried out in naturally or artificially fractured aquifers (e.g. Muskat, 1937; Warren and Root, 1963; Russell and Truitt, 1964; Bertrand et al., 1980; Barker, 1988; Moench, 1984; Hamm and Bidaux, 1996; Jourde et al., 2002; Tiab, 2005; Delay et al., 2007; Rafini and Larocque, 2012; Dewandel et al., 2014; Roques et al., 2016). This led to the development of several analytical solutions for understanding the flow behaviour created by a well intersecting and pumping one vertical, horizontal or inclined fracture with infinite or finite hydraulic conductivity embedded in a homogeneous aquifer (Gringarten and Ramey, 1973, 1974; Gringarten et al., 1974; Cinco-Ley et al., 1975, 1998, Thiery, 1980; see also the PetroWiki website). These solutions were obtained by applying the Green's and source functions and the Newman's product method (Newman,1936; Gringarten and Ramey, 1974). However, their domain of application, though proposed for a variety of fracture geometries, is restricted to isotropic or anisotropic infinite aquifers that may be limited in space by no-flow or constant-head boundaries (Gringarten et al., 1974), leaving several possibilities unsolved. For example, pumping a fracture in a multilayer or dual-porosity aquifer, where the fracture is not necessarily parallel or orthogonal to the aquifer boundaries, or a well intersecting and pumping fractures with various orientations, or the effect of  pumping from both fractures and also the aquifer in front of the screened interval of the well, etc.

 The aim of our research was to seek a new and alternative solution for computing drawdown while pumping in one or several fractures, based on existing line or point source solutions. We first demonstrate that the solutions developed for pumping in a vertical fracture can directly be found by integrating the Theis analytical solution (Theis, 1935) along the fracture axis. Then, by extension, line-source solutions for dual-porosity and multi-layer aquifers and a point-source solution for a leaky aquifer are used when pumping discrete fractures in such aquifers. Theoretical examples are given for pumping in: *i*) A fracture in a dual-porosity aquifer; *ii*) A fracture in a multi-layer system that fully or partially intersects one of the aquifer layers; *iii*) An inclined fracture in a homogeneous aquifer; *iv*) Two inclined fractures in a leaky-aquifer system with pumping only in the fractures (i.e. the well is only screened in front of the fractures); and *v*) The same as *iv)*, but now pumping in both the fractures and the aquifer trough the well itself (screened interval). As in *iv*) and *v*), the flowrate contributions of each pumped source (i.e., the two fractures and the screened interval of well) to the total pumping rate vary over time, analytical solutions for evaluating their relative contributions are also presented.

 We do not suggest that the proposed solution should replace existing models used for modelling drawdown in pumping tests performed in homogeneously fractured media (e.g., Barker, 1988; Moench, 1984; Hamm and Bidaux, 1996), or for pumping in fractures in homogeneous aquifers (e.g., Gringarten et al, 1974; Thiéry, 1980). Rather, they are meant to supplement existing models by providing additional hydrogeological information.

## **2. Mathematical demonstration**

 Here, we demonstrate that the well-known Theis analytical solution (1935), defined for a well fully penetrating an isotropic aquifer and integrated along a fracture axis, is strictly identical to the analytical solution of pumping in one vertical fracture proposed by Gringarten et al. (1974), with uniform flux distribution along the fracture plane. In this conceptual aquifer model (Fig. 1), the well intercepts the middle of a vertical fracture of length 2*xf* and negligible thickness, intersecting a homogeneous and infinite aquifer of transmissivity *T* and storage 90 coefficient *S*. Hereafter, we briefly show how we achieved this demonstration; more details 91 can be found in the *Supplemental Materials*.

92 Assuming that the conductivity of the fracture can be considered infinite, and integrating the Theis 93 well function along the fracture plane leads to the following equation:

94 
$$
s_t(x, y, t) = \frac{1}{4\pi T} \int_{-xf}^{+xf} q(x) E_1 \left( \frac{(x - x_{obs})^2 + y_{obs}^2}{4Tt} S \right) dx
$$
 Eq.1

95 Theis well function

96 where  $q(x)$  is the rate of pumping per unit length of the fracture, *t* the time since starting the 97 pumping, and  $E_1$  the exponential integral (see Fig. 1 for parameters that are not defined in the 98 text). Assuming that the pumping rate *Q* is uniformly distributed along the fracture, then *q(x)* 99 takes the following form:

100 
$$
Q = \int_{-x_f}^{+x_f} q(x) dx \Rightarrow q(x) = q = \frac{Q}{2x_f}
$$
 Eq.2

Using this statement and the following dimensionless variables:  $x_{f}^{2}S$  $t_p = \frac{Tt}{2}$ *f*  $D = \frac{1}{\sqrt{2 \pi}}$ ; *f obs*  $\frac{D}{x}$ *x* 101  $x_D = \frac{x_{obs}}{x}$ ;

102 
$$
y_D = \frac{y_{obs}}{x_f}
$$
, Eq.1 becomes:

103 
$$
s_1(x, y, t) = \frac{Q}{2\pi T} \frac{1}{4x_f} \int_{-x_f}^{+x_f} \left( \int_0^{t_D} \frac{e^{-\left[ \left( \frac{x}{x_f} - x_D \right)^2 + y_D^2 \right] \right/4\tau}}{\tau} d\tau \right) dx
$$
 Eq.3

104 According to the Fubini theorem ( $\tau$  and  $x$  being independent variables), the order of the 105 integration can be inverted, and as  $e^{-y_D^2/4\tau}$  does not depend on *x*, Eq.3 can be rearranged as:

106 
$$
s_1(x, y, t) = \frac{Q}{2\pi T} \frac{1}{x_f} \int_{0}^{t_p} e^{-y_D^2/4\tau} \left[ \int_{-x_f}^{+x_f} \frac{e^{-\left(\frac{y}{x_f} - x_D\right)^2/4\tau}}{4\tau} d\tau \right] d\tau
$$
 Eq.4

When changing variable  $v = \left(\frac{x}{x_a} - x_b\right)/2\sqrt{\tau}$ Ј  $\left(x\right)_{x}$   $-x_{D}$ L 108 When changing variable  $v = \left(\frac{x}{x_f} - x_p\right) / 2\sqrt{\tau}$ , the term (a) can be separated into two terms

109 related to the *Erf* function, such as:

110 
$$
a = \int_{-x_f}^{+x_f} \frac{e^{-\left(\frac{y}{x_f} - x_D\right)^2/4\tau}}{4\tau} dx = \frac{x_f \cdot \sqrt{\pi}}{4\sqrt{\tau}} \left[ -Erf\left(\frac{(1+x_D)}{2\sqrt{\tau}}\right) + Erf\left(\frac{(1-x_D)}{2\sqrt{\tau}}\right) \right]
$$
 Eq.5

111 Combining Eq.4 and Eq.5, we obtain:

112 
$$
s_1(x, y, t) = \frac{Q}{2\pi T} \int_0^{t_p} e^{-\frac{y_p^2}{4\tau}} \left[ Erf\left(\frac{1-x_p}{2\sqrt{\tau}}\right) - Erf\left(\frac{1+x_p}{2\sqrt{\tau}}\right) \right] \frac{\sqrt{\pi}}{4\sqrt{\tau}} d\tau
$$
 Eq.6

 Eq.6 demonstrates that the Theis analytical solution, when integrated along the fracture axis, corresponds exactly to the Gringarten et al. (1974) analytical solution (their Eq.20) using Green's and source functions and the Newman's product method for a vertical fracture with an uniform flux distribution that fully penetrates the aquifer.

 This implies that the integration of any line-source solution (i.e. a well function defined for a fully or partially penetrating well) along the fracture axis can be used for computing drawdown caused by pumping from a vertical fracture that partially or fully penetrates the aquifer. In addition, any point-source solution can be used for computing drawdown while pumping an inclined or a vertical fracture that is intersected by a well.

122

### 123 **3. Theoretical examples**

 This section presents theoretical examples based on the integration of known analytical line- or point-source solutions along the fracture plane. They are based on equation Eq.1 from which the inner integral (Theis well function in Eq.1) is replaced by another well function, assuming that flux is uniformly distributed along the fracture (Eq.2). Dimensionless 128 drawdown  $(s_D)$  and its logarithmic derivative  $(s_D)$  *vs.* dimensionless time with respect to fracture length (*tDxf* or *tDLl* 129 , see appendix A) were computed for creating Log-Log diagnostic plots and examining transient flow regimes (Bourdet et al., 1983; Deruyck et al., 1992; Renard et al., 2009; Rafini et al., 2017, etc.). As benchmarks, the solutions were compared to existing analytical solutions for pumping in a vertical and a horizontal fracture in a homogeneous aquifer (Gringarten et al., 1974; Thiéry, 1980).

#### 134 **3.1 Vertical fractures**

 These analytical solutions are found by integrating line-source solutions. Because the analytical integration of a given line-source solution along the fracture can be too difficult, we approximate it by first dividing the fracture length into smaller segments and then placing a line-source solution at each segment. Because of the linear properties of the diffusivity equation, we used the principle of superposition, summing up drawdowns induced by each segment to provide the total drawdown value due to pumping in the fracture. Then, this general formulation is used for integrating the solutions numerically:

142 
$$
S_{\varpi_{-f}}(x, y, t, x_f, \lambda, \alpha, \ldots) = \frac{Q}{4\pi T.M} \sum_{i=1,2,3,\ldots}^{M} S_{\varpi}(x_i, y, t, \lambda, \alpha, \ldots)
$$
 Eq.7

143 where  $s_{\varpi}$  (x, y, t,  $\lambda$ ,  $\alpha$ , ...) is a known line-source solution (e.g. dual-porosity, dual-144 permeability or partially penetrating well solutions) and  $s_{\varpi} f(x_i, y, t, x_f, \lambda, \alpha, ...)$  is the solution 145 for pumping from a vertical fracture with half-length *xf* (Fig. 1) in an aquifer whose 146 parameters are defined in the  $s_{\pi}$  solution ( $\lambda$ ,  $\alpha$ ,...). The fracture is along the *x*-axis, *M* is the 147 number of segments, and  $x_i$  is the abscissa of the middle of each segment  $(-x_f \leq x_i \leq +x_f)$ .

# 148 *3.1.1. Pumping at the centre of a vertical fracture fully penetrating a dual-porosity*  149 *medium*

 Figure 2 provides an example of the use of Eq.7 with Moench's (1984) dual-porosity model (a model frequently used for interpreting pumping tests performed in homogeneously fractured media). Here, the implemented fracture is vertical and fully penetrates the dual-porosity aquifer (e.g. Fig. 1). Figure 2 shows how pumping in a fracture behaves for a set of 154 interporosity flow coefficients  $\lambda$  (see caption for explanation of  $\lambda$ ). As expected, when  $\lambda=0$  (i.e. the matrix hydraulic conductivity is nil corresponding to a single-porosity aquifer), the solution is equivalent to the Gringarten et al. (1974) analytical solution, with an error of 0.2% 157 (*error* =  $|s_G - s_{\sigma} f|/s_G$ ,  $S_G$  being Gringarten's solution; Eq.6). The difference between the theoretical solution and its numerical evaluation shows that discretization (division of the fracture length into smaller elements) errors are very small.

160 At the start of pumping, as expected, the flow is linear (half-unit slope of derivative curves) 161 and corresponds to flow from the most permeable medium (i.e. the secondary porosity) to the 162 vertical fracture. For intermediate pumping stages, the derivative curves have a classical 'U'

 shape, characterizing the flow from primary (the block matrix) to secondary porosity of the dual-porosity aquifer. Then, for late pumping stages, the derivatives form a plateau that 165 corresponds to radial flow from the dual-porosity aquifer to the vertical fracture  $(s_D \neq 0.5)$ . In experimental data, this implies that drawdown values on a semi-logarithmic plot form a straight line from which the aquifer transmissivity can be deduced.

 *3.1.2. Pumping at the centre of a vertical fracture partially penetrating the deepest layer of a multi-layer aquifer*

 Our second example (Fig. 3a) corresponds to a multilayer aquifer, where pumping from the deepest aquifer induces depletion in the upper one. This system is characterized by a lower aquifer of thickness *B* with transmissivity *T* and storage coefficient *S*, and an upper aquifer 173 with transmissivity  $T_0$ , and storage coefficient *Sy*. Both aquifers are separated by an aquitard of hydraulic conductivity *k'* and thickness *B'*. The vertical fracture is located in the deeper 175 aquifer and is characterized by its location in the aquifer  $(z_f$ = vertical coordinate of the fracture centre), its height  $h_f$  and its length  $2x_f$ . The line-source solution for this conceptual model is an extension of the Hunt and Scott (2007) model for a partially penetrating well.

 Figure 3b shows type curves for various degrees of penetration of the fracture into the aquifer (*hf/B* ratio). In these examples, the fracture is located at the centre of the deepest aquifer 180 ( $z_f/B=0.5$ ), and there is no hydraulic conductivity anisotropy ( $kx=ky=kz$ ). As expected when  $h_f$  equals aquifer thickness (*B*), the solution is identical to the Gringarten et al. (1974) analytical solution until leakage from the upper aquifers starts. At the start of pumping, derivative curves follow the half-unit slope that shows linear flow from the aquifer to the fracture, before decreasing and following a negative slope tending to -½ that corresponds to ellipsoidal flow 185 because of the partial entry of the fracture into the aquifer (low  $h_f/B$  ratios, curve 5 in Fig. 3b for instance). For intermediate times, the derivative curves form a first plateau corresponding 187 to radial flow into the lower pumped aquifer  $(s_D^{\prime}=0.5)$ . Later, they have a 'V' shape characterizing leakage from the upper aquifers. Finally, for very late stages of pumping, derivative curves form a second plateau whose value depends on the transmissivity values of both upper and lower aquifers  $(s_D' = T/2(T + T_0))$ . 

- **3.2 Inclined fractures**
- 
- *3.2.1 Pumping at the centre of a single inclined fracture*

 Figure 4a presents the conceptual model. The aquifer is characterized by transmissivity *T* and 195 a storage coefficient *S*, and is anisotropic  $(k_x, k_y, k_z)$ . It is overlain by a leaky aquifer of thickness *B'* and hydraulic conductivity *k'*. This conceptual aquifer model is similar to that proposed by Hantush (1961) and assumes that the aquitard does not react to pumping (infinite storage). The fracture crosscuts the aquifer, and is characterized by length *L*, width *l*, and 199 angle  $\theta$  with the vertical axis. The analytical solution of drawdown for pumping in a fracture in such an aquifer is found by integrating the point-source analytical solution given by Hunt (2005) on the fracture plane:

202 
$$
s(x, y, z, t) = \frac{1}{2\pi T} \int_{-L/2 - L/2}^{L/2 - L/2} q(\xi, \theta) \sum_{n=1}^{\infty} \frac{\cos(\alpha_n \xi) \cos(\alpha_n z/B)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r/B)^2}{4t_D}, \alpha_n \frac{r}{B} \sqrt{k_z/k_{xy}}\right) d\xi d\theta \text{ Eq.8}
$$

203 where *r* is the radial distance to the well, *B* is aquifer thickness,  $k_z/k_x$  its vertical anisotropy in 204 hydraulic conductivity,  $t_D$  a dimensionless time,  $q(\xi, \theta)$  the pumping rate per unit area,  $\alpha_n$  the 205 root of an equation, and *W(a, b)* the Hantush leaky-aquifer well function; all parameters are 206 explained in Appendix B.

 As before, this equation is solved numerically by dividing the fracture in small elements along both *L* and *l* while assuming a uniformly distributed pumping rate per unit area (similarly to 209 Eq.2; i.e.  $q(\xi, \theta) = Q/(lL)$ ). Then the principle of superposition is applied for computing the drawdown at any location into the aquifer. Therefore, the solution yields:

211 
$$
s(x, y, z, t) = \frac{Q}{2\pi T.M.p} \sum_{j=1}^{p} \sum_{i=1}^{M} \sum_{n=1}^{\infty} \frac{\cos(\alpha_n z_{i,j}/B) \cos(\alpha_n z/B)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r_{i,j}/B)^2}{4t_D}, \alpha_n \frac{r_{i,j}}{B} \sqrt{k_z/k_{xy}}\right)
$$
 Eq.9

212 where  $r_{i,j}$  is the radial distance between the point  $(x, y, z)$  and the *i*, *j*<sup>th</sup> element, and *M* and *p* 213 are the number of segments along *L* and *l*.

214 Figure 4b presents type curves of Eq.9 for various values of  $\theta$ . For this example, the aquifer 215 has a vertical anisotropy of 10.0  $(k_x/k_z=10)$  and we ignore the leaky aquifer  $(k'/B'=0)$ . The 216 fracture is located at the centre of the aquifer  $(z/\sqrt{B}=0.5)$  and its width equals the aquifer 217 thickness  $(l=B)$ .

218 As expected, the results show that Eq.9 is identical to the benchmark solutions: the Gringarten 219 et al. (1974) solution for a vertical rectangular fracture when  $\theta = 0^\circ$ , and the Thiéry (1980) 220 solution for a horizontal rectangular fracture located at the centre  $(z \beta = 0.5)$  of the aquifer, 221 when  $\theta = 90^\circ$ . Derivative curves show that, regardless of the  $\theta$  value, the early stages of pumping always describe linear flow within the fracture (half-unit slope of derivative curves). 223 Later, derivatives create a hump, more or less pronounced according to the value of  $\theta$ . This behaviour corresponds to the transition from flow perpendicular to the fracture—controlled by the average hydraulic conductivity normal to the fracture plane—to horizontal flow controlled 226 by the horizontal aquifer transmissivity. In some cases (e.g.  $\theta = 90^{\circ}$ ), the derivatives follow a near-negative half-unit slope characterizing ellipsoidal flow because of the partial penetration 228 of the fracture into the aquifer. Note also that increasing the  $k_{xy}/k_z$  ratio with  $\theta \rightarrow 90$ , will result in a more pronounced hump of the derivative curve, because of the increased resistance to flow induced by the low vertical hydraulic conductivity. When the fracture is vertical and 231 fully penetrates the aquifer  $(\theta=0^{\circ})$ , the hump disappears as the drawdown no longer depends upon the vertical anisotropy in hydraulic conductivity. For the late stages of pumping, derivatives form a plateau characterizing the radial flow induced by flow towards the fracture 234 at the aquifer scale  $(s_D^{\text{th}}=0.5)$ . As mentioned before, in experimental data the plateau value will depend on the horizontal transmissivity of the aquifer.

### *3.2.2 Pumping intersecting two fractures – Pumping in the fractures only*

 One of the flexibilities inherent in our analytical/numerical solution is the possibility to consider several fractures with different locations and orientations intersecting the pumping well. Figures 5 and 6 show an example of application where two fractures intersect the well. 240 The characteristics of fracture 1 are *L*: 10 m; *l*: 20 m,  $\theta$ : 20°, and *zf*. 13 m, and of fracture 2 *L*: 241 10 m; *l*: 45 m,  $\theta$ : 100°, and  $z_f$ : 5 m. As before, only fractures are pumped (uniform flux).

 As shown previously, fractures with different orientations and locations are characterized by different flow behaviours. Consequently, their resistance to flow is different; this implies, when several fractures are pumped simultaneously, that the flowrates from each fracture vary over time. Therefore, the flowrate contributions of each fracture to the total pumping rate of the well have to be considered before evaluating the average drawdown in the well. To solve this problem, we again use the principle of superposition, i.e. the drawdown in a particular location is the sum of drawdown values of several pumping sources. This allows computing the flowrates of each fracture from the combination of dimensionless pumping-source solutions (see similar work in Lolon et al. 2008 and Lashgari et al., 2014). However, the boundary conditions of each individual solution do not include perturbation due to the 252 presence of the other pumping source, which is the only approximation of this solution. Note 253 that this approximation is the same as the one used for modelling drawdown (with analytical 254 solutions) created by several pumping wells.

255 As the well is screened in front of each fracture only and as it intersects the fracture centre, 256 the average dimensionless drawdown in the well,  $s_{Dot}$  can be derived from the following 257 equations:

258 
$$
s_{Dtot}Q = \frac{1}{2} [q_1 (s_{D1} + s_{D12}) + q_2 (s_{D2} + s_{D21})]
$$
 Eq. 10a

$$
259 \tQ = q_1 + q_2 \tEq. 10b
$$

260 where  $q_1$  and  $q_2$  are flowrates from fractures 1 and 2,  $Q$  is the total pumping flowrate.  $s_{D1}$  and 261 *sD2* are the dimensionless drawdown values at the centre of fractures 1 and 2, *sD12* is the 262 dimensionless drawdown induced by fracture 1 at the centre of fracture 2, and  $s_{D21}$ , the one 263 induced by fracture 2 at the centre of fracture 1.

264 At the intersection between fracture 1 and the well, the drawdown is:

265 
$$
S_{F1} = q_1 s_{D1} + q_2 s_{D21}
$$
 Eq. 10c

266 and at the intersection between fracture 2 and the well:

$$
S_{F2} = q_2 s_{D2} + q_1 s_{D12}
$$
 Eq. 10d

268 Assuming equal drawdown at the centre of both fractures  $(S_{F1}=S_{F2})$ , thus implicitly assuming 269 uniform drawdown at the well location, and solving the system of equations provided by Eqs. 270 – 10, one will find solutions for  $s_{Dtot}$ ,  $q_1$  and  $q_2$ :

271 
$$
s_{Dot} = \frac{1}{2}(A + BC)/(1 + C)
$$
 Eq. 11a

272 
$$
q_1 = Q/(1+C)
$$
, or  $q_{D1} = 1/(1+C)$  Eq.11b

273 
$$
q_2 = QC/(1+C)
$$
, or  $q_{D2} = C/(1+C)$  Eq.11c

with  $A = (s_{D1} + s_{D12})$ ,  $B = (s_{D2} + s_{D21})$  and  $C = (s_{D1} - s_{D12})/(s_{D2} - s_{D21})$ , and dimensionless 274 275 flowrate  $q_{Di} = q_i/Q$ .

276 For the computation, drawdown values  $s_{D1}$ ,  $s_{D2}$ ,  $s_{D12}$  and  $s_{D21}$  were calculated separately with 277 Eq.9 while assuming a unit flow rate  $(Q=1)$ .

 Figures 6a and b present the resulting drawdown and the flowrate contributions of each fracture, and figures 6c, d, e and f show drawdown distribution in the vertical plane at different times. In the early stages of pumping, the drawdown derivative is characterized by the classical half-unit slope corresponding to flow from the aquifer to both fractures (t <0.8 min; Fig. 6a and c). In the beginning, both fractures contribute almost equally to the 283 total pumping rate (Fig. 6b), though in the very early stage ( $t \le 0.1$  min) fracture 2 contributes slightly more because of its larger area. However, very rapidly flowrate from fracture 1 becomes dominant as its orientation offers less restriction to flow compared to fracture 2, which is oriented almost orthogonal to the hydraulic conductivity axis of the lower aquifer. Between 0.1 and 100 min, the derivative curve shows transitional flow behaviour, followed 288 by a near  $-1/2$  slope (100  $\lt t \lt 2000$  min) that characterizes the ellipsoidal flow towards both fractures as drawdown starts to progress in the whole aquifer (Fig. 6d). Later (2000<t<7000 min), the derivative stabilizes, forming a plateau characterizing radial flow from the aquifer to both fractures (Fig. 6e) whose value depends on the aquifer transmissivity. 292 For the last stages ( $t$ >7000 min), the derivative decreases because of leakage from the top 293 aquitard (Fig. 6f). At the end of pumping, fracture 1 (the smallest;  $200 \text{ m}^2$ ) contributes 74% of 294 the total flow rate while fracture 2, the largest  $(450 \text{ m}^2)$ , contributes only 26% (Fig. 6b).

 As a final point, this drawdown curve shows the difficulty of identifying the contribution of more than one fracture on a diagnosis plot, as this looks similar to one of the cases presented in Fig. 4 until leakage appears. In this case, only drawdown observation on other wells in the aquifer and/or few data on the geometry of fractures can provide unequivocal information about these fractures.

# *3.2.3. Pumping intersecting two fractures – pumping in the fractures and in the aquifer through the well itself*

 This last case is similar to the sketch presented on Figure 5, but pumping affects both fractures as well as the aquifer directly through a screened portion of the well. The well intersects the centres of both fractures. As for the previous case, flowrate contributions of each individual pumping source (the two fractures and the well) have to be evaluated before computing the average dimensionless drawdown,  $s_{Dtot}$ , in the well:

$$
307 \t SDolQ = q1sD1W + q2sD2W + qWsW
$$
 Eq. 12a

308 with 
$$
Q = q_1 + q_2 + q_w
$$
 Eq. 12b

309 where  $q_1$  and  $q_2$  are flowrates from fractures 1 and 2,  $q_W$  the flowrate from the aquifer through 310 the well and *Q* the total pumping flowrate.  $s_{D1W}$  and  $s_{D2W}$  are the dimensionless drawdown 311 values induced by fractures 1 and 2 in the screened well interval.  $s_{D1W}$  and  $s_{D2W}$  are the 312 average drawdown values along the screened interval, computed by integrating the analytical 313 solution (Eq.8) between the upper and lower limits of the pumped section.  $s_W$  is the 314 dimensionless drawdown induced by pumping in the screened interval (Hunt's (2005) 315 analytical solution).

316 At the intersection between fractures and the well, the drawdown in the well has to be 317 identical to that from the fractures. Therefore, at the intersection between the well and fracture 318 1:

319 
$$
q_w s_w = q_1 s_{D1} + q_2 s_{D21}
$$
 Eq. 12c

320 and between the well and fracture 2:

$$
321 \tqWsW = q2sD2 + q1sD12
$$
 Eq. 12d

322 with *sD1* and *sD2* the dimensionless drawdown values at the centre of fractures 1 and 2, *sD12* the 323 dimensionless drawdown induced by fracture 1 at the centre of fracture 2, and  $s_{D21}$  the one 324 induced by fracture 2 at the centre of fracture 1.

325 Assuming equal drawdown at the centre of both fractures (Eq.12c=Eq.12d) in the well in front 326 of both fractures, and solving the system of equations provided by Eqs.12, solutions are found 327 for  $s_{Dtot}$ ,  $q_1$ ,  $q_2$  and  $q_W$ :

328 
$$
s_{Dtot} = (s_{1W} + Cs_{2W} + D)/(1 + C + D/s_W)
$$
 Eq. 13a

329 
$$
q_1 = Q/(1 + C + D/s_w)
$$
, or  $q_{D1} = 1/(1 + C + D/s_w)$  Eq. 13b

330 
$$
q_2 = QC/(1 + C + D/s_w)
$$
, or  $q_{D2} = C/(1 + CD/s_w)$  Eq.13c

331 
$$
q_w = QD/[s_w(1 + C + D/s_w)], \text{ or } q_{DW} = D/[s_w(1 + C + D/s_w)]
$$
 Eq. 13d

where C is defined before (Eqs.11),  $D = (s_{D1}s_{D2} - s_{D12}s_{D21})/(s_{D2} - s_{D21})$ , and  $q_{Di} = q_i/Q$  are the dimensionless flowrates.

 Drawdown computations and flowrate of each pumping source were made for a well fully penetrating an aquifer with characteristics and fractures as sketched on Figure 5. Drawdown values were computed separately with Eq.9 for the fractures and with the Hunt (2005) analytical solution for the well (vertical, fully penetrating the aquifer and uniform flux distribution). Figures 7a and b show the results and Figures 7c, d, e and f the drawdown distribution (on a vertical axis) at different times.

 The drawdown behaviour is very different to the case where only the fractures are pumped (Fig. 6a). At the start of pumping (see also Fig. 7c), the half-unit slope persists as the contribution of the aquifer through the well in terms of flowrate is low compared to that from the fractures (Fig. 7b). After 1 to ~1000 minutes, the derivative curve slightly decreases, following a gentle near-negative slope. This period describes the transitional flow regime during which the contribution though the well starts to contribute significantly to the total flowrate. During this period, ellipsoidal flows from fractures still exist (Fig. 7d), but are masked by the well contribution. Between 1000 and ~7000 minutes (Fig. 7e), the derivative curve forms a plateau corresponding to near-radial flow induced by the three pumping sources. At the end of pumping, the well contributes up to 52% of the total flow rate, fracture 1 up to 36%, while fracture 2 yields only 12%.

 The high flowrate contribution from the aquifer to the well, here equivalent to both fractures, 352 is due to two reasons: *i*) The medium to high transmissivity of the aquifer (T=10<sup>-4</sup> m<sup>2</sup>/s), or more exactly its diffusivity; and *ii*) Because the screened interval covers the entire aquifer 354 thickness. If the aquifer transmissivity had been very low (e.g.  $\langle 10^{-8} \text{ m}^2/\text{s}$ ), or if the well had been screened near the fractures, the aquifer contribution through the well would have been negligible or drastically less, creating an average drawdown in the well similar to the one observed on Fig. 6. Appendix C provides an intermediate figure where the well is screened in front of both fractures only. Fractures and aquifer properties are those of Fig. 5 with a well- screened interval of 7 m. Drawdown and derivative behaviour are different from the previous case, and closer to those of Fig. 6 because of the lower contribution of the aquifer through the well to the total pumping flowrate (at the end of pumping, well: 20%, fracture1: 56% and fracture 2: 22%).

Journal of Hydrology 559 (2018) 89–99

 This result shows that where an aquifer has significant transmissivity, the fracture signature (i.e. the half-unit derivative slope in the early stages of pumping) may be very short, and even not seen. This may explain why pumping tests in fractured media commonly are not interpreted with solutions invoking discrete fractures intersected by the well, but with homogeneous aquifer solutions (e.g. dual-porosity models), even if fractures were seen on field data (e.g. Taylor and Howard, 2000; Maréchal et al., 2004; Dewandel et al., 2011, 2017).

#### **4. Conclusions**

 Integrating the well-known Theis analytical solution (1935) along a fracture axis is mathematically identical to the equally well-known solution of Gringarten et al. (1974) for a uniform-flux fracture fully penetrating an aquifer, obtained with Green's function and the product solution method. Though this result was mathematically expected (Green's function represents the distribution of instantaneous point-sources over length, area or volume), this had never before been demonstrated or highlighted as far as we are aware. This implies that any line- or point-source solution integrated along the fracture plane can be used for computing flow through a fracture. In view of the large number of existing analytical solutions, this provides a very wide range of applications for this generic analytical solution, and thus helps modelling discrete fractures that are intersected and pumped by a well, in most hydrogeological settings.

 We give several theoretical examples for dual-porosity or multilayer aquifer types, even when the pumping well intersects several fractures, and for pumping both in fractures and/or directly in the aquifer through the screened interval of a well. Where several sources are pumped (the aquifer trough the well and/or several fractures), solutions of flowrate contributions of each individual pumping source are given and can be extended to any number of fractures.

 As expected when the conceptual hydrogeological models respect the theoretical assumptions of benchmark analytical solutions (here for pumping in a vertical or a horizontal fracture), the proposed solutions do not show significant differences between models.

 Where a fracture is vertical and fully intersects the aquifer, the classical flow regime of a fracture can be recognized regardless of the conceptual aquifer model used. However, in the case of an inclined fracture, we show that flow behaviour is not unique. In that case, it  depends on the screened-interval length of the well, as well as on whether only the fracture, or both the fracture and the surrounding aquifer are screened. Even, when aquifer transmissivity (or diffusivity) is not low, the fracture signature (the classical half-unit slope on a derivative drawdown curve) can be masked by the contribution of the screened interval of the well.

 Our work is based on the assumption that flux is uniformly distributed along the fracture and the screened interval of the well. Even if this hypothesis is the closest approximation for an infinite conductivity fracture –it is nonetheless identical at early stages of pumping- solutions for infinite conductivity fracture can be derived from uniform flux analytical solutions (Gringarten et al., 1974). This can be computed by dividing the fracture into small sections, each with uniform flux per unit area, and then by evaluating the flow contribution of each section to arrive at identical drawdown values along the entire fracture, including the screened interval of the well if it is pumped as well.

 In addition to the possibility of using any line- or point-source solutions, our generic analytical solution offers several other advantages. For example, by using the superposition principle (image well theory) it allows developing solutions where no-flow and constant-head boundaries are not necessarily parallel to the fracture directions. Currently available analytical solutions, however, assume fractures parallel or at right angle to boundaries. In the case of a well field, our solution allows computing the overall drawdown created by several pumping wells intersecting fractures with various geometries, or when only some of the wells intersect such fractures, the others directly pumping the aquifer. Finally, though we established analytical solutions for rectangular fractures, these solutions can also be developed for any other geometrical configuration.

 From a practical point view, the proposed analytical solutions and few diagnosis plots as presented provide additional information on flow behaviour and drawdown in fractured media. Their application to real field data is expected soon to characterize fracture geometries, particularly in the case of fractured thermo-mineral aquifers (Maréchal et al., 2014). This should help defining, among other points, the location of a well within a given deep fracture, and then to establish if this fracture extends very deep or not. In addition, it should help characterizing fractures of the weathered Stratiform Fractured Layer of hard rock aquifers (Lachassagne et al., 2011) as well as their relationships with overlying saprolite (leakage effects).

#### **Acknowledgements**

This study was partly conducted under a research agreement between Evian-Volvic World

- (Water Institute by Evian, Danone Waters France) and BRGM. We are grateful to Dr. H.M. Kluijver for revising the English text.
- 

#### **References**

- Barker, J.A., 1988. A generalized radial flow model for hydraulic tests in fractured rock, Water Res. Res., 24, 1796-1804.
- Bertrand, L, Feuga, B., Noyer, M.L., Thiéry, D., 1980. Hot Dry Rocks, contribution to the methodology of determining the hydraulic properties of naturally or artificially fractured rocks. In French, BRGM report 80SGN029GEG, 229 p.
- Bourdet, D., Whittle, T.M., Dougals, A.A., Pirard, V.M., 1983. A new set of type curves simplifies well test analysis, World Oil., May, 95-106.
- Cinco-Ley, H., Ramey, H.J., Miller, F.G., 1975. Unsteady-state pressure distribution created by a well with an inclined fracture. Soc. Petr. Eng. SPE-5591 doi:10.2118/5591-MS.
- Cinco-Ley H, Samaniego F, Dominiguez N., 1998. Transient pressure behaviour for a well with a finite-conductivity vertical fracture. J. Soc. Petr. Eng., 253–264.
- Delay F., Kaczmaryk A., Ackerer Ph., 2007. Inversion of interference hydraulic pumping tests in both homogeneous and fractal dual media. Adv. Wat. Res., 30, 314–334. doi:10.1016/j.advwatres.2006.06.008
- Deruyck B., Ehlig-Economides C., Joseph J., 1992. Testing design and analysis. Oilfield and analysis. 28-45.
- Dewandel B., Lachassagne P., Zaidi F.K., Chandra S. 2011. A conceptual hydrodynamic
- model of a geological discontinuity in hard rock aquifers: Example of a quartz reef in granitic
- terrain in South India. J. of Hydrol., 405, 474-487.
- Dewandel, B., Aunay, B., Maréchal, J.C., Roques, C., Bour, O., Mougin, B., L. Aquilina,
- 2014. Analytical solutions for analysing pumping tests in a sub-vertical and anisotropic fault
- zone draining shallow aquifers. J. of Hydrol. 509, 115–131.
- Dewandel, B., Alazard, M., Lachassagne, P., Bailly-Comte, V., Couëffé, R., Grataloup, S.,
- Ladouche, B., Lanini, S., Maréchal, J.C., Wyns, R., 2017. Respective roles of the weathering
- profile and the tectonic fractures in the structure and functioning of crystalline thermo-mineral
- carbo-gaseous aquifers. J. of Hydrol. 547, 690–707.
- Gringarten, A.C., Ramey, H.J. 1973. The use of Source and Green's functions in solving unsteady-flow problems in reservoirs. J. Soc. Petr. Eng. 13, 285-296. SPE-3818-PA. http://dx.doi.org/10.2118/3818-PA
- Gringarten, A.C., Ramey, H.J., 1974. Unsteady pressure distribution created by a well with a single horizontal fracture, partial penetration and restricted entry. J. Soc. Petr. Eng. 14, 413– 426.
- Gringarten, A.C., Henry, J., Ramey, H.J., Raghavan, R., 1974. Unsteady state pressure distributions created by a well with a single infinite conductivity vertical fracture. J. Soc. Petr. Eng. 14, 347–360.
- Hamm, S.Y., and Bidaux, P., 1996. Dual-porosity fractal models for transient flow analysis in fissured rocks. Water Resources Research, 32: 2,733-2,745.
- Hantush, M.S., 1961. Aquifer tests on partially penetrating wells. Proc. Am. Soc. Civil Engin., 87, 171-195.
- Hunt, B., 2005. Flow to vertical and nonvertical wells in leaky aquifers. J. Hydrol. Eng., [doi.org/10.1061/\(ASCE\)1084-0699\(2005\),](https://doi.org/10.1061/(ASCE)1084-0699(2005)) vol.10, 477-484.
- Hunt, B., Scott, D., 2007. Flow to a well in a two-aquifer system. J. Hydrol. Eng., doi.org/10.1061/(ASCE)1084-0699(2007), vol. 12, 146–155.
- Jourde, H., Cornaton, F., Pistre, S., Bidaux, P., 2002. Flow behavior in a dual fracture network. J. of Hydrol., 266, 99-119.
- Lachassagne, P., Wyns, R., Dewandel, B., 2011. The fracture permeability of hard rock aquifers is due neither to tectonics, nor to unloading, but to weathering processes. Terra Nova, 23, 145-161
- Lashgari, H.R., El Rabaa, W., Chan, H., Vaidya, R., 2014. Estimation of hydraulic fracture
- contribution in medium to high permeability reservoirs. Soc. of Petr. Eng., n°169554-MS.
- Lolon, E.P., Archer, R.A., Ilk, D., Blasingame, T.A., 2008. New semi-analytical solutions for multilayer reservoirs. Soc. of Petr. Eng., n°114946.
- Maréchal J.C., Dewandel, B., Subrahmanyan K.,. 2004. Contribution of hydraulic tests at
- different scale to the characterisation of fracture network properties in hard-rock aquifers.
- Water Resourc. Res., 40 (W11508), 1-17
- Maréchal, J.C., Lachassagne, P., Ladouche, B., Dewandel, B., Lanini, S., Le Strat, P., Petelet-
- Giraud, E., 2014. Structure and hydrogeochemical functioning of a sparkling natural mineral
- water system determined using a multidisciplinary approach: a case study from southern
- France. Hydrogeol. J., 22, 47–68, DOI 10.1007/s10040-013-1073-1
- Newman, A.B. 1936. Heating and cooling rectangular and cylindrical solids. Ind. Eng. Chem. 28, 545–548. http://dx.doi.org/10.1021/ie50317a010.
- Moench, A.F., 1984. Double-porosity models for a fissured groundwater reservoir with fracture skin, Water Resourc. Res., 20, 831-846.
- Muskat, M., 1937. The Flow of Homogeneous Fluids through Porous Media, McGraw-Hill Book Company Inc., New York.
- PetroWiki. Solving unsteady flow problems with Green's and Source functions. [http://petrowiki.org;](http://petrowiki.org/) SPE International.
- Rafini, S., Larocque M., 2012. Numerical modeling of the hydraulic signatures of horizontal and inclined faults. Hydrogeol. J., 20, 337–350.
- Rafini, S., Chesnaux, R., Ferroud, A., 2017. A numerical investigation of pumping-test responses from contiguous aquifers. Hydrogeol J (2017) 25:877–894. DOI 10.1007/s10040- 017-1560-x.
- Renard, Ph., Glenz, D., Mejias, M., 2009. Understanding diagnostic plots for well-test interpretation. Hydrogeol. J., 17, 589–600.
- Roques, C., Bour, O., Aquilina, L., Dewandel, B., 2016. High-yielding aquifers in crystalline
- basement: insights about the role of fault zones, exemplified by Armorican Massif, France. Hydrogeol. J.. DOI 10.1007/s10040-016-1451-6.
- Russell, D.G., Truitt, N.E., 1964. Transient pressure behavior in vertically fractured reservoirs, J. Pet. Tech. 1159-1170; Trans., AIME, Vol. 231.
- Strelsova, T.D., 1988. Well Testing in Heterogeneous Formations. Exxon Monograph, Ed. John Wiley & Sons, Inc., 413 p.
- Taylor, R., Howard, K., 2000. A tectono-geomorphic model of the hydrogeology of deeply
- weathered crystalline rock: Evidence from Uganda. Hydrogeol. J., 8, 279-294.
- Theis, C.V., 1935. The relation between the lowering of the piezometric surface and the rate and duration of discharge of a well using groundwater storage. Trans. Am. Geoph. Union, 16, 519-524.
- Tiab, D., 2005. Analysis of pressure derivative data of hydraulically fractured wells by the Tiab's Direct Synthesis technique. J. of Petrol. Sci. and Eng., 49, 1–21.
- Thiéry, D., 1980. Analysis of a pumping test in a horizontal fracture. Thomas W Doe. Third
- Invitational Well-Testing Symposium well testing in low permeability environments, 26-28,
- 1980, March 1980, Berkeley (California), United States.
- Warren, J.E., P.J. Root, 1963. The behaviour of naturally fractured reservoirs, J. Soc. Petrol. Eng., 3, 245-255.

#### **Figure captions**

- **Figure 1.** Conceptual sketch of a well intercepting a single vertical fracture; it corresponds to the analytical solution of Gringarten et al. (1974) (plan view).
- **Figure 2.** A single fully penetrating fracture in a dual-porosity aquifer (dual-porosity
- conceptual model: Moench, 1984), pumping at the centre of the fracture. Type curves for
- 531 various  $\lambda$  coefficients  $(\lambda = K_m r^2 / (K_f b_m))$ . Dimensionless drawdown –s<sub>D</sub>: plain curves and
- 532 derivatives– $s_D$ : dotted curves. Circles correspond to the Gringarten et al. (1974) solution.
- 533 Dual-porosity model parameters:  $K_f$  and  $K_m$ , fracture and matrix hydraulic conductivity,  $b_m$ .
- 534 block size, *r*: radial distance from the well,  $s_f$  and  $s_m$ : fracture and matrix storage,  $sk_f$ = fracture skin.
- **Figure 3.** A single partially entering vertical fracture in a multilayer aquifer. Fracture in the
- deepest layer, and pumping at the centre of the fracture. a) Conceptual model with aquifer
- parameters (see text for explanation). b) Type curves for various *hf/B* ratios. Circles
- correspond to the Gringarten et al. (1974) analytical solution.
- **Figure 4.** Pumping at the centre of an inclined fracture in a leaky aquifer system. a)
- Conceptual model with aquifer parameters (see text for explanation). b) Type curves for
- 542 various angles of the fracture with the vertical,  $\theta$ , and  $k'/B' = 0$  (no-leakage). Dots correspond
- to the Gringarten et al. (1974) and Thiéry (1980) analytical solutions.
- **Figure 5.** Two inclined fractures in a leaky aquifer system: model parameters.

 **Figure 6.** Two inclined fractures in a leaky aquifer system; the well is screened in front of both fractures only. a) Average drawdown and its derivative at the well. b) Dimensionless flowrate contributions of each fracture. c), d), e) and f) Drawdown distributions at various 548 times along the x-axis. Pumping rate is  $1 \text{ m}^3/\text{h}$ .

- **Figure 7.** Two inclined fractures in a leaky aquifer system, pumping in the fractures and in
- the aquifer through a fully penetrating well. a) Average drawdown and its derivative at the
- well. b) Dimensionless flowrate contributions of each fracture and of the aquifer through the
- well. c), d), e) and f) Drawdown distributions at various times along the x-axis. Pumping rate
- 553 is  $1 \text{ m}^3/\text{h}$ .
- Caption for the figure in Appendix C: drawdown and its derivative (a) and flowrates (b)
- behaviour for a well screened in front of both fractures (length: 7 m); both fractures and the
- 556 aquifer though the screened-interval are pumped. Fractures and aquifer properties are the ones
- 557 of Figure 5. Pumping rate is  $1 \text{ m}^3/\text{h}$ .
- 558



560 **Figure 1.** Conceptual sketch of a well intercepting a single vertical fracture; it corresponds to 561 the analytical solution of Gringarten et al. (1974) (plan view).

562



563

564 **Figure 2.** A single fully penetrating fracture in a dual-porosity aquifer (dual-porosity 565 conceptual model: Moench, 1984), pumping at the centre of the fracture. Type curves for 566 various  $\lambda$  coefficients  $(\lambda = K_m r^2 / (K_f b_m))$ . Drawdown –s<sub>D</sub>: plain curves and derivatives–s<sub>D</sub>': 567 dotted curves. Circles correspond to the Gringarten et al. (1974) solution. Dual-porosity 568 model parameters:  $K_f$  and  $K_m$ , fracture and matrix hydraulic conductivity,  $b_m$ : block size, *r*: 569 radial distance from the well,  $s_f$  and  $s_m$ : fracture and matrix storage,  $sk_f$ = fracture skin.







572 **Figure 3.** A single partially entering vertical fracture in a multilayer aquifer. Fracture in the 573 deepest layer, and pumping at the centre of the fracture. a) Conceptual model with aquifer

- 574 parameters (see text for explanation). b) Type curves for various  $h_f/B$  ratios. Circles
- 575 correspond to the Gringarten et al. (1974) analytical solution.



**Figure 4.** Pumping at the centre of an inclined fracture in a leaky aquifer system. a)

Conceptual model with aquifer parameters (see text for explanation). b) Type curves for

579 various angles of the fracture with the vertical,  $\theta$ , and  $k'/B' = 0$  (no-leakage). Dots correspond

to the Gringarten et al. (1974) and Thiéry (1980) analytical solutions.





582



585

586 **Figure 6.** Two inclined fractures in a leaky aquifer system; the well is screened in front of 587 both fractures only. a) Average drawdown and its derivative at the well. b) Dimensionless 588 flowrate contributions of each fracture. Pumping rate is  $1 \text{ m}^3/\text{h}$ .



591 **Figure 6. cont'd.** c), d), e) and f) Drawdown distributions at various times along the x-axis.



593 **Figure 7.** Two inclined fractures in a leaky aquifer system, pumping in the fractures and in 594 the aquifer through a fully penetrating well. a) Average drawdown and its derivative at the 595 well. b) Dimensionless flowrate contributions of each fracture and, of the aquifer though the 596 well. Pumping rate is  $1 \text{ m}^3/\text{h}$ .



599 **Figure 7. cont'd.** c), d), e) and f) Drawdown distributions at various times along the x-axis.

## 601 **Appendices:**

- 602 **Appendix A: Dimensionless drawdown and time**
- 603 Dimensionless drawdown, *s<sup>D</sup>* :

$$
604 \qquad s_D = \frac{2\pi T}{Q} s
$$

605 With *s*, the drawdown, *T*, the transmissivity, and *Q*, the pumping rate.

- 606 Dimensionless time, *tDxf* and *tDLl* :
- 607 For a vertical fracture fully penetrating the aquifer, *tDxf* :
- 2  $\frac{Dxf}{Sx}$ 608  $t_{Dxf} = \frac{tT}{f(x)}$ ; with *x<sub>f</sub>*, the half-fracture length, *t*, the time and *S*, the storage coefficient of

609 the aquifer.

610 - For a fracture partially penetrating the aquifer, *tDLl* :

611 
$$
t_{DII} = \frac{tT}{S(L/2.l/2)};
$$
 with *L* and *l*, the length and the width of the fracture.

- 612
- 

#### 613 **Appendix B: point-source solution of Hunt (2005)**

614 Point-source (or sink source) solution for the aquifer described in Figure 4 (Hunt, 2005), from 615 which 3-D anisotropy in hydraulic conductivity has been implemented:

616 
$$
s_{\text{sink}}(x_D, y_D, z_D, t_D) = \frac{Q}{2\pi T} \sum_{n=1}^{\infty} \frac{\cos(\alpha_n Z/B)\cos(\alpha_n z_D)}{1 + \frac{\sin(2\alpha_n)}{2\alpha_n}} W\left(\frac{(r/B)^2}{4t_D}, \alpha_n \frac{r}{B} \sqrt{k_z/k_{xy}}\right)
$$

where  $x_D = \frac{x}{B} \sqrt{k_y/k_{xy}}$  $x_D = \frac{x}{B} \sqrt{k_y / k_{xy}}$ ,  $y_D = \frac{y}{B} \sqrt{k_x / k_{xy}}$  $y_p = \frac{y}{R} \sqrt{k_x/k_{xy}}$ , *B*  $z_D = \frac{z}{R}$ , are dimensionless variables of the point 617

618 location  $(x, y, z)$ ;  $k_x$ ,  $k_y$ ,  $k_z$  are the hydraulic conductivity along the *x*, *y*, *z* directions; *Z* is the *z* 619 coordinate of the point source where pumping  $Q$  takes place;  $k_{xy}$  is the horizontal hydraulic 620 conductivity of the aquifer ( $k_{xy} = \sqrt{k_x k_y} = T/B$ ); *B* is the aquifer thickness and  $t_D = tk_{xy}/SB$ ;

*t* is time;  $\alpha_n$  is the root of equation  $\alpha_n \tan(\alpha_n)$  $k, \sqrt{B}$  $k$   $\!/B$ *z*  $n \cdot \mathbf{un}(\mathcal{U}_n)$  $\alpha_n \tan(\alpha_n) = \frac{k/B'}{k}$ , with *k'* and *B'* the hydraulic 621 622 conductivity and the thickness of the leaky aquifer; *W(a, b)* is the Hantush leaky-aquifer well 623 function.

624

 **Appendix C:** drawdown and its derivative (a) and flowrates (b) behaviour for a well screened in front of both fractures (length: 7 m); both fractures and the aquifer though the screened- interval are pumped. Fractures and aquifer properties are the ones of Figure 5. Pumping rate is  $1 \text{ m}^3/\text{h}$ .



629

# 631 **Supplemental material:**

632 **Demonstrating that the well Theis funtion integrated along a fracture plane is identical** 

633 **to the analytical solution of pumping in a vertical fracture proposed by Gringarten et al.**  634 **(1974) with uniform flux distribution.**

- 635 Definition conceptual model: the borewell intercepts a vertical fracture of length *2xf* in a
- 636 homogeneous and infinite aquifer of transmissivity, *T*, and storativity, *S*. The width of the
- 637 fracture is assumed to be negligible.





## 639 **Notations**

640 
$$
r^2 = (x - x_{obs})^2 + y_{obs}^2
$$
 with x a point in the fracture ;  $t_D = \frac{Tt}{x_f^2} ; x_D = \frac{x_{obs}}{x_f} ; y_D = \frac{y_{obs}}{x_f}$ ;

641 
$$
u = \frac{r^2.S}{4Tt}
$$
 = >  $u = \frac{(x - x_{obs})^2 + y_{obs}^2}{4Tt}S = \frac{(x/x_f - x_D)^2 + y_D^2}{4t_D}$  and  $a = u.t_D = \frac{(x/x_f - x_D)^2 + y_D^2}{4}$ 

642

# 643 **Properties of the well function** *W(u)***:**

644 
$$
W(u) = E_1(u) = -E_i(-u) = \int_u^{\infty} \frac{e^{-y}}{y} dy
$$

With the following change of variable:  $y = \frac{a}{\tau}$  $y = \frac{a}{\tau}$ , it becomes that  $\int_{a}^{\infty} \frac{e^{-y}}{y} dy = \int_{a}^{b}$  $\infty$   $-y$   $t_{D}$   $=$  $\int_{0}^{b}$ *D*  $t_{D}$ <sub> $a$ </sub> $-a$ *a t*  $\int d^y = \int^{\frac{t_D}{2}} \frac{e^{-a/\tau}}{t} dt$ *y e* 0 / / τ τ τ 645

## 646 **Properties of the** *Erf* **function:**

647 
$$
Erf(x) = \frac{2}{\sqrt{\pi}} \int_{0}^{x} e^{-u^{2}} du = -\frac{2}{\sqrt{\pi}} \int_{-x}^{0} e^{-u^{2}} du
$$

$$
648 \quad \quad Erf(x) = -Erf(-x)
$$

### 650 **Well Theis function for a well**

651 The aquifer is pumped in a single vertical well fully penetrating the aquifer with a pumping 652 rate  $q$ . The drawdown at a radial distance  $r$  from the well can be expressed as:

653 
$$
\begin{cases} s(r,t) = \frac{q}{4\pi T} W(u) = \frac{q}{4\pi T} \int_{u}^{\infty} \frac{e^{-y}}{y} dy \\ u = \frac{r^2 S}{4Tt} \end{cases}
$$
 [Eq.1]

654

#### 655 **Gringarten analytical solution (Gringarten et al., 1974)**

656 For a pumping in a single vertical fracture intersecting the pumped well, with uniform flux 657 distribution, the drawdown can be expressed as:

658 
$$
s(x, y, t) = \frac{Q}{2\pi T} s_D
$$
 with  $s_D = \int_0^{t_D} e^{-\frac{y_D^2}{4t_D}} \left[ Erf\left(\frac{1-x_D}{2\sqrt{\tau}}\right) - Erf\left(\frac{1+x_D}{2\sqrt{\tau}}\right) \right] \frac{\sqrt{\pi}}{4\sqrt{\tau}} d\tau$   
\n[Eq.2]

659 Where *Q* is the pumping rate. As *Q* is assumed to be uniformly distributed along the fracture, 660 therefore:

661 
$$
Q = \int_{-x_f}^{+x_f} q(x) dx \Rightarrow q(x) = q = \frac{Q}{2x_f}
$$

#### 662 **Mathematical demonstration**

663 Integration of the Theis' solution along the fracture yields:  
\n664 
$$
s_1(x, y, t) = \int_{-xf}^{+xf} s(x, y, t) dx = \frac{1}{4\pi T} \int_{-xf}^{+xf} qE_1 \left( \frac{(x - x_{obs})^2 + y_{obs}^2}{4Tt} S \right) dx
$$

665 with the simplified notation and using the property of the  $E_i$  function, it becomes:

665 with the simplified notation and using the property of the E<sub>i</sub> function, it becomes:  
\n666 
$$
s_{I} = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-x}^{+xf} E_{1}\left(\frac{a}{t_{D}}\right) dx = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} \left(\int_{a/t_{D}}^{+\infty} \frac{e^{-u}}{u} du\right) dx = \frac{Q}{2\pi T} \frac{1}{4x_{f}} \int_{-xf}^{+xf} \left(\int_{0}^{t_{D}} \frac{e^{-a/r}}{\tau} d\tau\right) dx
$$

667 which can be rewritten as : 
$$
s_I = \frac{Q}{2\pi T} \frac{1}{4x_f} \int_{-x_f}^{+x_f} \left( \int_{0}^{t_D} \frac{e^{-\left[ \left( \frac{x}{\lambda x_f} - x_D \right) + y_D^2 \right] \right/4\tau}}{\tau} d\tau \right) dx
$$

668 According to the Fubini theorem (i.e.  $\tau$  and  $x$  are independent), the order of integration can be 669 inverted, yielding to:

 *d x d e e T x Q <sup>d</sup> d x <sup>e</sup> e T x Q s D f f D f D f f D D f D t x x x x x y f x x x x x t y f <sup>I</sup>* 0 4 4 4 0 4 4 1 4 2 1 2 2 2 2 2 670

671 and as  $e^{-y_D^2/4\tau}$  does not depend on x, it can be rearranged as:

672 
$$
s_{I} = \frac{Q}{2\pi T} \frac{1}{x_{f}} \int_{0}^{t_{D}} e^{-y_{D}^{2}/4\tau} \left( \int_{-x_{f}}^{+x_{f}} \frac{e^{-\left(\frac{x}{x_{f}} - x_{D}\right)^{2}}/4\tau}{4\tau} d\tau \right) d\tau
$$
 [Eq.3]

With the following change of variable:  $v = \left(\frac{x}{x_f} - x_b\right) / 2\sqrt{\tau}$  $\left(\frac{x}{x_f} - x_D\right)$ 674 With the following change of variable:  $v = \left(\frac{x}{x_f} - x_D\right) \left/ 2\sqrt{\tau}$ , the (a) term of [Eq.3] can be 675 rewritten as :

676 
$$
\int_{-x_f}^{+x_f} \frac{e^{-\left(\frac{y}{\lambda x_f} - x_D\right)^2/4\tau}}{4\tau} dx = \frac{x_f}{4\sqrt{\tau}} \int_{-(1+x_D)/2\sqrt{\tau}}^{(1-x_D)/2\sqrt{\tau}} 2e^{-v^2} dv
$$
 [Eq.4]

677 The right part of [Eq.4] can be decomposed into two terms related to the *Erf* function:

678 The right part of [Eq.4] can be decomposed into two terms related to the *Erf* function:  
\n678 
$$
\frac{x_f}{4\sqrt{\tau}} \int_{-(1+x_D)/2\sqrt{\tau}}^{(1-x_D)/2\sqrt{\tau}} 2e^{-v^2} dv = \frac{x_f}{4\sqrt{\tau}} \left( \int_{-(1+x_D)/2\sqrt{\tau}}^{0} 2e^{-v^2} dv + \int_{0}^{(1-x_D)/2\sqrt{\tau}} 2e^{-v^2} dv \right) = \frac{x_f \sqrt{\pi}}{4\sqrt{\tau}} \left[ -Erf \left( \frac{(1+x_D)}{2\sqrt{\tau}} \right) + Erf \left( \frac{(1-x_D)}{2\sqrt{\tau}} \right) \right]
$$
\n679 [Eq.5]

680 Combining [Eq.3] and [Eq.5], we obtain:

$$
681 \t sI = \frac{Q}{2\pi T} \frac{1}{x_f} \int_0^{t_p} e^{-\frac{y_D^2}{4\tau}} \frac{x_f \sqrt{\pi}}{4\sqrt{\tau}} \left[ -Erf \left( \frac{(1+x_D)}{2\sqrt{\tau}} \right) + Erf \left( \frac{(1-x_D)}{2\sqrt{\tau}} \right) \right] d\tau
$$
  
\n682 And thus : 
$$
s_I = \frac{Q}{2\pi T} \int_0^{t_p} e^{-\frac{y_D^2}{4\tau}} \left[ Erf \left( \frac{1-x_D}{2\sqrt{\tau}} \right) - Erf \left( \frac{1+x_D}{2\sqrt{\tau}} \right) \right] \frac{\sqrt{\pi}}{2\sqrt{\tau}} d\tau
$$
 [Eq.6]

τ τ  $\frac{Q}{\pi T} \int_{0}^{t_{D}} e^{-\frac{y_{D}}{4\tau}} \Bigg[ Erf\Bigg(\frac{1-x_{D}}{2\sqrt{\tau}}\Bigg) - Erf\Bigg(\frac{1+x_{D}}{2\sqrt{\tau}}\Bigg) \Bigg| \frac{\sqrt{\pi}}{4\sqrt{\tau}} d\tau$ *T*  $s_I = \frac{Q}{2\pi T} \int_0^{t_D} e^{-\frac{y_D}{4\tau}} \left[ Erf\left(\frac{1-x_D}{2\sqrt{\tau}}\right) - Erf\left(\frac{1+x_D}{2\sqrt{\tau}}\right) \right]$  $\mathsf{L}$ L  $\overline{\phantom{a}}$ J  $\left(\frac{1+x_D}{\sqrt{2}}\right)$  $\setminus$  $-\mathrm{Erf}\left(\frac{1+}{2}\right)$ J  $\left(\frac{1-x_D}{\sqrt{1-x_D}}\right)$  $\setminus$  $=\frac{Q}{2\pi}\int_{0}^{t_{D}}e^{-\frac{y_{D}}{4\tau}}Erf(\frac{1}{\tau})$  $\mathbf{0}$  $2\sqrt{\tau}$  ) | 4 2 2 682 And thus:  $s_l = \frac{Q}{2\pi} \left[ e^{-4\tau} \left[ Erf \right] \frac{1 - x_D}{2\pi} \left[ -Erf \left[ \frac{1 + x_D}{2\pi} \right] \right] \frac{v}{\sqrt{n}} d\tau \right]$  [Eq.6]

683 [Eq.6] proves that the Theis analytical solution integrated along the fracture plane is exactly 684 the same as the Gringarten analytical solution, for a vertical fracture with a uniform flux 685 distribution (Eq. 2 or the Eq. 20 in Gringarten et al. 1974).