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Predicting the surface tension of aqueous 1-1 electrolyte solutions at high salinity

Philippe Leroy ^{a*}, Arnault Lassin ^a, Mohamed Azaroual ^a, Laurent André ^a

^a BRGM, 3 avenue C. Guillemin, BP 6009, 45061 Orléans, France

*** Corresponding author:**

Philippe Leroy

BRGM

Water Division (EAU/M2H)

3 Avenue Claude Guillemin

45060 Orléans Cedex 2, France

E-mail: p.leroy@brgm.fr

Tel: +33 (0)2 38 64 39 73

Fax: +33 (0)2 38 64 37 19

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ABSTRACT

The surface tension of the air/water interface is a phenomenon of particular interest in the water-unsaturated zone of porous media because it influences the contact angle and consequently the capillary water volume. A mechanistic model based on the modified Poisson-Boltzmann equation and the Pitzer theory is described and used to predict, under isothermal and isobaric conditions, the surface tension of 1:1 electrolytes at high salinity. These theories enable the determination of the electrical potential at the air/water interface and the activity coefficient of the ionic species in the bulk pore water, respectively. Hydration free energies of the structure-making and structure-breaking ions that influence the surface tension at high salinity are taken into account. Structure-making ions flee the air/water surface because they can better organize the water dipoles in bulk water than at the interface. Structure-breaking ions are positively adsorbed at the air/water interface because the bulk water can better organize their hydrogen bonding network without these ions. The resulting surface tension increases and decreases, respectively, compared to the surface tension of pure water. The model predictions are in good agreement with the surface tension data for 1:1 electrolytes (NaCl, KCl, HCl, NaNO₃, KNO₃, HNO₃ electrolytes) and the optimized parameters depend on the effective electrostatic diameters of cations and on the hydration free energies of the ions at the interface.

1. INTRODUCTION

The surface tension of aqueous electrolyte solutions is an important physical and chemical property of the porous medium, particularly in the unsaturated zone (UZ), that retains aqueous solutions against gravity and evaporation by capillary forces. For example, unsaturated flow of water induced by hydraulic head gradients caused by concentration-dependent surface tension changes has been demonstrated in the laboratory (Smith and Gillham, 1999; Henry et al., 2001). These authors studied the transport properties of porous media contaminated with organic solutes that reduce the surface tension of water. The surface tension of the gas/water interface combined with the mineral/water and mineral/gas interfacial energies also enables the calculation of the contact angle (Douillard and Médout-Marère, 2000; Mercury et al., 2003; Ward and Sefiane, 2009). Environmental problems like salt weathering need to characterize the drying kinetics of porous media (Sghaier et al., 2006). In this context, crystallisation of salt in the pores are responsible for the deterioration of buildings materials. According to Prat (2007), pore shape and contact angle can greatly affect the drying kinetics. This author used pore network simulations to explore the influence of these geometrical and physico-chemical parameters on drying rates. Predicting the surface tension of the gas/water interface might also be of crucial importance in some industrial contexts. For example, capillary failure in the context of CO₂ geological storage is, to a large extent, controlled by the gas/water surface tension. The higher the surface tension, the higher the pressure difference between the non-wetting (gas) and wetting (brine) phases is to attain capillary failure (André et al., 2007; Chiquet et al., 2007; Pettenati et al., 2008). Another application is the study of heterogeneous chemistry involving sea-salt aerosols, which influences the chemical composition of the atmosphere (Jungwirth and Tobias, 2001, 2006).

The surface tensions of 1:1 aqueous electrolytes are usually higher than that of pure water. Any solute that increases the surface tension of water may exhibit a negative total adsorption at the air/water interface, according to the Gibbs adsorption equation (Weissenborn and Pugh, 1996). The depletion of such ions near the water's surface was long attributed to repulsive electrostatic image forces (Onsager and Samaras, 1934). Their theory was criticised, particularly at high salinity (> 0.2 M), for its oversimplification, mathematical treatment and neglect of water's structure. Recent molecular dynamics simulations have shown that the common belief that the interfacial region is depleted of ions might not be true (Jungwirth and Tobias, 2001, 2006). At high ionic strengths, the effect of ions on water's structure (structure-making or structure-breaking ions) may be responsible for the variation of surface tension with salinity (Horne and Young, 1972; Ralston and Healy, 1973; Manciu and Ruckenstein, 2003, 2005). The effect of ion hydration on interfacial tension can be easily understood qualitatively. Structure-making ions flee the interface because they can better organize the water dipoles in bulk water than at the interface (Fig. 1). They orient the hydrogen of the water molecules of the solvation shell towards their neighbouring water molecules in the aqueous solution, with which they can create hydrogen bonds. The opposite is true for the structure-breaking ions. The total free energy of the system is minimized by pushing the structure-breaking ions toward the interface because the bulk water can better organize their hydrogen bonding network without these ions.

Some macroscopic models can explain the variation of the surface tension with the salinity of the solution. Yu et al. (2000) studied the ionic concentration dependence of the surface tension of single and mixed electrolyte solutions. They utilized the mean spherical approximation method (Lu et al., 1993) to estimate the activity coefficients of water in the air/water interface and bulk liquid phase. For that purpose, they used the phenomenological surface-phase method (Guggenheim and Adam, 1933) on the basis of Gibbs thermodynamic

analysis. Li and Lu (2001) employed a combination of the Gibbs chemical potential concept and the Langmuir adsorption model (Desnoyer et al., 1997) to study the surface tension of single and mixed electrolyte solutions. The activity coefficients were calculated thanks to the Pitzer theory (Pitzer, 1973). However, Yu et al. (2000) and Li and Lu (2001) used some empirical parameters in the adsorption equilibrium constants that are not specific to the effect of ions on water's structure. Recently, continuum electrostatic models (in which the solvent is considered to be a continuum) have been developed based on the modified Poisson-Boltzmann (MPB) approach (Bhuiyan et al., 1991; Bhuiyan et al., 1992; Karraker and Radke, 2002; Markin and Volkov, 2002; Manciu and Ruckenstein, 2003, 2005). These models consider ion-specific behaviour including, in some cases, a propensity of certain types of ions for the air/water interface. Bhuiyan et al. (1991) studied the surface tension of aqueous electrolytes within the framework of the primitive model of the planar electric double layer. The excess surface tension was determined by numerical integration of the Gibbs adsorption isotherm. The primitive model they used can describe very well the surface properties of the air/water interface. The MPB allows the determination of the singlet density profiles and the surface deficit of the ions repelled from the polarized interface. The bulk MPB equation was also used for the calculation of the activity coefficients in the bulk water. Karraker and Radke (2002) developed a self-consistent electrostatic theory to predict successfully disjoining pressure isotherms of aqueous thin-liquid films stabilized by non-ionic surfactants, air/water surface tensions and zeta potentials of single electrolyte solutions. Their theory is based on specific adsorption of hydroxide ions at the interface, and image plus dispersion forces due to Van der Waals interactions on ions. Markin and Volkov (2002), in their model, included image forces, ions of finite radii and the hydrophobic effect at liquid interfaces. Manciu and Ruckenstein (2003, 2005) considered specific ion effects via the change in the ion hydration between the bulk and the vicinity of the surface. Their approach can account for both the

depletion of structure-making ions as well as for the accumulation of structure breaking ions near the air/water surface. According to the models mentioned above, electrostatic forces are of key importance for explaining the distribution of ionic species at the interface and hence the resulting surface tension. For example, at high salinity (> 0.2 M), differences in the hydration free energies of the ions lead to asymmetric ion depletions (Manciu and Ruckenstein, 2003) and therefore generate a macroscopic electrical potential. This electrical potential modifies the ionic distribution, and, consequently, the surface tension.

In some contexts, such as the geological sequestration of CO_2 in deep saline aquifers, massive evaporation of native water may occur, generating residual brines having very high ionic strengths (Azaroual et al., 2004; see Andre et al., 2007 for further details). Furthermore, many systems encountered in industrial processes contain aqueous solutions of highly concentrated electrolytes (Azaroual et al., 1997; Azaroual et al., 2004). According to the Gibbs adsorption equation, the activity coefficient of the ionic species in the bulk water also influences the surface tension. Consequently, there is a need, especially at high ionic strength (> 0.2 M), to take non-ideality into account. The Pitzer theory (Pitzer, 1973, 1995) and the resulting chemical equilibrium models of solution behaviour and solubility (Harvie and Weare, 1980; Harvie et al., 1984; Plummer et al., 1988; Felmy et al., 1994; Christov and Moller, 2004; Carroll et al., 2005) are particularly well-adapted to taking into account short-range interactions between ions like, for example, the interaction of pairs of oppositely charged ions in the calculation of activity coefficients.

To our knowledge there has been no attempt, to date, to predict the surface tension of 1:1 electrolytes at high salinity using a mechanistic approach based on both the modified Poisson-Boltzmann and Pitzer theories. We propose here a unified and consistent model of the electrochemical properties of the air/water interface at high electrolyte concentrations. We take into account the specific influence of anions and cations close to the water's surface

(structure-breaking and structure-making ions) on the calculation of the surface tension. After a brief description of the theoretical background, the model is validated by comparison with experimental surface tension data of NaCl, KCl, HCl, NaNO₃, KNO₃, HNO₃ aqueous electrolytic solutions. The specific physico-chemical properties of the ions constituting the different electrolytes at the interface and in the bulk are investigated.

2. THEORETICAL BACKGROUND

2.1. Gibbs adsorption equation

Under isobaric and isothermal conditions, the dependence of the surface (or interfacial) tension σ (in N m⁻¹) on the electrolyte concentration can be calculated using the Gibbs adsorption equation (Weissenborn and Pugh, 1996; Li and Lu, 2001; Manciu and Ruckenstein, 2003)

$$d\sigma = -\sum_i \Gamma_i d\mu_i = -\sum_{j \neq w} \Gamma_j d\mu_j - \Gamma_w d\mu_w, \quad (1)$$

where Γ_i is the surface excess/deficiency of the species i , i.e. of ion (“ j ”) and water (“ w ”) (in sites m⁻²), and μ_i is the chemical potential of the species at the interface (in J mol⁻¹). The surface excess/deficiency of the species is expressed by (Manciu and Ruckenstein, 2005)

$$\Gamma_i = \int_0^{5\lambda_d} [c_i(x) - c_i^b] dx, \quad (2)$$

where c is the concentration (molarity in mol m^{-3}) and x the distance of the species from the air/water surface corresponding to the Gibbs dividing surface (in m). The superscript “ b ” refers to concentration in the bulk water where electroneutrality occurs (not influenced by the air/water interface). The parameter χ_d represents the Debye length associated with the thickness of the diffuse layer (in m^{-1} , see Hunter, 1981 for further details). We assume here a sharp air/water surface and step distributions for the ions at the interface.

A cross-sectional area of the Gibbs dividing surface contains the same number of moles of water between the air and water phases. Therefore, the Gibbs dividing surface is chosen so that $\Gamma_w = 0$ sites m^{-2} . This implies, according to equation (1):

$$d\sigma = -\sum_{j \neq w} \Gamma_j d\mu_j . \quad (3)$$

The chemical potential of the ionic species at the air/water interface is considered equal to their chemical potential in the bulk electrolyte (Lyklema, 1991). This implies

$$\mu_j = \mu_j^b, \quad (4)$$

$$\mu_j = \mu_{0j}^b + k_b T \ln(a_j^b), \quad (5)$$

where k_b is the Boltzmann constant ($1.381 \times 10^{-23} \text{ J K}^{-1}$), T the temperature in K, and a the activity ($a_j^b = \gamma_j^b m_j^b$ where γ_j is the activity coefficient and m the molality in $\text{mol kg}_{\text{water}}^{-1}$ of species j). The subscript “0” refers to the chemical potential in the standard state, i.e. the unit molar concentration of the ionic species at hypothetical infinite dilution (Helgeson et al., 1981). By combining equations (3) and (5), we obtain, at constant temperature and pressure,

$$d\sigma = -k_b T \sum_j \Gamma_j da_j^b / a_j^b . \quad (6)$$

According to equation (6), the surface excess/deficiency of the ionic species and their activity in the bulk are key parameters to predict the surface tension of the solution at a given salinity. In the next chapters, we will explain how to calculate these parameters using the relevant theories.

2.2. Ionic distributions at the air/water interface

The modified Poisson-Boltzmann (MPB) equation enables the calculation of the ionic concentrations at the air/water interface (in mol m⁻³). At thermodynamic equilibrium, they are described using a modified Boltzmann distribution (Manciu and Ruckenstein, 2003, 2005):

$$c_j(x) = c_j^b \exp\left\{\frac{-[q_j \psi(x) + W_j(x)]}{k_b T}\right\}, \quad (7)$$

where $q_j = \pm e z_j$ is the charge of species j with e the electronic charge (1.602×10^{-19} C) and z_j the valency, $\psi(x)$ the “mean electrical potential” (in V), and $W_j(x)$ an interaction free energy term (in J) that includes all the other interactions of the ion with the surrounding medium. Hereafter, we consider only the effect of the electrical potential and the ionic hydration free energies on interfacial tension because our study is focused on highly saline solutions (ionic strengths >1 M). Image charge interaction free energy (Onsager and Samaras, 1934) is not taken into account in our approach. Indeed, because of its dependence on the Debye length, the image force depends on the ionic strength of the solution and its amplitude decreases as the ionic strength increases. We also believe, using the approach of Manciu and Ruckenstein

(2003), that specific ion effects on water structure can be explained by ionic hydration free energies. Other authors, such as Ninham and Yaminsky (1997) and Karraker and Radke (2002) consider that the Van der Waals interactions between an ion and the interface between two dielectric media (air and water) are not screened by the electrolyte and therefore might become greater than the image force at high salinity. However, as stated by Manciu and Ruckenstein (2003), this hypothesis leads to the debatable prediction that the cations would draw nearer to the interface than the anions. Also, in our modeling approach, the ion correlation effects associated with the exclusion volume of the ions, which could be very important at high electrolyte concentration, is neglected. To include the ion correlation effects, both the modified Poisson-Boltzmann (Bhuiyan et al., 1991; Kralchevsky and Paunov, 1992) and density functional theories (Forsman, 2009; Peng and Yu, 2009) can be used.

According to Poisson's equation, which considers the electrostatic interactions associated with the volume charge density ρ (in C m^{-3}) (Hunter, 1981), the Laplacian of the electrical potential can be described by

$$\nabla^2 \psi = -\frac{\rho}{\varepsilon_w} = -\frac{\sum_j q_j c_j}{\varepsilon_w}, \quad (8)$$

where ε_w is the dielectric permittivity of water ($\varepsilon_w = \sim 78.3 \times 8.85 \times 10^{-12} \text{ F m}^{-1}$ at 298 K and at a pressure of 1 bar). For simplicity, like Manciu and Ruckenstein (2003, 2005), we consider a uniform dielectric constant of value ε_w at the air/water interface. Large electrical fields and concentrations of counterions near the surface actually induce a decrease of the dielectric constant. Indeed, the presence of an electrolyte hinders the polarizability of the water molecules, thus decreasing the dielectric constant. In our model, the influence of the dielectric constant of water at the interface on the ionic distribution is taken into account by the

interaction free energy term in equation (7). Note that the substantial reorientation and restructuring of surface water molecules significantly diminishes the dielectric constant (Jungwirth and Tobias, 2006). Combining equations (7) and (8) leads to the non-linear modified Poisson-Boltzmann equation:

$$\nabla^2 \psi(x) = -\frac{1}{\epsilon_w} \sum_j q_j c_j^b \exp\left\{\frac{-[q_j \psi(x) + W_j(x)]}{k_b T}\right\}. \quad (9)$$

The electrical potential near the surface is solved numerically using equation (9) and the following boundary conditions

$$\nabla \psi|_{x=0} = 0, \quad (10)$$

$$\psi(x) = 0 \text{ at } x = x_{\max}, \quad (11)$$

where x_{\max} is about 5 Debye-Hückel lengths from the air/water surface (Manciu and Ruckenstein, 2005). The previous system of equations (equations (9) to (11)) is solved in our study using a Partial Differential Equations (PDE) solver based on the finite-element method (the Electrostatic module of COMSOL MultiphysicsTM 3.4).

Equation (9) assumes that there is no surface charge density at the air/water surface. However, at low salinity ($<10^{-2}$ M), the surface adsorption of OH⁻ ions may not be neglected (Karraker and Radke, 2002; these authors specifically consider adsorbed hydroxide ions, which lose part of their hydration shells). This generates a surface charge. According to Jungwirth and Tobias (2006), on average, surface water molecules have a net orientation of their dipoles and contribute, therefore, to a macroscopic electrical potential at the air/water interface. Indeed, surface water molecules are reoriented and restructured to maximize their

hydrogen bonding in the under-coordinated environment, as opposed to their isotropic orientation in the bulk. A positive value of the electrical potential of pure water indicates that the liquid is positive relative to vapour and that surface water molecules are oriented with their hydrogen atoms pointing towards the water phase. Consequently, adsorption sites for hydroxide ions may be sparsely spaced groups of surface water molecules available for hydrogen bonding with the hydroxide oxygen. On the other hand, Jungwirth and Tobias (2006) stated that a significant percentage of surface water molecules (~19%) might possess free OH bonds. Nevertheless, by adding an electrolyte, other interactions may dominate the surface charge generated by adsorption of hydroxide ions or by free OH bonds. According to Karraker and Radke (2002), the presence of OH⁻ ions at the interface generates a surface charge of extremely low amplitude. Therefore, at high salinity, following the approach of Manciu and Ruckenstein (2003), the surface charge density can be neglected.

After determining the electrical potential $\psi(x)$, we can calculate the ionic concentrations $c_j(x)$ at the interface (equation (7)) using a given value for $W_j(x)$. By integrating the difference between the ionic concentrations at a distance x from the surface and the ionic concentrations in the bulk electrolyte, we can predict the excess/deficiency of the ionic species near the surface (equation (2)).

2.3. Surface tension of 1:1 electrolytes

For binary electrolytes, the surface tension variation with the electrolyte concentration can be obtained using equation (3):

$$d\sigma = -\Gamma_M d\mu_M^b - \Gamma_X d\mu_X^b. \quad (12)$$

By incorporating equation (5) and the equation $a_j^b = \gamma_j^b m_j^b$ in equation (12), where γ_j is the activity coefficient of species j , we find:

$$d\sigma = -k_b T \left[\Gamma_M (d \ln \gamma_M^b + d \ln m_M^b) + \Gamma_X (d \ln \gamma_X^b + d \ln m_X^b) \right]. \quad (13)$$

For 1:1 electrolytes like NaCl or KCl, the molalities of the cations and anions are equal in the bulk water, i.e. $m_M^b = m_X^b = m_{MX}^b$. This implies, according to equation (13):

$$d\sigma = -k_b T \left[\Gamma_M d \ln \gamma_M^b + \Gamma_X d \ln \gamma_X^b + (\Gamma_M + \Gamma_X) d \ln m_{MX}^b \right]. \quad (14)$$

We also assume that the activity coefficients of cations and anions are equal (see Appendix A). This gives:

$$d\sigma = -k_b T (\Gamma_M + \Gamma_X) (d \ln \gamma_{MX}^b + d \ln m_{MX}^b). \quad (15)$$

The equation for the estimation of the surface tension as a function of the electrolyte concentration can be obtained by integrating equation (15):

$$\Delta \sigma = \sigma - \sigma_w = \int_{\sigma_w}^{\sigma} d\sigma = -k_b T \left[\int_1^{\gamma_{MX}^b} (\Gamma_M + \Gamma_X) d \ln \gamma_{MX}^b + \int_0^{m_{MX}^b} (\Gamma_M + \Gamma_X) d \ln m_{MX}^b \right], \quad (16)$$

$$\Delta \sigma = -k_b T \int_0^{m_{MX}^b} (\Gamma_M + \Gamma_X) \left(\frac{d \ln \gamma_{MX}^b}{dm_{MX}^b} + \frac{1}{m_{MX}^b} \right) dm_{MX}^b, \quad (17)$$

where σ is the surface tension at a given electrolyte concentration and σ_w the surface tension of pure water. According to equation (17), the excess/deficiency of the ionic species at the interface and the variation of the natural logarithm of their activity coefficient with their molality are key parameters for predicting surface tension variation. The molarities of cations and anions (in mol m⁻³) are also equal in the bulk electrolyte and this gives, using equation (2):

$$\Gamma_M + \Gamma_X = \int_0^{5\lambda_d} [c_M(x) + c_X(x) - 2c_{MX}^b] dx. \quad (18)$$

By combining equations (17) and (18), we have:

$$\Delta\sigma = -k_b T \int_0^{m_{MX}} \left\{ \int_0^{5\lambda_d} [c_M(x) + c_X(x) - 2c_{MX}^b] dx \left(\frac{d \ln \gamma_{MX}^b}{dm_{MX}^b} + \frac{1}{m_{MX}^b} \right) \right\} dm_{MX}^b. \quad (19)$$

To summarize, in order to predict the surface tension variation with the salinity in the case of 1:1 electrolytes, we must (i) determine the surface excess/deficiency of the ionic species (molarity, in mol m⁻³), (ii) calculate the activity coefficients in the bulk water (see equations (A1) to (A11), (B1) to (B5), and Table 1 for further details), and (iii) integrate equation (15) from the molality of pure water to the molality of the bulk electrolyte (equation (19)). We also need to convert molarities (mol m⁻³) into molalities (mol kg_{water}⁻¹) (equations (C1) and (C2)). Standard partial molal volumes of water and ionic species at a given temperature and pressure (1 bar) were calculated using the ‘‘Thermo-ZNS’’ computer code (Lassin et al., 2005) (Table 2). This code was developed to determine the relevant thermodynamic properties of water, aqueous species, minerals and gases in capillary systems where interfacial phenomena are of primary importance. Standard partial molal volumes of

water and ionic species as a function of temperature and pressure are determined using the revised Helgeson-Kirkham-Flowers (HKF) model (Helgeson et al., 1981) for predicting the standard partial molal properties of the species (Tanger and Helgeson, 1988). This thermodynamic model takes into account the temperature and pressure dependence of the solvent structure (up to 1000 °C and a pressure of 5 kbar). According to Tanger and Helgeson (1988), the standard partial molal volume of an aqueous ion or electrolyte can be expressed by its nonsolvation and solvation contribution. The electrostatic solvation properties of an ion can be discussed by considering the sign of the solvation contribution to the standard partial molal volume. For example, as seen in Table 2, sodium has a negative standard partial molal volume and is considered to be an electrostatic “negative solvating ion”.

2.4. Structure-making and structure-breaking ions

The surface tension of an inorganic aqueous electrolyte can increase or decrease with the electrolyte concentration. According to the Gibbs adsorption equation, these variations depend on the ionic distribution at the air/water interface, i.e. the possible depletion of the structure-making ions and the positive adsorption of the structure-breaking ions. Are the cations and/or the anions negatively or positively adsorbed near the surface? Positive or negative $d\sigma/dm$ values do not necessarily mean that both cations and anions are negatively or positively adsorbed, but that one ion may be predominate in terms of overall adsorption and effect on surface tension. Recent molecular dynamics simulations have shown that the common belief that the interfacial region is depleted of ions might not be true (Jungwirth and Tobias, 2001).

The close correlation between the magnitude of $d\sigma/dm$ and hydration entropy and enthalpy of the salts in the bulk water has been observed (Hey et al., 1981; Weissenborn and

Pugh, 1996). The most important property of liquid water is its unique ability to form a network of self-associated molecules through hydrogen bonding. Small ions of high valency are highly hydrated and increase the organization of water molecules by compacting water molecules around themselves. They are commonly called structure-making ions. Large ions of low valency are weakly hydrated, have the opposite effect on water's structure and are commonly called structure-breaking ions. It is generally accepted that the oxygen atoms of water molecules surrounding the solvated ion are oriented toward the cations or the hydrogen atoms toward the anions. In this case, the water molecules in the first hydration shell are oriented in the strong electric field of the ion (see Fig. 1), and the 3D network of water molecules is interrupted. Structure-making ions impose more local ordering in surrounding water, i.e. they strengthen the hydrogen bonds between them, while structure breaking ions disorder them. In fact, the surface properties of anions and cations are due to several factors, of which ion size and ion and water polarizability dominate (Boström et al., 2001; Hribar et al., 2002; Karraker and Radke, 2002; Boström et al., 2005; Jungwirth and Tobias, 2006). According to Collins (1985) and Hribar et al. (2002) charge densities and a balance of electrostatic and hydrogen bonding forces control the effects of ions on water structure. These forces are due to water's dipole interacting with ions and neighbouring water molecules, respectively. Ion-water interactions are dominated by charge density effects, and water-water interactions are dominated by hydrogen bonding. Structure-making ions are expected to stay in the bulk water because they can better organize the water dipoles in the bulk water than at the interface, and structure-breaking ions are expected to be pushed to the surface by the bulk water.

Two of the model's parameters are of primary importance if we wish to describe the effect of ion hydration on interfacial tension: a length that corresponds either to the distance of closest approach of a structure-making ion or to the thickness of a layer of positive

adsorption of structure-breaking ions, d_j (in Å), and an interaction free energy term, W_j , of the ion j (in J; see Fig. 2). We use the approach of Manciu and Ruckenstein (2003). These authors explain the physico-chemical meaning of these parameters. The distribution of ions in the vicinity of the surface can be calculated if reasonable estimations of d_j and W_j are available. The hydration free energy per ion, in bulk water, is of the order of ~ 100 - $200 k_bT$ (Markin and Volkov, 2002; Manciu and Ruckenstein, 2003). The hydration free energy of a structure-making ion increases by a few k_bT from its bulk hydration free energy at a distance d_j from the air/water surface (Manciu and Ruckenstein, 2003). This variation of the hydration free energy corresponds to only a fraction of the total ion hydration free energy, but can still be sufficient to cause a significant depletion of ions in the aqueous interface (Markin and Volkov, 2002; Jungwirth and Tobias, 2006). Conversely, the hydration free energy of a structure-breaking ion decreases by a few k_bT from its bulk hydration free energy at a distance d_j from the air/water surface (Manciu and Ruckenstein, 2003, 2005).

3. COMPARISON WITH EXPERIMENTAL DATA

Surface tension data for aqueous electrolyte solutions available in the literature are used to validate the model. Abramzon and Gaukhberg (1993) reported surface tension data for 179 simple inorganic electrolytes (at various temperatures between 0 and 100 °C and at a pressure of 1 bar) over a wide concentration range (from 0 to 36 M).

3.1. NaCl electrolyte

The NaCl-electrolyte surface tension data from Abramzon and Gaukhberg (1993) show a near-linear tendency of the surface tension to increase with the electrolyte

concentration up to a salinity of 3 mol kg_{water}⁻¹. There may, therefore, be a zone depleted of ions at the air/water interface.

Recent molecular dynamics simulations (Jungwirth and Tobias, 2001, 2006) have shown the accumulation of large negative anions (e.g. Br⁻ or I⁻) near the air/water surface and the depletion in this area of some cations (e.g. Na⁺). According to these simulations, chloride ions penetrate the interface but are slightly repelled from the surface. Anions near the surface are still solvated, maintaining a substantial solvation shell. Chlorides do not, however, exhibit a significant concentration peak at the air/water interface. On the other hand, sodium is strongly repelled from the surface. Consequently, a resulting electrical double layer exists at the interface. We, therefore, consider that the cations are negatively adsorbed at the interface (structure-making) and that the anions can penetrate it, i.e. $d_M > d_X$, $W_X = \infty k_b T$ for $0 < x < d_X$ and $W_M = \infty k_b T$ for $0 < x < d_M$. There are three layers at the air/water interface: the uppermost layer having a thickness d_X (I) depleted of ions, a layer having a thickness $d_M - d_X$ (II) made up solely of chlorides, and a layer having a thickness $5\chi_d$ (III) where ionic concentrations are described according to the classical Gouy-Chapman theory (Gouy, 1910; Chapman, 1913).

In order to determine the electrical potential at the air/water interface, we solve the modified Poisson-Boltzmann equation (equation (9)) and the classical Poisson-Boltzmann equation numerically using the boundary conditions (equations (10), (11), and continuity of the electrical potential between the layers I, II, and III). The electrical potential is calculated using the following equations:

$$\frac{d^2\psi_1}{dx^2} = 0, \quad 0 < x < d_X, \quad (20)$$

$$\frac{d^2\psi_{II}}{dx^2} = \frac{ec_{MX}^b N_a}{\epsilon_w} \exp(e\psi_{II}/k_b T), \quad d_X < x < d_M, \quad (21)$$

$$\frac{d^2\psi_{III}}{dx^2} = \frac{2ec_{MX}^b N_a}{\epsilon_w} \sinh(e\psi_{III}/k_b T), \quad d_M < x < x_{\max} = 5\chi_d, \quad (22)$$

$$\left. \frac{d\psi_I}{dx} \right|_{x=0} = 0, \quad (23)$$

$$\psi_I = \psi_{II} \text{ and } \psi_{II} = \psi_{III} \text{ for } x = d_X \text{ and } x = d_M, \text{ respectively, and} \quad (24)$$

$$\psi_{III} = 0 \text{ at } x = x_{\max}. \quad (25)$$

The surface excess/deficiency of anions and cations are calculated using the electrical potential and equations (7), (18) thus determined. Using the Pitzer parameters (Table 1), we estimate the variation of the activity coefficients with the molality in the bulk water (see equations (A1) to (A11) and equations (B1) to (B5)). Molarities (mol m^{-3}) must also be converted into molalities ($\text{mol kg}_{\text{water}}^{-1}$) (equations (C1) and (C2)). Partial molal volumes of water and ionic species at a given temperature and pressure (1 bar here) were calculated using the ‘‘Thermo-ZNS’’ computer code (Lassin et al., 2005) (Table 2). The surface tension variation with the salinity was determined using equation (19).

A sensitivity analysis is done by varying the value of d_M and keeping d_X constant (Fig. 3a). As expected, the surface tension increases significantly with the distance of closest approach d_M . We observe the same trend by varying the value of d_X and keeping d_M constant (Fig. 3b). However, in this case, the predictions are more sensitive to the distance of closest approach of cations than that of anions, particularly at high salinity.

Only the values of d_X and d_M are optimized in the case of the NaCl electrolyte. We fit the surface tension data with the Simplex algorithm (Caceci and Cacheris, 1984) to obtain the posterior values of the parameters. The algorithm minimizes a cost function R defined in a least square sense. The forward problem, solved by COMSOL MultiphysicsTM 3.4, is coupled with an optimization routine written in a MatLab routine (Fig. 4). We use the a-priori values of Manciu and Ruckenstein (2003) for d_X and d_M , i.e. $d_X = 2 \text{ \AA}$ and $d_M = 5 \text{ \AA}$.

The model predictions of the surface tension variations and the experimental data are in very good agreement for the entire salinity range, especially if the increasing value of the activity coefficient at high salinity is taken into account (Fig. 5). Indeed, using the Pitzer theory (Pitzer and Mayorga, 1973) enables us to predict the experimental data at high salinity ($>3 \text{ mol kg}_{\text{water}}^{-1}$) rather well. Fig.6 presents also the cost function R for different values of d_{Cl} and d_{Na} .

We use a distance of minimum approach for the cations (d_M) of $4.25 \pm 0.5 \text{ \AA}$. This value is similar to the value obtained by Manciu and Ruckenstein (2003) (5 \AA), and corresponds approximately to the effective diameter of sodium ($\sim 4.09 \text{ \AA}$ according to Marcus, 1987, and Shock and Helgeson, 1988) (Table 3). This distance of closest approach corresponds also to the thickness of the “layer” where the water density profile changes (thickness of two water molecules, roughly between 4 \AA and 6 \AA , Jungwirth and Tobias, 2001; Boström et al., 2005). The ion hydration free energy increases when the ion moves into this region (with a lower water density than in the bulk) and structure-making ions, therefore, prefer to stay in the bulk. We use $d_X = 1 \pm 0.25 \text{ \AA}$ for the distance of closest approach of anions. Note that, by using $d_M = 4.1 \text{ \AA}$ and $d_X = 1.5 \text{ \AA}$, we obtain also a very good agreement between the model predictions and the experimental data (see Fig. 6).

Fig. 7a and 7b show the corresponding electrical potential and ionic concentration profiles. The asymmetric ion depletions create a surface electrical potential, even in the

absence of any “initial surface charge” (by comparison with ion complexation processes on solid surfaces, Sverjensky, 2006). The charge generating the diffuse double layer is located a few Ångströms from the surface and is due to the presence of the anions. The resulting electrical potential is negative, and its amplitude at interface increases with the salinity.

In order to check the validity of the previously optimized parameters, we use the same value of d and W for chlorides to predict the surface tension variation of KCl and HCl aqueous solutions (Fig. 8).

3.2. KCl and HCl electrolytes

The model predictions were compared to the surface tension data for KCl and HCl aqueous electrolyte solutions. In our model, chlorides are absent in a layer 1 Å thick and we also assume that the effect of the potassium cation on the structure of water is similar to that of sodium (structure-making ions). Indeed, sodium and potassium have the same electrical charge and their effective electrostatic radii are relatively similar (Table 3). Furthermore, recent molecular dynamics simulations for KNO₃ electrolytes have predicted a thick layer (~ 5 Å) depleted of potassium (Dang et al., 2006).

The calculation of the surface distribution of hydrogen is a particular case because H⁺ cannot exist alone, except in isolation inside a high vacuum. The ionic radius of H⁺ is several orders of magnitude smaller than that of other ions (Li et al., 2008). When hydrogen approaches another atom or molecule, it can distort the other’s electron cloud. H⁺ must be attached to another atom or molecule that possesses a lone pair of electrons. The proton, as an electron acceptor, can be stabilized as in pyramidal hydronium ions (or hydroxonium) H₃O⁺. According to some authors (e.g. Weissenborn and Pugh, 1996; Manciu and Ruckenstein,

2005; Jungwirth and Tobias, 2006), H_3O^+ ions can be positively adsorbed at the air/water interface. Consequently, we consider that these ions are structure-breaking ions.

For the KCl electrolyte, we use equations (20) to (25) to determine the electrical potential at a given salinity, and equations (7), (B1) to (B5), (C1), (C2) and (19) to calculate the surface tension variations with the salinity. For the HCl electrolyte, the electrical potential in the three surface layers is calculated using equation (9) with, respectively, $W_M = -1 k_b T$ (a-priori value from Manciu and Ruckenstein, 2005), $W_X = \infty k_b T$ for $0 < x < d_X$, $W_M = -1 k_b T$ and $W_X = 0 k_b T$ for $d_X < x < d_M$, $W_M = W_X = 0 k_b T$ for $x > d_M$ and the boundary conditions (equations (10) and (11)).

There is a good agreement between the model calculations and the data for the entire salinity range (Figs. 9 and 10). In the case of the KCl electrolyte, we use a distance of closest approach for the cations of $4.5 \pm 0.5 \text{ \AA}$. This corresponds approximately to the effective electrostatic diameter of potassium ($\sim 4.4 \text{ \AA}$ according to Marcus, 1987 and Shock and Helgeson, 1988, Table 3). We also compare the effect of the activity coefficient γ on the surface tension predictions using the Pitzer model (Pitzer and Mayorga, 1973 and Christov and Moller, 2004, for KCl and HCl electrolytes, respectively). For a KCl aqueous solution, the effect of the activity coefficient on the surface tension variations is negligible. This is not the case for the HCl electrolyte whose activity coefficient significantly decreases the predicted surface tension values at high ionic strength ($> 1 \text{ mol kg}_{\text{water}}^{-1}$), and taking into account this correction strongly improves the predictions. For the HCl electrolyte, the Pitzer coefficients are adjusted using an iterative procedure (see Appendix A for more details). Note that we obtain also a good agreement between the model predictions and the data using $d_K = 4.4 \text{ \AA}$ and $d_{\text{Cl}} = 1.5 \text{ \AA}$. In the case of the HCl electrolyte, an increasing value of the distance of closest approach leads to a diminution of the amplitude of the interaction free energy term. For example, the model predictions are in good agreement with the experimental data using

$d_H = 2 \text{ \AA}$ and $W_H = -1 k_b T$ or $d_H = 8 \text{ \AA}$ and $W_H = -0.2 k_b T$. Consequently, we decide to take the value of the effective electrostatic diameter of hydronium for d_H (see Table 3) and we optimize the value of W_H . Finally, we obtain $d_H = 6.2 \text{ \AA}$ and $W_H = -0.3 k_b T$. For HCl, the distances of minimum approach of cations and anions are very close each other and the amplitude of the depth of the potential well is low. Therefore, the “double layer effect” associated with the asymmetric distribution of charges at interface is weak comparing to the NaCl electrolyte. This hypothesis is in agreement with recent molecular dynamics simulations (Mucha et al., 2005): the singlet density profiles of hydronium and chlorides are very close each other and the resulting surface electrostatic potential is weak (Marcelja, 2006). In our approach, we have also neglected ion correlation effects which can be very important at high ionic strength. This assumption may lead to some uncertainties concerning the optimized parameter values.

3.3. NaNO₃, KNO₃, and HNO₃ electrolytes

Chloride anions were then replaced by nitrates. The study of the behaviour of nitrates at the air/water interface for sea salt aerosols is a topic of great interest to atmospheric scientists (Thomas et al., 2007; Brown et al., 2009). Pegram and Record (2006), in their analysis of surface tension data, have suggested that the concentration of nitrate ions at the interface is the same as in the bulk. Computational studies of NO₃⁻ at the aqueous interface have produced conflicting results concerning its distribution. In 2003, molecular dynamics (MD) simulations (polarizable force field) predicted that nitrate anions prefer interfacial to bulk solvation in aqueous solutions (Salvador et al., 2003). Recent MD simulations at high electrolyte concentration (up to 6.8 M of NaNO₃, Thomas et al., 2007), however, predict a decreased interfacial nitrate concentration (Dang et al., 2006; Thomas et al., 2007). Is NO₃⁻

present at the surface of aqueous NaNO_3 solutions? As opposed to chloride, which is a monoatomic ion, nitrate is a polyatomic ion composed of one nitrogen and three oxygen atoms. Thus, nitrate presents an anisotropic molecular polarizability with a large difference between the in-plane and out-of-plane components. The factors stabilizing the bulk versus interfacial solvation of NO_3^- are more complex than those in the case of spherical ions due to multiple hydrogen-bonding sites and also differences in size and shape. The back-and-forth nature of the simulations leaves the debate over the surface propensity of nitrate anions unresolved, and quantitative experimental results have been slow to follow. Brown et al. (2009) used depth resolved X-ray photoemission spectroscopy (XPS) of deliquesced single-crystal samples to provide the most quantitative determination of ion concentration profiles at the liquid-vapor interface of sodium nitrate solutions. They found that the top surface layers of 3 M of NO_3^- solutions are only partially depleted of anions. According to all the studies previously cited, the question regarding the surface propensity of nitrate anions remains unresolved.

Consequently, nitrate is initially assumed to behave like chloride at the air/water interface (a-priori value of $d_X = 1 \text{ \AA}$ and $W_X = \infty k_b T$ in the uppermost layer). We use the same parameter values as previously for the cations, i.e. $d_M = 4.25 \text{ \AA}$ for Na^+ (structure-making), 4.5 \AA for K^+ (structure-making), and 6.2 \AA for H^+ (structure-breaking with $W_M = -0.3 k_b T$).

The model predictions for the surface tension, the activity coefficient, and the corresponding experimental data are shown in Fig. 11, 12 and 13. The model calculations are in good agreement with the experimental data. All the surface tension data for nitrates are lower than those for chlorides, especially at high salinity ($>2 \text{ mol kg}_{\text{water}}^{-1}$). The optimized value for the distance of closest approach of nitrate (d_X) is 0. Table 4 summarizes the values of the optimized parameters for each ionic species. In the case of the KNO_3 electrolyte, the effect of the activity coefficient on the surface tension is very strong and divides it

approximately in two at high salinity ($>3 \text{ mol kg}^{-1}$ water). The activity coefficient significantly influences the surface tension predictions for NaNO_3 , KNO_3 , and HNO_3 electrolytes. This is probably due to the polyatomic nature of nitrate. For HNO_3 electrolyte, the model predictions correspond to the experimental data up to a salinity of $30 \text{ mol kg}_{\text{water}}^{-1}$.

4. CONCLUSION

We have developed a new model to predict the surface tension variation of 1:1 electrolytes at high salinity. According to the Gibbs adsorption equation, the excess/deficiency of the ionic species at the air/water interface and their activities in the bulk water are responsible for the variation of the surface tension with the salinity.

The modified Poisson-Boltzmann and Pitzer equations are used to determine the ionic distributions at interface and the activity coefficients, respectively. The resulting electrochemical model needs only two parameters (two parameters for each ion; four parameters for 1:1 electrolytes) which depend on the effect of ions on water's structure: a length corresponding to the distance of closest approach of structure-making ion or to the thickness of a layer of positive adsorption of structure-breaking ions, and an interaction free energy term (positive for structure-making and negative for structure-breaking ions).

The model predictions are in good agreement with the surface tension data for NaCl , KCl , HCl , NaNO_3 , KNO_3 , and HNO_3 electrolytes. We consider that sodium and potassium are absent in a layer corresponding to the size of their effective electrostatic diameter. The thickness of this layer also corresponds, according to molecular dynamics, to the size of a layer where the density and orientation of water molecules are different from that of the bulk water. Chlorides and nitrates, which have larger ionic radii than sodium and potassium, are assumed to be present at the air/water surface. On the other hand, protons, because of their

very small ionic radii, are complexed with water molecules to form hydronium ions. These ions are considered to be structure-breaking ions and are consequently positively adsorbed at the interface.

The study of the influence of the activity coefficient on the surface tension behaviour is of particular importance in the salinity domain investigated here. Taking the activity coefficient into account leads to a substantial decrease in the surface tension (except for the KCl electrolyte and the case of the NaCl electrolyte where the activity coefficient significantly increases the surface tension at high salinity). This effect is very strong for electrolytes containing hydronium and nitrate. This might be related to the polyatomic nature of hydronium and nitrate.

In future work, we will develop our model to study the surface tension of mixed aqueous solutions, and multivalent ionic species. The model will also include ion correlation effects which can be very important at high ionic strength.

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APPENDIX A

The Pitzer theory (Pitzer, 1973, 1995) enables the determination of the activity coefficients γ_j in the bulk electrolyte. In the case of completely dissociated 1:1 electrolytes and according to Harvie and Weare (1980) and Plummer et al. (1988):

$$\ln \gamma_X^b = \ln \gamma_M^b = \ln \gamma_{MX}^b = F + m_X^b \left[2B_{MX} + (m_M^b + m_X^b)C_{MX} \right] + m_M^b m_X^b C_{MX}, \quad (\text{A1})$$

$$F = -A^\phi \left[\frac{\sqrt{I}}{1 + b\sqrt{I}} + \frac{2}{b} \ln(1 + b\sqrt{I}) \right] + m_M^b m_X^b B'_{MX}, \quad (\text{A2})$$

$$A^\phi = \frac{1}{3} \left(\frac{2\pi N_a d_w}{1000} \right)^{1/2} \left(\frac{e^2}{\epsilon_w k_b T} \right)^{3/2}, \quad (\text{A3})$$

$$z_M m_M^b + m_{H^+}^b = z_X m_X^b + m_{OH^-}^b, \quad (\text{A4})$$

$$I = \frac{1}{2} \sum_j z_j^2 m_j^b = \frac{1}{2} (m_M^b + m_X^b + m_{H^+}^b + m_{OH^-}^b), \quad (\text{A5})$$

where I is the ionic strength (in mol kg_{water}⁻¹), b an empirical parameter with a value of 1.2 and A^ϕ the Debye-Hückel coefficient (0.388 and 0.391 for $T = 20^\circ\text{C}$ and 25°C , respectively, and a pressure of 1 bar using the ‘‘Thermo-ZNS’’ computer code; Lassin et al., 2005). The Debye-Hückel coefficient depends on the density d_w (in kg m⁻³) and the dielectric constant ϵ_w (in F m⁻¹) of pure water. The subscript ‘‘M’’ denotes cations, ‘‘X’’ denotes anions, and N_a is the Avogadro number ($\sim 6.02 \times 10^{23}$ sites mol⁻¹). According to equations (A4) and (A5), the ionic

strength at high salinity can be written $I \cong 0.5(m_M^b + m_X^b) = m_{MX}^b$. B_{MX} and B'_{MX} describe the interaction of pairs of oppositely charged ions, which are defined as explicit functions of the ionic strength by:

$$B_{MX} = \beta_{MX}^0 + \beta_{MX}^1 g(x_1) + \beta_{MX}^2 g(x_2), \quad (A6)$$

$$x_1 = \alpha_1 \sqrt{I} \quad x_2 = \alpha_2 \sqrt{I}, \quad (A7)$$

$$g(x) = 2 [1 - (1+x)e^{-x}] / x^2, \quad (A8)$$

$$B'_{MX} = \beta_{MX}^1 g'(x_1) / I + \beta_{MX}^2 g'(x_2) / I, \quad (A9)$$

$$g'(x) = -2 [1 - (1+x+0.5x^2)e^{-x}] / x^2, \quad (A10)$$

where $\beta_{MX}^0, \beta_{MX}^1, \beta_{MX}^2$ are the Pitzer parameters, and $\alpha_1 = 2$ for 1:1 electrolytes ($\alpha_2 = 0$, except for HNO₃, $\alpha_2 = 1$, Christov, 2005). The term C_{MX} describes ternary short range interactions between two cations and one anion, and one cation and two anions. It is expressed by:

$$C_{MX} = 0.5C_{MX}^\phi. \quad (A11)$$

We, therefore, need four coefficient values, $\beta_{MX}^0, \beta_{MX}^1, \beta_{MX}^2$ and C_{MX}^ϕ , to predict the activity coefficient using the Pitzer theory for 1:1 electrolytes. These values are available in Pitzer databases for many electrolytes (Pitzer, 1973; Pitzer and Mayorga, 1973; Harvie and Weare, 1980; Harvie et al., 1984; Plummer et al., 1988; Felmy et al., 1994; Pitzer, 1995;

Christov and Moller, 2004; Carroll et al., 2005; Christov, 2005; see Table 1 for the electrolytes investigated in our study). In the case of HCl and HNO₃ electrolytes, we use a methodology for determining Pitzer interaction parameters. This methodology is based on the use of PEST (Doherty, 1994), a nonlinear parameter estimation and optimization package, coupled to the PHREEQC code (Parkhurst and Appelo, 1999). PEST takes control of PHREEQC, running it as many times as necessary while adjusting the Pitzer interaction parameters. The fitting procedure stops when the discrepancies between selected model outputs and a complementary set of field or laboratory measurements are reduced to a minimum. The computations were performed using measurements of water activities and mineral solubilities (Hamer and Wu, 1972).

APPENDIX B

The following set of equations describes the variation of the mean activity coefficient for 1:1 electrolytes with the molality (the superscript “*b*” has been deleted for clarity):

$$\frac{d \ln \gamma_{\text{MX}}}{dm_{\text{MX}}} = \frac{d \ln \gamma_{\text{MX}}}{dI} \times \frac{dI}{dm_{\text{MX}}} = \frac{d \ln \gamma_{\text{MX}}}{dI}, \quad (\text{B1})$$

$$\frac{d \ln \gamma_{\text{MX}}}{dI} = \frac{dF}{dI} + 2B_{\text{MX}} + 2I \left(\frac{dB_{\text{MX}}}{dI} + 3C_{\text{MX}} \right), \quad (\text{B2})$$

$$\frac{dF}{dI} = -\frac{A^\phi}{2} \left[\frac{3I^{-1/2}(1+b\sqrt{I})-b}{(1+b\sqrt{I})^2} \right] + 2IB'_{\text{MX}} + I^2 \frac{dB'_{\text{MX}}}{dI}, \quad (\text{B3})$$

$$\frac{dB_{\text{MX}}}{dI} = \sum_{i=1}^2 \frac{\alpha_i \beta_{\text{MX}}^i}{2\sqrt{I}} \left[-\frac{2}{x_i} g(x_i) + \frac{2}{x_i} e^{-x_i} \right], \quad (\text{B4})$$

$$\frac{dB'_{\text{MX}}}{dI} = \sum_{i=1}^2 \frac{\beta_{\text{MX}}^i}{I^2} \left\{ \frac{\alpha_i \sqrt{I}}{2} \left[-\frac{2}{x_i} g'(x_i) - e^{-x_i} \right] - g'(x_i) \right\}. \quad (\text{B5})$$

APPENDIX C

The ionic concentrations c_j in mol m^{-3} are converted to concentrations m_j in $\text{mol kg}_{\text{water}}^{-1}$ by:

$$m_j = \frac{1000c_j}{M_w c_w}, \quad (\text{C1})$$

$$c_w = [10^6 - (c_M V_{Mm} + c_X V_{Xm})] / V_{wm}, \quad (\text{C2})$$

where M_w is the molar mass of water (in g mol^{-1}), c_w the concentration of water (in mol m^{-3}), and V_{im} the standard partial molal volume (in $\text{cm}^3 \text{mol}^{-1}$) of the species.

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Table 1. Pitzer parameters used in this study

Solutions	β_{MX}^0	β_{MX}^1	β_{MX}^2	C_{MX}^ϕ	m_{max}^b
NaCl ^a (25°C)	0.0765	0.2644	0	1.27×10^{-3}	5.50 ^{**}
KCl ^a (25°C)	0.04835	0.2122	0	-8.40×10^{-4}	3.35 ^{**}
HCl ^c (25 °C)	0.1973	0.1691	0	-3.175×10^{-3}	14.98 ^{**}
NaNO ₃ ^b (20°C)	1.443×10^{-3}	0.1861	0	1.212×10^{-4}	9.84 ^{**}
KNO ₃ ^d (25°C)	0.0816	0.0494	0	6.60×10^{-3}	2.63 ^{**}
HNO ₃ ^e (25°C) [*]	-0.1560	0.2847	0.3083	-4.006×10^{-7}	29.09 ^{**}

^a Pitzer and Mayorga (1973)

^b Felmy et al. (1994)

^c by adjusting the Pitzer parameters

^d Carroll et al. (2005)

^e by adjusting the Pitzer parameters

* In the case of HNO₃, we need a four parameters parametrization to predict correctly the activity coefficients ($\alpha_1 = 2$, $\alpha_2 = 0.1$).

** The m_{max}^b values are given by Li and Lu (2001).

Table 2. Standard partial molal volumes of the species calculated with the “Thermo-ZNS” computer code (Lassin et al., 2005)

Species	Volume (cm ³ mol ⁻¹)	
	<i>T</i> = 20°C	<i>T</i> = 25°C
H ₂ O	18.05	18.07
H ⁺	0.00	0.00
Na ⁺	-1.44	-1.13
K ⁺	8.84	9.05
Cl ⁻	17.54	17.68
NO ₃ ⁻	28.25	28.92

Table 3. Ionic and effective electrostatic radii of the species ($T = 25\text{ }^{\circ}\text{C}$)

Ions	Ionic radius (Å)	Electrostatic radius (Å)
H^+ (H_3O^+)	2.14 ^a	3.08 ^a
Na^+	0.97 ^c 1.02 ^b	1.91 ^c 2.18 ^b
K^+	1.33 ^c 1.38 ^b	2.13 ^b 2.27 ^c
Cl^-	1.81 ^c	1.81 ^c
NO_3^-	2.81 ^a	2.81 ^a 2.97 ^c

^a Helgeson et al. (1981)

^b Marcus (1987)

^c Shock and Helgeson (1988)

Table 4. Optimized parameter values for the distance d and the interaction free energy W

	Na ⁺	K ⁺	H ⁺	Cl ⁻	NO ₃ ⁻
d (Å)	4.25 ± 1	4.5 ± 1	6.2 ± 2	1 ± 0.5	0
W (k_bT)	∞	∞	-0.3 ± 0.1	∞	∞

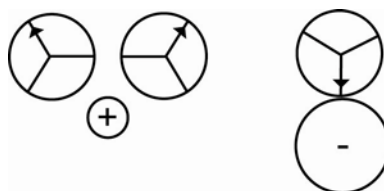


Figure 1. Schematic representation of the structure-making and structure-breaking ions (from Hribar et al., 2002). Structure-making ions (like Na^+) promote hydrogen bonding of neighbouring waters. Structure-breaking ions (like I^-) promote interaction with the neighbouring waters.

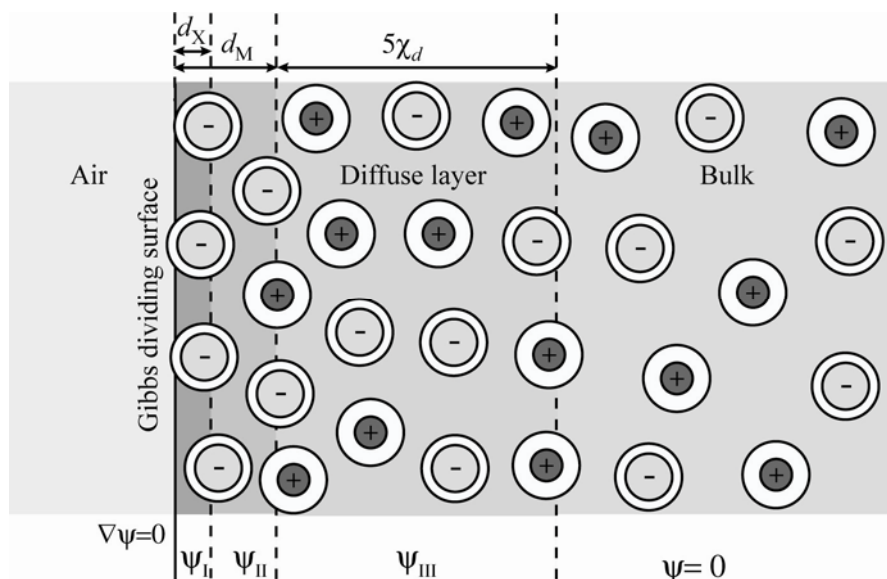


Figure 2. The distribution of ionic species at the air/water interface according to the approach of Manciu and Ruckenstein (2003) (for 1:1 electrolytes). Structure-making ions (cations here, $W_M > 0$) can better organize the water dipoles in bulk water than at the interface. The structure-breaking ions, on the other hand, are pushed toward the interface by the bulk water ($W_X < 0$) because the bulk water molecules can better organize their hydrogen bonding network without these ions. The thicknesses of the interface layers are exaggerated compared to the size of the ions.

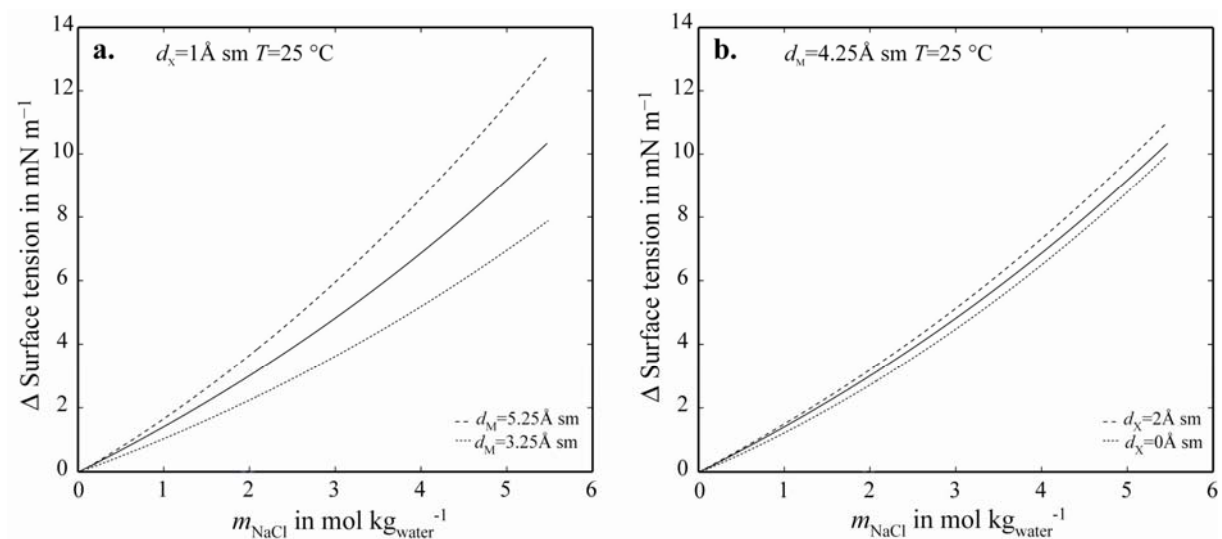


Figure 3. Sensitivity analysis for the prediction of the surface tension of an aqueous NaCl solution (sm: structure-making). The solid lines are the model calculations using $d_X = 1 \text{ \AA}$ and $d_M = 4 \text{ \AA}$.

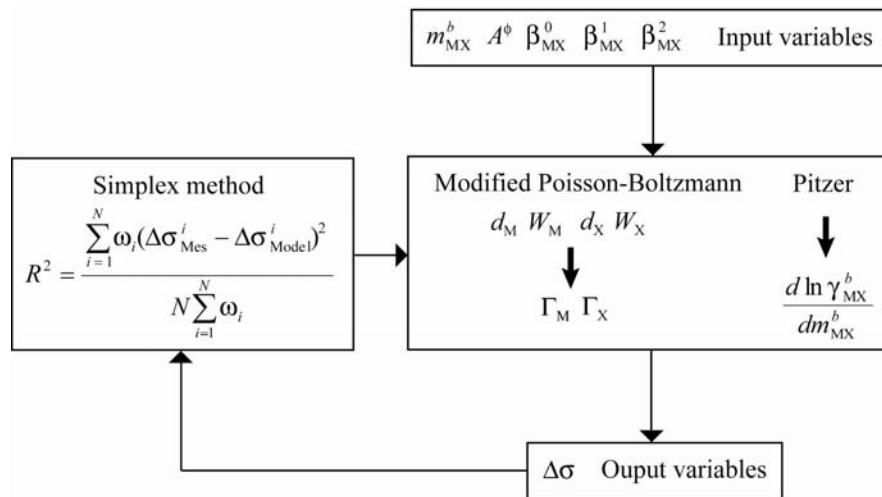


Figure 4. The optimization procedure used to determine the value of d and W for each ionic species. The Simplex algorithm is used to minimize the cost function R^2 . The parameters ω_i are the weight coefficients. By default, $\omega_i = 1$.

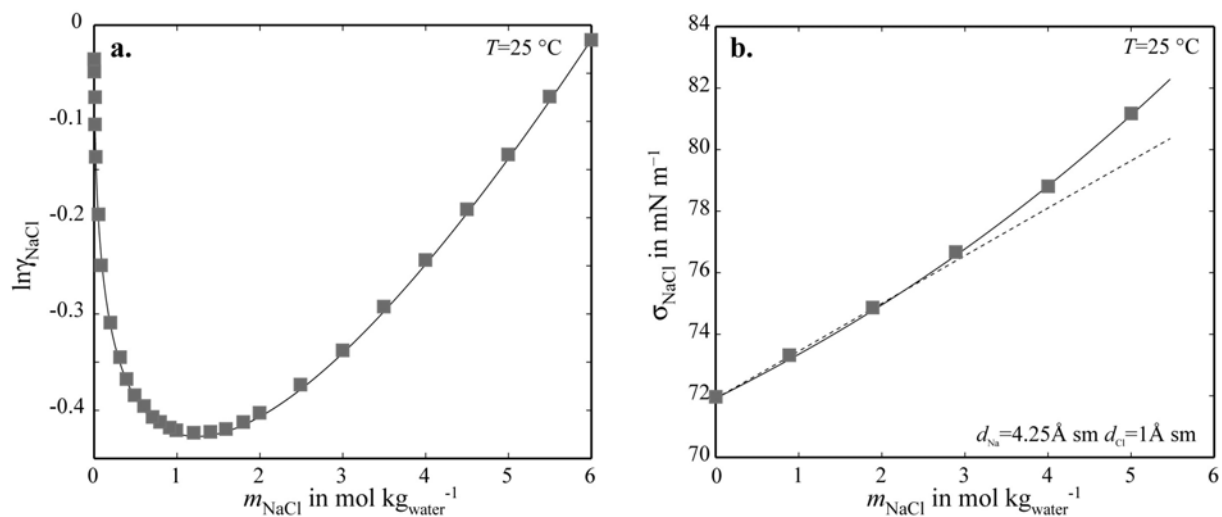


Figure 5. The surface tension (a) and mean activity coefficient (b) of an aqueous NaCl solution versus its molality in the bulk water (“sm”: structure-making). The uncertainties are smaller than the size of the symbols. Symbols correspond to the experimental data of Hamer and Wu (1972) for the activity coefficients and of Abramzon and Gaukhberg (1993) for the surface tension. Plain lines correspond to the model predictions and dashed line to the predictions with $\gamma_{\text{MX}} = 1$.

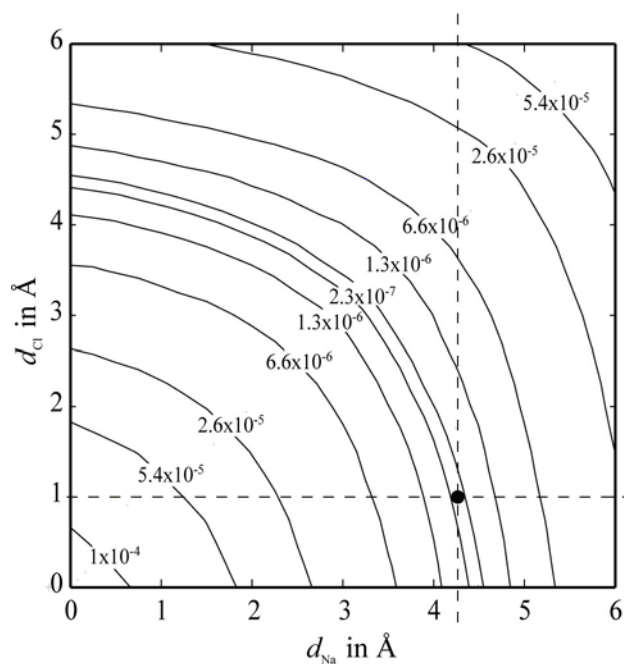


Figure 6. Shape of the cost function R^2 in the case of the NaCl electrolyte surface tension. This function has a unique minimum corresponding to the position of the filled circle. The optimized value of d_{Na} (4.25 Å) is very close to the value of its mean effective electrostatic diameter (4.09 Å, using an arithmetic mean of the values of Marcus, 1987 and Shock and Helgeson, 1988; see Table 3).

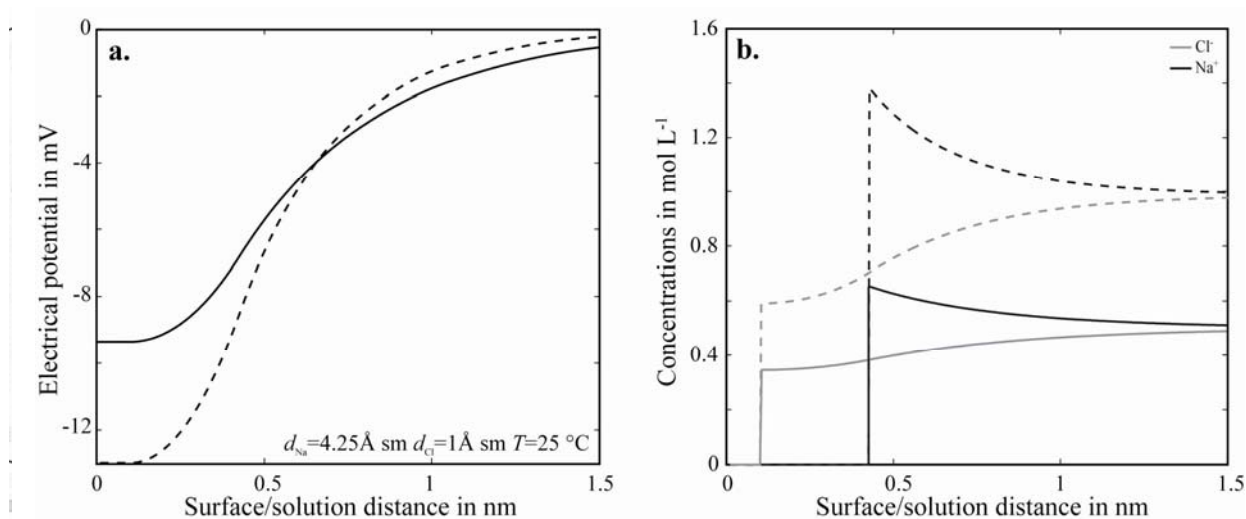


Figure 7. Electrical potential at the air/water interface from the model in the case of NaCl electrolytes using the parameters described in Fig. 5a (a) and the corresponding concentration profiles of chlorides and sodiums (b) (the plain and dashed line correspond to a molarity of 0.5 and 1 mol L⁻¹, respectively). The deficit of ionic species at interface are responsible for the observed surface tension at the air/water interface.

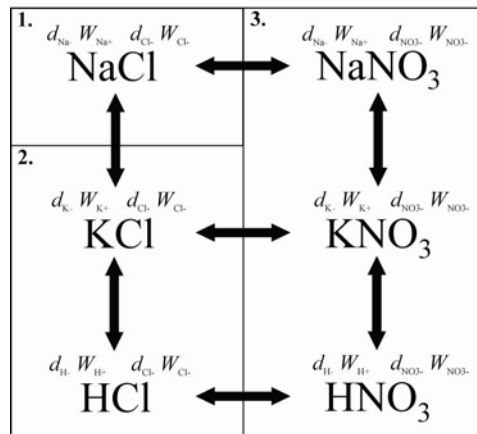


Figure 8. The optimization procedure used to determine the value of d and W for each ionic species and for different electrolytes. At first, the Simplex algorithm is used to optimize the values of d and W for NaCl-electrolyte. Secondly, the optimized values of d_{Cl^-} and W_{Cl^-} are utilized to constrain the parameter values for KCl and HCl electrolytes. The gradient descent method is used to optimize d_{K^+} and the Simplex to estimate d_{H^+} and W_{H^+} . Thirdly, the optimized value of d_{Na^+} constrains the parameter values for NaNO₃ electrolyte. Finally, all the optimized parameter values for K^+ , H^+ , and NO_3^- are used to predict the surface tension variation for KNO₃ and HNO₃ electrolytes.

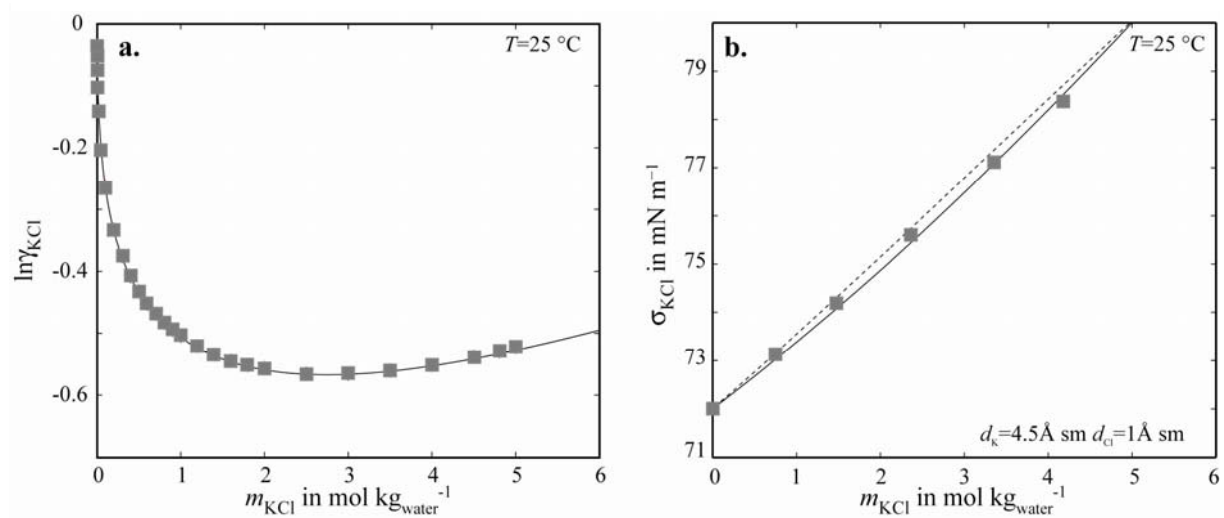


Figure 9. The surface tension (a) and mean activity coefficient (b) of an aqueous KCl solution versus its molality in the bulk water.

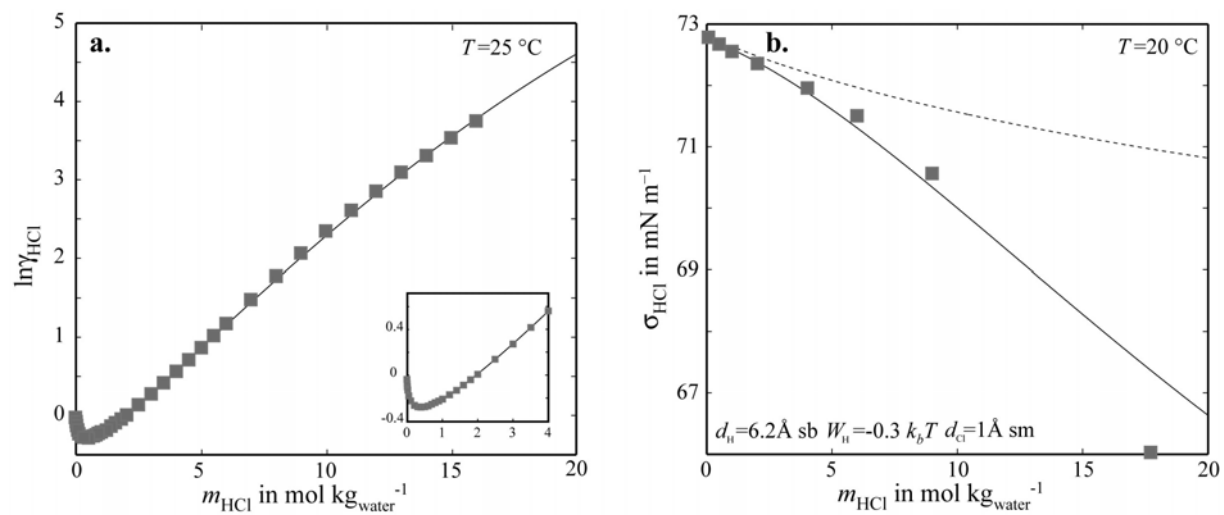


Figure 10. The surface tension (a) and mean activity coefficient (b) of an aqueous HCl solution versus its molality in the bulk water (“sm”: structure-making; “sb”: structure-breaking).

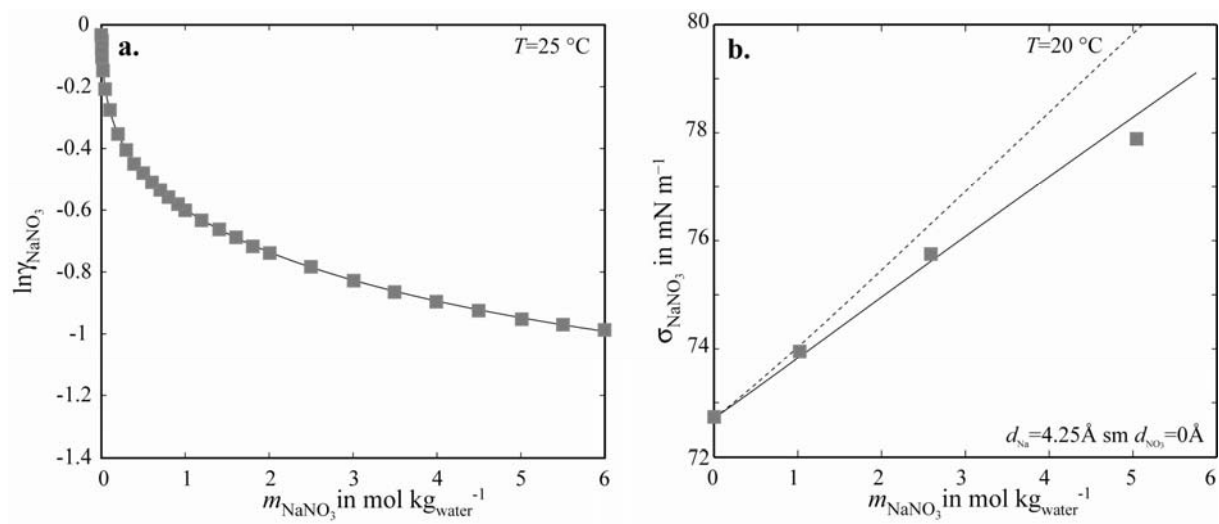


Figure 11. The surface tension (a) and mean activity coefficient (b) of an aqueous NaNO₃ solution versus its molality in the bulk water.

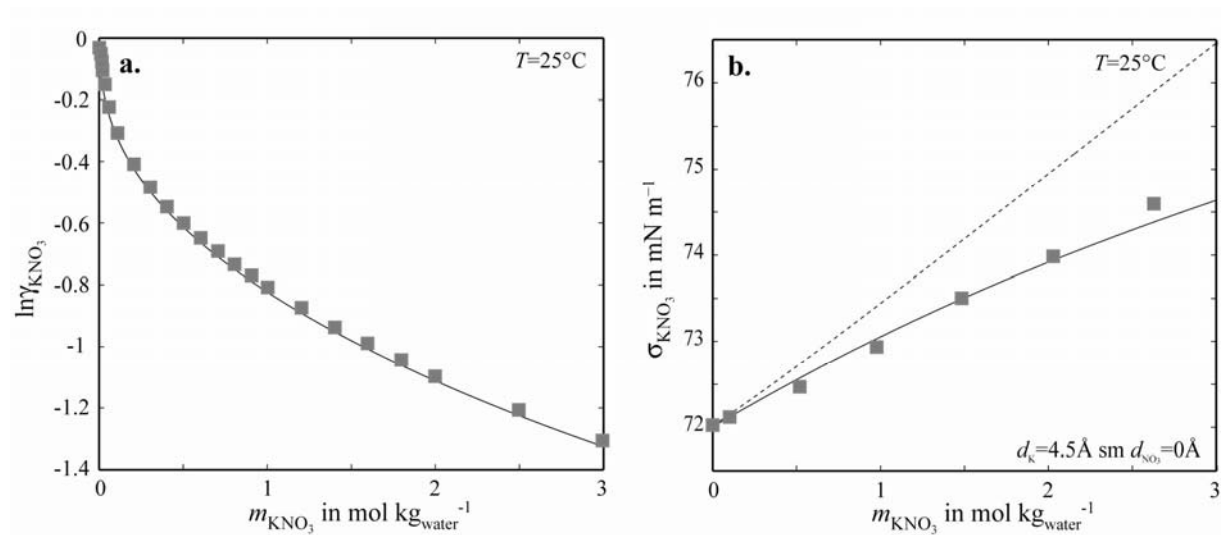


Figure 12. The surface tension (a) and mean activity coefficient (b) of an aqueous KNO₃ solution versus its molality in the bulk water.

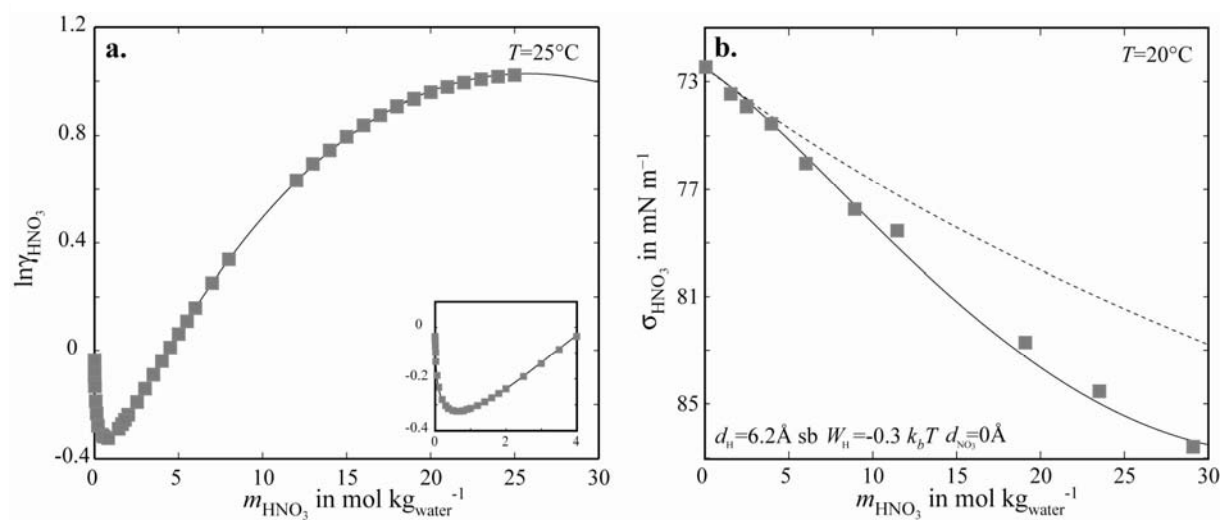


Figure 13. The surface tension (a) and mean activity coefficient (b) of an aqueous HNO_3 solution versus its molality in the bulk water (sm: structure-making; sb: structure-breaking).